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Efficient Multistep Algorithms for First-Order IVPs with Oscillating Solutions: II Implicit and Predictor–Corrector Algorithms

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Abstract: This research introduces a fresh methodology for creating efficient numerical algorithms to solve first-order Initial Value Problems (IVPs). The study delves into the theoretical foundations of these methods and demonstrates their application to the Adams–Moulton technique in a five-step process. We focus on developing amplification-fitted algorithms with minimal phase-lag or phase-lag equal to zero (phase-fitted). The request of amplification-fitted (zero dissipation) is to ensure behavior like symmetric multistep methods (symmetric multistep methods are methods with zero dissipation). Additionally, the stability of the innovative algorithms is examined. Comparisons between our new algorithm and traditional methods reveal its superior performance. Numerical tests corroborate that our approach is considerably more effective than standard methods for solving IVPs, especially those with oscillatory solutions.



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1. Introduction

Equations or systems of Equations of the form

$$\mathbf{s}'(\mathbf{t}) = \mathbf{q}(\mathbf{t}, \mathbf{s}), \quad \mathbf{s}(\mathbf{t}_0) = \mathbf{s}_0 \quad (1)$$

are utilized across a range of fields such as astrophysics, chemistry, physics, electronics, nanotechnology, materials science, and more. Equations with oscillatory or periodic solutions are of particular interest (refer to [1,2]).

Extensive research has been dedicated to studying the numerical solutions for the aforementioned equation or system of equations over the past two decades (see, for instance [3–11], and references therein). For a detailed examination of techniques for solving (1) with oscillating solutions, refer to [3,7,12], as well as the works of Quinlan and Tremaine [5,6,8,13], among others. Most existing numerical methods for solving (1) share a common feature of being multistep or hybrid approaches and were primarily developed for second-order differential equations. Some of the key method categories and their bibliography include:

- Exponentially-fitted, Trigonometrically-Fitted, Phase-Fitted, and Amplification-Fitted Runge–Kutta and Runge–Kutta–Nyström Methods with minimal phase-lag (refer to [14–49]);

- Exponentially-fitted and Trigonometrically-Fitted Phase-Fitted and Amplification-Fitted Multistep Methods and Multistep Methods with minimal phase-lag (refer to [50–114]).

Recently, Simos [115] developed the theory for the development of multistep methods with minimal phase-lag or phase-fitted multistep methods for first-order IVPs. More specifically, he developed the theory for computing the phase-lag and amplification error of multistep methods for first-order IVPs. In this paper, we will extend this theory to the implicit multistep methods for first-order IVPs. This paper introduces the theory.

The rest of the paper is structured as follows:

- Section 2 presents the general theory for calculating the phase-lag and amplification error of implicit multistep methods for first-order IVPs. In this section, we produce the direct formulae for the calculation of the phase-lag and amplification factor;
- Section 3 introduces the methodologies and the methods which will be developed in Sections 4–14. In this section, we present the methodologies for the development of efficient multistep methods for first-order initial value problems;
- Section 4 introduces the Adams–Bashforth five-step method and presents the methodology for the minimization of the phase-lag. In this section, we present the explicit Adams–Bashforth five-step method and we study its phase-lag and amplification error;
- Section 5 presents the development of the amplification-fitted Adams–Bashforth five-step method of fourth algebraic order with phase-lag of order four. Based on the theory developed in Section 2, we eliminate the amplification error and we calculate the coefficients of the method in order for the method to have phase-lag of order four;
- Section 6 presents the development of the amplification-fitted Adams–Bashforth five-step method of third algebraic order with phase-lag of order six. Based on the theory presented in Section 2, we eliminate the amplification error and we calculate the coefficients of the method in order for the method to have phase-lag of order six.
- Section 7 presents the development of the amplification-fitted Adams–Bashforth five-step method of fourth algebraic order. We calculate the coefficients of the method, in order for its amplification error to be equal to zero;
- Section 8 presents the development of the amplification-fitted Adams–Bashforth five-step method of fourth algebraic order. We calculate the coefficients of the method, in order for its phase-lag and amplification error to be eliminated;
- Section 9 introduces the Adams–Moulton five-step method and presents the methodology for the minimization of the phase-lag. In this section, we present the implicit Adams–Moulton five-step method and we study its phase-lag and amplification error;
- Section 10 presents the development of the amplification-fitted Adams–Moulton five-step method of fifth algebraic order with phase-lag of order four. Based on the theory developed in Section 2, we eliminate the amplification error and we calculate the coefficients of the method in order for the method to have phase-lag of order four;
- Section 11 presents the development of the amplification-fitted Adams–Moulton five-step method of second algebraic order with phase-lag of order six. Based on the theory presented in Section 2, we vanish the amplification error and we calculate the coefficients of the method in order for the method to have phase-lag of order six;
- Section 12 presents the development of the amplification-fitted Adams–Moulton five-step method of second algebraic order with phase-lag of order eight. Based on the theory presented in Section 2, we demand the amplification error to be equal to zero and we calculate the coefficients of the method in order for the method to have phase-lag of order eight;
- Section 13 presents the development of the amplification-fitted Adams–Moulton five-step method of fifth algebraic order. We calculate the coefficients of the method, in order for its amplification error to be equal to zero;
- Section 14 presents the development of the amplification-fitted and phase-fitted Adams–Moulton five-step method of fifth algebraic order. We calculate the coefficients of the method, in order for its phase-lag and amplification error to be eliminated;

- Section 15 discusses a stability analysis for the newly proposed methods in Sections 4–14. We examine the stability of the developed methods for several values of v .
- Section 16 presents numerical results. We examine the efficiency of the proposed methods in their application on seven well-known problems. For each problem, we give conclusion for the behavior of the developed methods;
- Finally, Section 17 presents the conclusions of this research.

The numerical results demonstrate that the methodology for developing phase-fitted and amplification-fitted multistep methods has yielded the most effective solutions for problems with oscillating behavior.

2. The Theory

Using the following scalar test equation, we can examine the phase-lag of multistep approaches for the problems (1).

$$s'(t) = I \omega s(t). \quad (2)$$

This problem is solved by using the following formula:

$$s(t) = \exp(I \omega t). \quad (3)$$

Taking into consideration the multistep approaches that enable the numerical solution of the problem that was discussed before (1):

$$s_{n+k} - s_{n+k-1} = h \sum_{j=0}^{k-1} [A_{n+k-j}(\omega h) q_{n+k-j}], \quad (4)$$

where $A_{n+k-j}(\omega h)$, $j = 1, 2, \dots, k$ are polynomials of ωh and h is the step length of the integration.

The following result is obtained by applying Equations (2)–(4):

$$s_{n+k} - s_{n+k-1} = I \omega h \sum_{j=0}^{k-1} [A_{n+k-j}(\omega h) s_{n+k-j}]. \quad (5)$$

While considering:

$$vs. = \omega h, \quad (6)$$

(5) gives:

$$s_{n+k} - s_{n+k-1} = I vs. \sum_{j=0}^{k-1} [A_{n+k-j}(v) s_{n+k-j}], \quad (7)$$

and

$$(1 - I v A_{n+k}(v)) s_{n+k} - (1 + I v A_{n+k-1}(v)) s_{n+k-1} - I vs. \sum_{j=2}^{k-1} [A_{n+k-j}(v) s_{n+k-j}] = 0. \quad (8)$$

What follows is the characteristic equation of the difference equation that was mentioned in (8):

$$[1 - I v A_{n+k}(v)] \lambda^k - [1 + I v A_{n+k-1}(v)] \lambda^{k-1} - I vs. \sum_{j=2}^{k-1} [A_{n+k-j}(v) \lambda^{k-j}] = 0. \quad (9)$$

Definition 1. Considering that the theoretical solution of the scalar test Equation (2) for $t = h$ is $\exp(I \omega h)$, which can also be written as $\exp(I v)$ (refer to (6)), and the numerical solution of the scalar test Equation (2) for $t = h$ is $\exp(I \theta(vs.))$, we can define the phase-lag as follows:

$$\Phi = vs. - \theta(vs.). \quad (10)$$

Assuming $v \rightarrow 0$, the phase-lag order is q if and only if $\Phi = O(v^{q+1})$.

Considering the following:

$$\lambda^n = \exp^{nI\theta(v)} = \cos[n\theta(v)] + I \sin[n\theta(v)] \quad n = 1, 2, \dots, \quad (11)$$

we obtain:

$$\begin{aligned} & (1 - Iv A_{n+k}) \left\{ \cos[k\theta(v)] + I \sin[k\theta(v)] \right\} \\ & - (1 + Iv A_{n+k-1}) \left\{ \cos[(k-1)\theta(v)] + I \sin[(k-1)\theta(v)] \right\} \\ & - \sum_{j=2}^{k-1} Iv A_{n+k-j}(v) \left\{ \cos[(k-j)\theta(v)] + I \sin[(k-j)\theta(v)] \right\} = 0. \end{aligned} \quad (12)$$

To understand the connection mentioned above (12), you must use the next lemmas.

Lemma 1. *The following relations hold:*

$$\cos[\theta(v)] = \cos(vs.) + c v^{q+2} + O(v^{q+4}). \quad (13)$$

$$\sin[\theta(v)] = \sin(vs.) - c v^{q+1} + O(v^{q+3}). \quad (14)$$

For the proof, see [115].

Lemma 2. *This relation holds:*

$$\cos[j\theta(v)] = \cos(j vs.) + c j^2 v^{q+2} + O(v^{q+4}). \quad (15)$$

$$\sin[j\theta(v)] = \sin(j vs.) - c j v^{q+1} + O(v^{q+3}). \quad (16)$$

For the proof, see [115].

When the relations (15) and (16) are considered, relation (12) shifts to:

$$\begin{aligned} & (1 - Iv A_{n+k}) \left\{ \cos[k vs.] + c k^2 v^{q+2} + I \left[\sin[k vs.] - c k v^{q+1} \right] \right\} \\ & - (1 + Iv A_{n+k-1}) \left\{ \left[\cos[(k-1) vs.] + c (k-1)^2 v^{q+2} \right] \right. \\ & \quad \left. + I \left[\sin[(k-1) vs.] - c (k-1) v^{q+1} \right] \right\} \\ & - \sum_{j=2}^{k-1} Iv A_{n+k-j}(v) \left\{ \left[\cos[(k-j) vs.] + c (k-j)^2 v^{q+2} \right] \right. \\ & \quad \left. + I \left[\sin[(k-j) vs.] - c (k-j) v^{q+1} \right] \right\} = 0. \end{aligned} \quad (17)$$

It is possible to split connection (17) into two halves, the real and the imaginary.

The Real Part

The real part gives:

$$\begin{aligned} & \cos[k vs.] + c k^2 v^{q+2} + v A_{n+k} \left[\sin[k vs.] - c k v^{q+1} \right] \\ & - \cos[(k-1) vs.] - c (k-1)^2 v^{q+2} \\ & + v A_{n+k-1} \left[\sin[(k-1) vs.] - c (k-1) v^{q+1} \right] \\ & + \sum_{j=0}^{k-1} v A_{n+k-j}(v) \left[\sin[(k-j) vs.] - c (k-j) v^{q+1} \right] = 0. \end{aligned} \quad (18)$$

Relation (18) gives:

$$\begin{aligned} \cos[k \text{ vs.}] - \cos[(k-1) \text{ vs.}] + v \sum_{j=0}^{k-1} A_{n+k-j}(v) \sin[(k-j) \text{ vs.}] \\ = -c v^{q+2} \left[k^2 - (k-1)^2 - \sum_{j=0}^{k-1} (k-j) A_{n+k-j}(v) \right] \Rightarrow \\ -c v^{q+2} = \frac{\cos[k \text{ vs.}] - \cos[(k-1) \text{ vs.}] + v \sum_{j=0}^{k-1} A_{n+k-j}(v) \sin[(k-j) \text{ vs.}]}{2k - 1 - \sum_{j=0}^{k-1} A_{n+k-j}(v) (k-j)}. \end{aligned} \quad (19)$$

According to technique(4), this is the direct formula for calculating the phase-lag of the multistep approach. We will outline the steps to calculate the phase-lag of technique (4) below.

The Imaginary Part

The imaginary part gives:

$$\begin{aligned} \sin[k \text{ vs.}] - c k v^{q+1} - v A_{n+k} \left[\cos[k \text{ vs.}] + c k^2 v^{q+2} \right] \\ - \sin[(k-1) \text{ vs.}] + c (k-1) v^{q+1} \\ - v A_{n+k-1} \left[\cos[(k-1) \text{ vs.}] + c (k-1)^2 v^{q+2} \right] \\ - \sum_{j=0}^{k-1} v A_{n+k-j}(v) \left[\cos[(k-j) \text{ vs.}] + c (k-j)^2 v^{q+2} \right] = 0. \end{aligned} \quad (20)$$

Relation (20) gives:

$$\begin{aligned} \sin[k \text{ vs.}] - \sin[(k-1) \text{ vs.}] - v \sum_{j=0}^{k-1} A_{n+k-j}(v) \cos[(k-j) \text{ vs.}] \\ = -c v^{q+1} \left[-1 - v^2 \sum_{j=0}^{k-1} A_{n+k-j}(v) (k-j)^2 \right] \Rightarrow \\ -c v^{q+1} = \frac{\sin[k \text{ vs.}] - \sin[(k-1) \text{ vs.}] - v \sum_{j=0}^{k-1} A_{n+k-j}(v) \cos[(k-j) \text{ vs.}]}{-1 - v^2 \sum_{j=0}^{k-1} A_{n+k-j}(v) (k-j)^2}. \end{aligned} \quad (21)$$

In the multistep technique (4), this is the straightforward approach to calculating the amplification factor.

Definition 2. We refer to the method with eliminated phase-lag as the **phase-fitted** method.

Definition 3. We refer the method with eliminated amplification factor as the **amplification-fitted** method.

3. Procedures for the Methodologies for Achieving the Minimum Phase-Lag, Minimum Amplification Factor, Phase-Fitted, and Amplification-Fitted

In the following sections, we will present several procedure for:

- Procedures for the methodologies for achieving the minimum phase-lag;
- Procedures for the methodologies for achieving the minimum amplification factor;
- Procedures for the methodologies for achieving phase-fitted and amplification-fitted algorithms.

Methodologies for the Development of the Newly Introduced Methods

We can divided the methodologies for the development of efficient multistep methods into the following categories:

- Methods with minimization of the phase-lag (see Sections 5, 6, 10–12);
- Amplification-fitted methods (see Sections 7 and 13);
- Phase-fitted and amplification-fitted methods (see Sections 8 and 14).

4. Explicit Method: Adams–Bashforth Five-Step Method

In particular, we shall illustrate the famous Adams–Bashforth approach of fourth algebraic order, which is the following:

$$s_{n+1} - s_n = \frac{h}{24} (55 s'_n - 59 s'_{n-1} + 37 s'_{n-2} - 9 s'_{n-3}), \quad (22)$$

together with the local truncation error (*LTE*) provided by:

$$LTE = \frac{251}{720} h^5 s^{(5)}(t) + O(h^6). \quad (23)$$

We use the theory from Section 2 to obtain the method's phase-lag and amplification error.

The difference Equation (7) with $k = 4$ is obtained by applying algorithm (22) to the test Equation (2) with:

$$A_4(vs.) = 0, A_3(vs.) = \frac{55}{24}, A_2(vs.) = -\frac{59}{24}, A_1(vs.) = \frac{37}{24}, A_0(vs.) = -\frac{3}{8}. \quad (24)$$

By applying the Taylor series expansion to the above Equation (19) and setting $m = 1(1)4$, we can obtain the following:

$$\begin{aligned} \frac{\cos(4 vs.) - \cos(3 vs.) + v \sum_{j=0}^4 A_{k-j}(v) \sin[(k-j) vs.]}{2k-1 - \sum_{j=0}^4 A_{k-j}(v)(k-j)} &= \\ &- \frac{977}{5040} v^6 + \frac{611}{4032} v^8 + \dots \end{aligned} \quad (25)$$

Consequently, $q = 4$ and $c = -\frac{977}{5040}$. The fourth algebraic order Adams–Bashforth method is of fourth order phase-lag.

By applying the Taylor series expansion to the above Equation (21) and setting $m = 1(1)4$, we can obtain the following:

$$\begin{aligned} \frac{\sin(4 vs.) - \sin(3 vs.) - v \sum_{j=0}^4 A_{k-j}(v) \cos[(k-j) vs.]}{-1 - v^2 \sum_{j=0}^4 A_{k-j}(v)(k-j)^2} &= \\ &- \frac{251}{720} v^5 + \frac{151,577}{30,240} v^7 + \dots \end{aligned} \quad (26)$$

Consequently, $q = 4$ and $c = -\frac{251}{720}$. The Adams–Bashforth approach, which is of fourth order algebraic, has an amplification error of the same order. We will refer to the fourth algebraic order Adams–Bashforth algorithm as **Algorithm I** for our computing purposes.

4.1. Minimal Phase-Lag

We examine the following basic five-step algorithm to learn more about the methodologies to minimize the phase-lag:

$$s_{n+1} - s_n = h \left(K_0(vs.), s'_n + K_1(vs.) s'_{n-1} + K_2(vs.) s'_{n-2} + K_3(vs.) s'_{n-3} \right). \quad (27)$$

Procedure to Minimize the Phase-Lag

Below is the procedure that minimizes the phase-lag:

- Eliminating the amplification factor;
- Phase-lag calculation using the coefficient acquired in the preceding stage;
- Expanding the phase-lag calculated before using a Taylor series;
- Defining the set of equations that minimizes the phase-lag;
- Calculation of the updated coefficients.

The following two phase-lag-minimizing algorithms are derived from the aforementioned procedure.

5. Amplification-Fitted Method of Fourth Algebraic Order with Phase-Lag of Order Four

Let us consider method (27) with $K_1(vs.) = -\frac{59}{24}$, $K_3(vs.) = -\frac{9}{24}$.

5.1. Eliminating the Amplification Factor

We obtain the following result when we use the straightforward approach for computing the amplification factor (21):

$$AF = \frac{\sin(4v) - \sin(3v) - K_0(vs.)v \cos(3v) + \frac{59}{24} \cos(2v)v - K_2(vs.)v \cos(vs.) + \frac{3}{8}v}{-v^2 K_2(vs.) - 9v^2 K_0(vs.) + \frac{59}{6}v^2 - 1}, \quad (28)$$

where AF denotes the amplification factor.

Assuming that the amplification factor must be eliminated, or that $AF = 0$, we derive:

$$K_0(vs.) = -\frac{1}{24} \frac{24K_2v \cos(vs.) - 59v \cos(2v) - 24 \sin(4v) + 24 \sin(3v) - 9v}{v \cos(3v)}. \quad (29)$$

5.2. Procedure for Minimizing the Phase-Lag

We can obtain the phase-lag by plugging the values of $K_0(vs.)$ and $K_2(vs.)$ that were previously provided into the direct formula for calculating it (19):

$$PhErr = -\frac{v \left(18 \sin(vs.) (\cos(vs.))^2 v - 24 \sin(vs.) \cos(vs.) v K_2(vs.) + 25 v \sin(vs.) + 12 \cos(vs.) - 12 \right)}{\Xi_1(vs.)}, \quad (30)$$

where

$$\begin{aligned} \Xi_1(vs.) &= 48(\cos(vs.))^3 v K_2(vs.) + 288 \sin(vs.) (\cos(vs.))^3 - 572 (\cos(vs.))^3 vs. \\ &- 144 \sin(vs.) (\cos(vs.))^2 + 177 v (\cos(vs.))^2 - 72 K_2(vs.) v \cos(vs.) \\ &- 144 \sin(vs.) \cos(vs.) + 429 v \cos(vs.) + 36 \sin(vs.) - 75 v, \end{aligned} \quad (31)$$

and $PhErr$ denotes the phase-lag.

By applying the Taylor series expansion to the Formula (30), we are able to retrieve:

$$\begin{aligned} PhErr &= -\frac{(37 - 24 K_2(vs.))v^2}{-24 K_2(vs.) - 5} \\ &- \frac{v^4}{-24 K_2(vs.) - 5} \left[-\frac{74}{3} + 16 K_2(vs.) - \frac{-37 + 24 K_2(vs.)}{24 K_2(vs.) + 5} \left(-36 K_2(vs.) + \frac{489}{2} \right) \right] \\ &- \frac{v^6}{-24 K_2(vs.) - 5} \left[\frac{1121}{120} - \frac{16 K_2(vs.)}{5} - \frac{-37 + 24 K_2(vs.)}{24 K_2(vs.) + 5} \left(39 K_2(vs.) - \frac{7573}{40} \right) \right] \\ &+ \frac{1}{6} \frac{7488 K_2(vs.)^2 - 46,272 K_2(vs.) + 53,539}{(24 K_2(vs.) + 5)^2} \left(-36 K_2(vs.) + \frac{489}{2} \right) + \dots \end{aligned} \quad (32)$$

Requiring the phase-lag to be minimized, we obtain the equation mentioned below:

$$-\frac{(37 - 24 K_2(vs.))v^2}{-24 K_2(vs.) - 5} = 0 \implies K_2(vs.) = \frac{37}{24}. \quad (33)$$

This novel algorithm has the following features:

$$\begin{aligned} K_0(vs.) &= \frac{1}{24} \frac{\Xi_2(vs.)}{v \cos(3v)}, \\ K_1(vs.) &= -\frac{59}{24}, \\ K_2(vs.) &= \frac{37}{24}, \\ K_3(vs.) &= -\frac{9}{24}, \\ LTE &= \frac{251}{720} h^5 \left(s^{\{(5)\}}(t) - \omega^4 s'(t) \right) + O(h^6) \\ PhErr &= \frac{529}{5040} v^6 + \frac{20,551}{47,040} v^8 + \dots, \\ AF &= 0, \end{aligned} \quad (34)$$

where

$$\begin{aligned} \Xi_2(vs.) &= 192 \sin(vs.) (\cos(vs.))^3 - 96 \sin(vs.) (\cos(vs.))^2 \\ &+ 118 v (\cos(vs.))^2 - 96 \sin(vs.) \cos(vs.) \\ &- 37 v \cos(vs.) + 24 \sin(vs.) - 50 v. \end{aligned} \quad (35)$$

$K_0(vs.)$ may be expressed as a Taylor series expansion:

$$K_0(vs.) = \frac{55}{24} + \frac{251 v^4}{720} + \frac{647 v^6}{756} + \dots \quad (36)$$

From a computational standpoint, we shall refer to the aforementioned new technique as **Algorithm II**.

Remark 1. If we chose three free parameters (for example K_j , $j = 0(1)2$), the resulting algorithm will be the same as above.

6. Amplification-Fitted Method of Third Algebraic Order with Phase-Lag of Order Six

Let us consider method (27) with the parameter $K_j(vs.)$, $j = 0(1)3$ free.

For the development of the method, see Appendix A.

This novel algorithm has the following features:

$$\begin{aligned} K_0(vs.)(vs.) &= \frac{\Psi_{12}(vs.)}{720 v \cos(3v)}, \\ K_1(vs.) &= -\frac{179}{288}, \\ K_2(vs.) &= \frac{13}{180}, \\ K_3(vs.) &= -\frac{11}{1440}, \\ LTE &= \frac{529}{1440} h^4 \left(s^{\{(3)\}}(t) - \omega^2 s'(t) \right) + O(h^5), \\ PhErr &= -\frac{191}{423,360} v^8 - \frac{426,379}{237,081,600} v^{10} + \dots, \\ AF &= 0, \end{aligned} \quad (37)$$

where

$$\begin{aligned}\Psi_{12}(vs.) &= 5760 \sin(vs.) (\cos(vs.))^3 - 2880 \sin(vs.) (\cos(vs.))^2 \\ &+ 895 v (\cos(vs.))^2 - 2880 \sin(vs.) \cos(vs.) \\ &- 52 v \cos(vs.) + 720 \sin(vs.) - 442 v.\end{aligned}\quad (38)$$

$K_0(vs.)(vs.)$ may be expressed as a Taylor series expansion:

$$K_0(vs.)(vs.) = \frac{1121}{720} - \frac{529}{1440} v^2 + \frac{41}{3456} v^4 - \frac{8623}{3,628,800} v^6 + \dots\quad (39)$$

From a computational standpoint, we shall refer to the aforementioned new technique as **Algorithm III**.

7. Amplification-Fitted Method of Fourth Algebraic Order

Let us consider method (27) with $K_1(vs.) = -\frac{59}{24}$, $K_2(vs.) = \frac{37}{24}$, $K_3(vs.) = -\frac{9}{24}$

7.1. Eliminating the Amplification Factor

We obtain the following result when we use the straightforward approach for computing the amplification factor (21):

$$AF = \frac{\sin(4v) - \sin(3v) - K_0(vs.) v \cos(3v) + \frac{59}{24} v \cos(2v) - \frac{37}{24} v \cos(vs.) + \frac{3}{8} v}{-9 v^2 K_0(vs.) + \frac{199}{24} v^2 - 1}, \quad (40)$$

where AF denotes the amplification factor.

Assuming that the amplification factor must be eliminated, or that $AF = 0$, we derive:

$$K_0(vs.) = \frac{1}{24} \frac{59 v \cos(2v) - 37 v \cos(vs.) + 24 \sin(4v) - 24 \sin(3v) + 9v}{v \cos(3v)}. \quad (41)$$

Phase-Lag of the Method

We can obtain the phase-lag by plugging the value of $K_0(vs.)$ that was previously provided into the direct formula for calculating it (19):

$$PhErr = -\frac{1}{3} \frac{v \Psi_{13}(vs.)}{\Psi_{14}(vs.)}. \quad (42)$$

where

$$\begin{aligned}\Psi_{13}(vs.) &= 18 (\cos(vs.))^2 \sin(vs.) v - 37 \cos(vs.) \sin(vs.) v \\ &+ 25 v \sin(vs.) + 12 \cos(vs.) - 12, \\ \Psi_{14}(vs.) &= 96 \sin(vs.) (\cos(vs.))^3 - 166 v (\cos(vs.))^3 \\ &- 48 (\cos(vs.))^2 \sin(vs.) + 59 v (\cos(vs.))^2 \\ &- 48 \cos(vs.) \sin(vs.) + 106 v \cos(vs.) + 12 \sin(vs.) - 25 v,\end{aligned}\quad (43)$$

and $PhErr$ denotes the phase-lag.

By applying the Taylor series expansion to the Formula (42), we are able to retrieve:

$$PhErr = \frac{529}{5040} v^6 + \frac{20,551}{47,040} v^8 + \dots\quad (44)$$

This novel algorithm has the following features:

$$\begin{aligned}
 K_0(vs.) &= \frac{1}{24} \frac{59v\cos(2v) - 37v\cos(vs.) + 24\sin(4v) - 24\sin(3v) + 9v}{v\cos(3v)}, \\
 K_1(vs.) &= -\frac{59}{24}, \\
 K_2(vs.) &= \frac{37}{24}, \\
 K_3(vs.) &= -\frac{9}{24}, \\
 LTE &= \frac{251}{720}h^5 \left(s^{\{(5)\}}(t) - \omega^4 s'(t) \right) + O(h^6), \\
 PhErr &= \frac{529}{5040}v^6 + \frac{20,551}{47,040}v^8 + \dots, \\
 AF &= 0.
 \end{aligned} \tag{45}$$

$K_0(vs.)$ may be expressed as a Taylor series expansion:

$$K_0(vs.) = \frac{55}{24} + \frac{251v^4}{720} + \frac{647v^6}{756} + \dots \tag{46}$$

From a computational standpoint, we shall refer to the aforementioned new technique as **Algorithm IV**.

8. Phase-Fitted and Amplification-Fitted Fourth Order Adams–Bashforth Method

Procedure (27) is taken into account, with $K_1(vs.) = -\frac{59}{24}$, and $K_3(vs.) = -\frac{9}{24}$

The following is the result that we obtain when we use the straightforward approach for calculating the phase-lag and the amplification factor:

$$PhErr = \frac{\cos(4v) - \cos(3v) + K_0(vs.)v\sin(3v) - \frac{59v\sin(2v)}{24} + K_2(vs.)v\sin(vs.)}{\frac{143}{12} - 3K_0(vs.) - K_2(vs.)}, \tag{47}$$

$$AF = \frac{\sin(4v) - \sin(3v) - K_0(vs.)v\cos(3v) + \frac{59v\cos(2v)}{24} - K_2(vs.)v\cos(vs.) + \frac{3}{8}v}{-v^2K_2(vs.) - 9v^2K_0(vs.) + \frac{59}{6}v^2 - 1}, \tag{48}$$

with $PhErr$ denoting the phase-lag and AF denoting the amplification factor.

After the phase-lag and amplification factors have been eliminated, or $PhErr = 0$ and $AF = 0$, the following result is obtained:

$$K_0(vs.) = \frac{1}{24} \frac{48(\sin(vs.))^2 \cos(vs.) + 25v\sin(vs.) - 24(\sin(vs.))^2 - 12\cos(vs.) + 12}{v\cos(vs.)\sin(vs.)}, \tag{49}$$

$$K_2(vs.) = -\frac{1}{24} \frac{18v(\sin(vs.))^3 - 43v\sin(vs.) - 12\cos(vs.) + 12}{v\cos(vs.)\sin(vs.)}. \tag{50}$$

By expanding the aforementioned formulas using the Taylor series, we obtain:

$$\begin{aligned}
 K_0(vs.) &= \frac{55}{24} + \frac{95}{576}v^4 + \frac{2935}{48,384}v^6 + \frac{14,417}{580,608}v^8 \\
 &\quad + \frac{27,559}{2,737,152}v^{10} + \frac{121,989,367}{29,889,699,840}v^{12} + \dots,
 \end{aligned} \tag{51}$$

$$\begin{aligned}
 K_2(vs.) &= \frac{37}{24} + \frac{529}{2880}v^4 + \frac{14,621}{241,920}v^6 + \frac{10,321}{414,720}v^8 \\
 &\quad + \frac{4,824,823}{479,001,600}v^{10} + \frac{610,012,553}{149,448,499,200}v^{12} + \dots
 \end{aligned} \tag{52}$$

The characteristics of this new method are:

$$\begin{aligned}
 K_0(\text{vs.}) &\text{ see (49),} \\
 K_1(\text{vs.}) &= -\frac{59}{24}, \\
 K_2(\text{vs.}) &\text{ see (50),} \\
 K_3(\text{vs.}) &= -\frac{9}{24}, \\
 LTE &= \frac{251}{720} h^5 \left(s^{\{(5)\}}(t) - \omega^4 s'(t) \right) + O(h^6), \\
 PhErr &= 0, \\
 AF &= 0.
 \end{aligned} \tag{53}$$

From a computational standpoint, we shall refer to the aforementioned new technique as **Algorithm V**.

9. Implicit Method: Adams–Moulton Five-Step Method

In particular, we shall illustrate the famous Adams–Moulton approach of fifth algebraic order, which is the following:

$$s_{n+1} - s_n = \frac{h}{720} \left(251 s'_{n+1} + 646 s'_n - 264 s'_{n-1} + 106 s'_{n-2} - 19 s'_{n-3} \right). \tag{54}$$

together with the local truncation error (*LTE*) provided by:

$$LTE = -\frac{3}{360} h^6 s^{\{(6)\}}(t) + O(h^7). \tag{55}$$

We use the theory from Section 2 to obtain the method's phase-lag and amplification error.

Difference Equation (7) with $k = 4$ is obtained by applying the algorithm (54) to the test Equation (2) with:

$$A_4(\text{vs.}) = \frac{251}{720}, A_3(\text{vs.}) = \frac{323}{360}, A_2(\text{vs.}) = -\frac{11}{30}, A_1(\text{vs.}) = \frac{53}{360}, A_0(\text{vs.}) = -\frac{19}{720}. \tag{56}$$

By applying the Taylor series expansion to the above Equation (19) and setting $m = 0(1)4$, we can obtain:

$$\begin{aligned}
 \frac{\cos(4 \text{ vs.}) - \cos(3 \text{ vs.}) + v \sum_{j=0}^4 A_{k-j}(v) \sin[(k-j) \text{ vs.}]}{2k-1 - \sum_{j=0}^4 A_{k-j}(v)(k-j)} &= \\
 \frac{3}{560} v^6 - \frac{25}{1728} v^8 + \dots
 \end{aligned} \tag{57}$$

Consequently, $q = 4$ and $c = \frac{3}{560}$. The fifth algebraic order Adams–Moulton method is of fourth order phase-lag.

By applying the Taylor series expansion to the above Equation (21) and setting $m = 1(1)4$, we can obtain the following:

$$\begin{aligned}
 \frac{\sin(4 \text{ vs.}) - \sin(3 \text{ vs.}) - v \sum_{j=0}^4 A_{k-j}(v) \cos[(k-j) \text{ vs.}]}{-1 - v^2 \sum_{j=0}^4 A_{k-j}(v)(k-j)^2} &= \\
 -\frac{641}{15,120} v^7 + \frac{56,953}{100,800} v^9 + \dots
 \end{aligned} \tag{58}$$

Consequently, $q = 6$ and $c = -\frac{641}{15,120}$. The Adams–Moulton approach, which is of fifth order algebraic, has an amplification error of sixth order. We will refer to the fifth algebraic order Adams–Moulton algorithm as **Algorithm VI** for our computing purposes.

9.1. Minimal Phase-Lag

We examine the following basic five-step algorithm to learn more about the methodologies to minimize the phase-lag:

$$s_{n+1} - s_n = h \left(Q_0(v.s.), s'_{n+1} + Q_1(v.s.), s'_n + Q_2(v.s.) s'_{n-1} + Q_3(v.s.) s'_{n-2} + Q_4(v.s.) s'_{n-3} \right). \quad (59)$$

Procedure to Minimize the Phase-Lag

Below is the procedure that minimizes the phase-lag:

- Eliminating the amplification factor;
- Phase-lag calculation using the coefficient acquired in the preceding stage;
- Expanding the phase-lag calculated before using a Taylor series;
- Defining the set of equations that minimizes the phase-lag;
- Calculation of the updated coefficients.

The following three phase-lag-minimizing algorithms are derived from the aforementioned procedure.

10. Amplification-Fitted Adams–Moulton Five-Step Method of Fifth Algebraic Order with Phase-Lag of Order Four

Let us consider method (59) with $Q_1(v.s.) = \frac{323}{360}$, $Q_2(v.s.) = -\frac{11}{30}$. For the development of this algorithm, see Appendix B.

This novel algorithm has the following features:

$$\begin{aligned} Q_0(v.s.) &= \frac{\Psi_{21}(v.s.)}{360 v \cos(4v)}, \\ Q_1(v.s.) &= \frac{323}{360}, \\ Q_2(v.s.) &= -\frac{11}{30}, \\ Q_3(v.s.) &= \frac{53}{360}, \\ Q_4(v.s.) &= -\frac{19}{720}, \\ LTE &= -\frac{3}{160} h^6 \left(s^{(6)}(t) \right) + O(h^7), \\ PhErr &= \frac{3}{560} v^6 + \dots, \\ AF &= 0, \end{aligned} \quad (60)$$

where

$$\begin{aligned} \Psi_{21}(v.s.) &= -53v \cos(v.s.) - 323v \cos(3v) + 132v \cos(2v) \\ &+ \frac{19}{2}v + 360 \sin(4v) - 360 \sin(3v). \end{aligned} \quad (61)$$

$Q_0(v.s.)$ may be expressed as a Taylor series expansion:

$$Q_0(v.s.) = \frac{251}{720} + \frac{641v^6}{15,120} + \frac{269,443v^8}{907,200} + \frac{77,375,869v^{10}}{39,916,800} + \dots \quad (62)$$

From a computational standpoint, we shall refer to the aforementioned new technique as **Algorithm VII**.

11. Amplification-Fitted Adams–Moulton Five-Step Method of Second Algebraic Order with Phase-Lag of Order Six

Let us consider method (59) with $Q_2(vs.) = \frac{323}{360}$.

For the development of this algorithm, see Appendix C.

This novel algorithm has the following features:

$$\begin{aligned}
Q_0(vs.) &= \frac{\Psi_{34}(vs.)}{20,160 v \cos(4v)}, \\
Q_1(vs.) &= \frac{323}{360}, \\
Q_2(vs.) &= -\frac{167}{480}, \\
Q_3(vs.) &= \frac{317}{2520}, \\
Q_4(vs.) &= -\frac{397}{20,160}, \\
LTE &= -\frac{3}{560} h^3 \left(s^{\{(3)\}}(t) + \omega^2 s'(t) \right) + +O(h^4), \\
PhErr &= -\frac{313}{60,480} v^8 + \dots, \\
AF &= 0,
\end{aligned} \tag{63}$$

where

$$\begin{aligned}
\Psi_{34}(vs.) &= 7014 v \cos(2v) - 2536 v \cos(vs.) - 18,088 v \cos(3v) \\
&\quad + 397 v + 20,160 \sin(4v) - 20,160 \sin(3v).
\end{aligned} \tag{64}$$

$Q_0(vs.)$ may be expressed as a Taylor series expansion:

$$Q_0(vs.) = \frac{6947}{20,160} - \frac{3}{560} v^2 - \frac{13}{1120} v^4 - \frac{4397}{302,400} v^6 - \frac{2,983,453}{50,803,200} v^8 + \dots \tag{65}$$

From a computational standpoint, we shall refer to the aforementioned new technique as **Algorithm VIII**.

12. Amplification-Fitted Adams–Moulton Five-Step Method of Second Algebraic Order with Phase-Lag of Order Eight

Let us consider method (59).

For the development of this algorithm, see Appendix D.

This novel algorithm has the following features:

$$\begin{aligned}
Q_0(vs.) &= \frac{1}{120,960} \frac{\Psi_{59}(vs.)}{v \cos(4v)}, \\
Q_1(vs.) &= \frac{5561}{8640}, \\
Q_2(vs.) &= -\frac{163}{1728}, \\
Q_3(vs.) &= \frac{23}{1344}, \\
Q_4(vs.) &= -\frac{191}{120,960}, \\
LTE &= -\frac{1867}{60,480} h^3 \left(s^{\{(3)\}}(t) + \omega^2 s'(t) \right) + +O(h^4), \\
PhErr &= -\frac{2497}{25,401,600} v^{10} + \dots, \\
AF &= 0,
\end{aligned} \tag{66}$$

where

$$\begin{aligned}\Psi_{59}(vs.) &= -77,854 v \cos(3v) + 11,410 v \cos(2v) + 191 v \\ &\quad - 2070 v \cos(vs.) + 120,960 \sin(4v) - 120,960 \sin(3v).\end{aligned}\quad (67)$$

$Q_0(vs.)$ may be expressed as a Taylor series expansion:

$$Q_0(vs.) = \frac{52,637}{120,960} + \frac{1867}{60,480} v^2 + \frac{2531}{725,760} v^4 + \frac{14,257}{21,772,800} v^6 + \frac{448,163}{1,219,276,800} v^8 + \dots\quad (68)$$

From a computational standpoint, we shall refer to the aforementioned new technique as **Algorithm IX**.

13. Amplification-Fitted Adams–Moulton Five-Step Method of Fifth Algebraic Order

Let us consider method (59) with $Q_1(vs.) = \frac{323}{360}$, $Q_2(vs.) = -\frac{11}{30}$, $Q_3(vs.) = \frac{53}{360}$, $Q_4(vs.) = -\frac{19}{720}$.

13.1. Eliminating the Amplification Factor

We obtain the following result when we use the straightforward approach for computing the amplification factor (21):

$$AF = \frac{\Psi_{60}(vs.)}{-16 v^2 Q_0(vs.) - \frac{304 v^2}{45} - 1}. \quad (69)$$

where AF denotes the amplification factor, and

$$\begin{aligned}\Psi_{60}(vs.) &= \sin(4v) - \sin(3v) - Q_0(vs.) v \cos(4v) \\ &\quad - \frac{323}{360} v \cos(3v) + \frac{11}{30} v \cos(2v) \\ &\quad - \frac{53}{360} v \cos(vs.) + \frac{19v}{720}.\end{aligned}\quad (70)$$

Assuming that the amplification factor must be eliminated, or that $AF = 0$, we derive:

$$Q_0(vs.) = \frac{\Psi_{61}(vs.)}{720 v \cos(4v)}. \quad (71)$$

where

$$\begin{aligned}\Psi_{61}(vs.) &= -646 v \cos(3v) + 264 v \cos(2v) - 106 v \cos(vs.) \\ &\quad + 720 \sin(4v) - 720 \sin(3v) + 19v.\end{aligned}\quad (72)$$

Phase-Lag of the Method

We can obtain the phase-lag by plugging the value of $Q_0(vs.)$ that was previously provided into the direct formula for calculating it (19):

$$PhErr = \frac{1}{2} \frac{v \Psi_{62}(vs.)}{\Psi_{63}(vs.)}. \quad (73)$$

where

$$\begin{aligned}\Psi_{62}(vs.) &= 38 (\cos(vs.))^3 \sin(vs.) vs. - 106 (\cos(vs.))^2 \sin(vs.) v \\ &\quad + 113 \cos(vs.) \sin(vs.) v - 135 v \sin(vs.) \\ &\quad - 180 \cos(vs.) + 180, \\ \Psi_{63}(vs.) &= 3524 v (\cos(vs.))^4 - 2880 (\cos(vs.))^3 \sin(vs.) \\ &\quad + 1292 v (\cos(vs.))^3 + 1440 (\cos(vs.))^2 \sin(vs.) \\ &\quad - 3788 v (\cos(vs.))^2 + 1440 \cos(vs.) \sin(vs.) \\ &\quad - 916 v \cos(vs.) - 360 \sin(vs.) + 563 v.\end{aligned}\quad (74)$$

and $PhErr$ denotes the phase-lag.

By applying the Taylor series expansion to Formula (42), we are able to retrieve:

$$PhErr = \frac{3}{560} v^6 + \frac{14,387}{423,360} v^8 + \dots \quad (75)$$

This novel algorithm has the following features:

$$\begin{aligned} Q_0(vs.) &= \frac{\Psi_{61}(vs.)}{720 v \cos(4v)}, \\ Q_1(vs.) &= \frac{323}{360}, \\ Q_2(vs.) &= -\frac{11}{30}, \\ Q_3(vs.) &= \frac{53}{360}, \\ Q_4(vs.) &= -\frac{19}{720}, \\ LTE &= -\frac{3}{160} h^6 s^{\{(6)\}}(t) + O(h^7), \\ PhErr &= \frac{3}{560} v^6 + \frac{14,387}{423,360} v^8 + \dots, \\ AF &= 0. \end{aligned} \quad (76)$$

$Q_0(vs.)$ may be expressed as a Taylor series expansion:

$$Q_0(vs.) = \frac{251}{720} + \frac{641}{15,120} v^6 + \frac{269,443}{907,200} v^8 + \frac{77,375,869}{39,916,800} v^{10} + \dots \quad (77)$$

From a computational standpoint, we shall refer to the aforementioned new technique as **Algorithm X**.

14. Phase-Fitted and Amplification-Fitted Fifth Order Adams–Moulton Method

Let us considering the method (59) with $Q_1(vs.) = \frac{323}{360}$, $Q_2(vs.) = -\frac{11}{30}$, and $Q_4(vs.) = -\frac{19}{720}$.

The following is the result that we obtain when we use the straightforward approach for calculating the phase-lag and the amplification factor:

$$\begin{aligned} PhErr &= \frac{\Psi_{64}(vs.)}{-1815 + 1440 Q_0(vs.) + 360 Q_3(vs.)}, \\ AF &= \frac{\Psi_{65}(vs.)}{11,520 v^2 Q_0(vs.) + 720 v^2 Q_3(vs.) + 4758 v^2 + 720} \end{aligned} \quad (78)$$

with $PhErr$ denoting the phase-lag and AF denoting the amplification factor, and

$$\begin{aligned} \Psi_{64}(vs.) &= -360 Q_0(vs.) v \sin(4v) - 323 v \sin(3v) \\ &- 360 \cos(4v) + 132 v \sin(2v) \\ &- 360 Q_3(vs.) v \sin(vs.) + 360 \cos(3v), \\ \Psi_{65}(vs.) &= 5760 v Q_0(vs.) (\cos(vs.))^4 - 5760 \sin(vs.) (\cos(vs.))^3 \\ &+ 2584 v (\cos(vs.))^3 - 5760 (\cos(vs.))^2 v Q_0(vs.) \\ &+ 2880 \sin(vs.) (\cos(vs.))^2 - 528 (\cos(vs.))^2 v \\ &+ 720 Q_3(vs.) v \cos(vs.) + 2880 \sin(vs.) \cos(vs.) \\ &- 1938 v \cos(vs.) + 720 v Q_0(vs.) \\ &- 720 \sin(vs.) + 245 v. \end{aligned} \quad (79)$$

After the phase-lag and amplification factors have been eliminated, or $PhErr = 0$ and $AF = 0$, the following result is obtained:

$$Q_0(vs.) = \frac{1}{720} \frac{\Psi_{66}(vs.)}{v \left(4(\cos(vs.))^3 + 4(\cos(vs.))^2 - \cos(vs.) - 1 \right)}. \quad (80)$$

$$Q_3(vs.) = \frac{1}{360} \frac{\Psi_{67}(vs.)}{v \left(4(\cos(vs.))^3 + 4(\cos(vs.))^2 - \cos(vs.) - 1 \right)}. \quad (81)$$

where

$$\begin{aligned} \Psi_{66}(vs.) &= 2880 \sin(vs.) (\cos(vs.))^2 - 1292 (\cos(vs.))^2 v \\ &+ 1440 \sin(vs.) \cos(vs.) - 1047 v \cos(vs.) \\ &- 720 \sin(vs.) + 245 v, \\ \Psi_{67}(vs.) &= 76 v (\cos(vs.))^4 + 76 v (\cos(vs.))^3 \\ &+ 226 (\cos(vs.))^2 v - 97 v \cos(vs.) \\ &+ 360 \sin(vs.) - 323 v. \end{aligned} \quad (82)$$

By expanding the aforementioned formulas using the Taylor series, we obtain:

$$\begin{aligned} Q_0(vs.) &= \frac{251}{720} - \frac{1}{160} v^4 - \frac{1271}{362,880} v^6 - \frac{52,901}{21,772,800} v^8 \\ &- \frac{362,891}{179,625,600} v^{10} - \frac{2,012,951,791}{1,120,863,744,000} v^{12} + \dots \end{aligned} \quad (83)$$

$$\begin{aligned} Q_3(vs.) &= \frac{53}{360} + \frac{1}{160} v^4 - \frac{71}{72,576} v^6 - \frac{39,667}{21,772,800} v^8 \\ &- \frac{1,351,639}{718,502,400} v^{10} - \frac{1,977,358,457}{1,120,863,744,000} v^{12} + \dots \end{aligned} \quad (84)$$

The characteristics of this new method are:

$$\begin{aligned} Q_0(vs.) &\text{ see (80),} \\ Q_1(vs.) &= \frac{323}{360}, \\ Q_2(vs.) &= -\frac{11}{30}, \\ Q_3(vs.) &\text{ see (81),} \\ Q_4(vs.) &= -\frac{19}{720}, \\ LTE &= -\frac{3}{160} h^6 \left(s^{\{(6)\}}(t) - \omega^4 s^{\{(2)\}}(t) \right) + O(h^7), \\ PhErr &= 0, \\ AF &= 0. \end{aligned} \quad (85)$$

From a computational standpoint, we shall refer to the aforementioned new technique as **Algorithm XI**.

15. Stability Analysis

In this section, we will study the stability of the methods developed in Sections 4–14.

15.1. Adams–Bashforth Algorithm

A general description of the five-step methods proposed by Adams–Bashforth (Explicit) and Adams–Moulton (Implicit) is as follows:

$$s_{n+1} - s_n = h \left(T_3(v.s.), s'_n + T_2(v.s.) s'_{n-1} + T_1(v.s.) s'_{n-2} + T_0(v.s.) s'_{n-3} \right). \quad (86)$$

Algorithms (22), (34), (37), (45), and (53) that were studied in Sections 4–8 constitute the general algorithm (86).

By combining the scalar test equation:

$$s' = \lambda s \text{ where } \lambda \in \mathcal{C}, \quad (87)$$

with the scheme (86), we can obtain the subsequent difference equation

$$s_{n+1} - S_3(H) s_n - S_2(H) s_{n-1} - S_1(H) s_{n-2} - S_0(H) s_{n-3} = 0, \quad (88)$$

with $H = \lambda h$ and

$$S_3(H) = 1 + T_3 H, \quad S_2(H) = T_2 H, \quad S_1(H) = T_1 H, \quad S_0(H) = T_0 H. \quad (89)$$

Presenting the characteristic equation of (88), we have:

$$r^4 - S_3(H) r^3 - S_2(H) r^2 - S_1(H) r - S_0(H) = 0. \quad (90)$$

15.2. Adams–Moulton Five-Step Algorithm

A general description of the five-step methods proposed by Adams–Moulton (Implicit) is as follows:

$$s_{n+1} - s_n = h \left(G_4(v.s.), s'_{n+1} + G_3(v.s.) s'_n + G_2(v.s.) s'_{n-1} + G_1(v.s.) s'_{n-2} + G_0(v.s.) s'_{n-3} \right). \quad (91)$$

Algorithms (54), (60), (66), (76), (85), and (A31) that were studied in Sections 4–8 constitute the general algorithm (91).

By combining scalar test Equation (87) with scheme (91), we can obtain the subsequent difference equation

$$W_4(H) s_{n+1} - W_3(H) s_n - W_2(H) s_{n-1} - W_1(H) s_{n-2} - W_0(H) s_{n-3} = 0, \quad (92)$$

with $H = \lambda h$ and

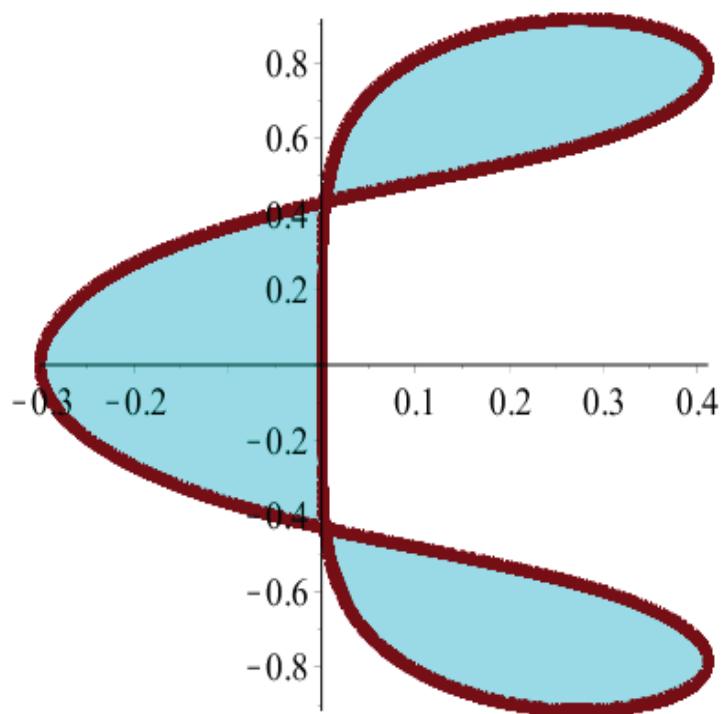
$$\begin{aligned} W_4(H) &= 1 - G_4 H, \quad W_3(H) = 1 + G_3 H, \\ W_2(H) &= G_2 H, \quad W_1(H) = G_1 H, \\ W_0(H) &= G_0 H. \end{aligned} \quad (93)$$

Presenting the characteristic equation of (92), we have:

$$W_4(H) r^4 - W_3(H) r^3 - W_2(H) r^2 - W_1(H) r - W_0(H) = 0. \quad (94)$$

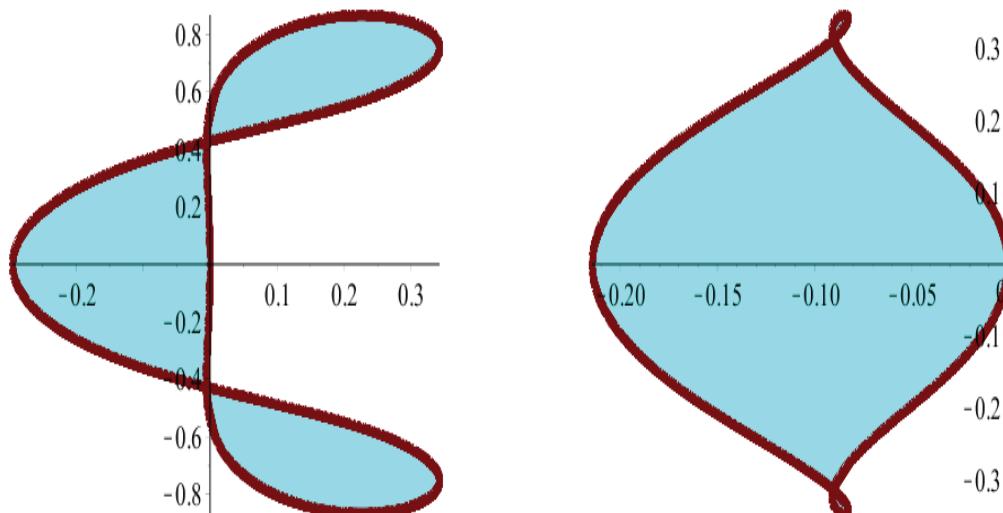
15.3. Stabilities of Adams–Bashforth and Adams–Moulton Algorithms

We can visualize the stability areas for $\theta \in [0, 2\pi]$ by solving the original Equations (90) and (94) in H and inserting $r = \exp(i\theta)$, where $i = \sqrt{-1}$. We display the stability areas for the accomplished **Methods I–V** in Figures 1–11. We show the stability areas for $v = 1$, $v = 30$, and $v = 1000$ for the instances of Methods II–V.



Stability Region of Adams-Bashforth 4th algebraic order method (Algorithm I)

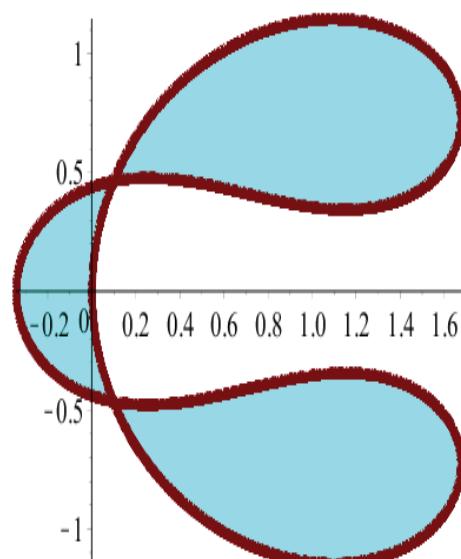
Figure 1. Stability Region for the classical Fourth-Order Adams–Bashforth Method (Algorithm I). The axes of the stability region are H and θ .



Stability Region of the Amplification Fitted Adams-Bashforth Method of Algebraic Order 4 with Phase-Lag of Order 4 (Algorithm II)- Case
 $v=1$

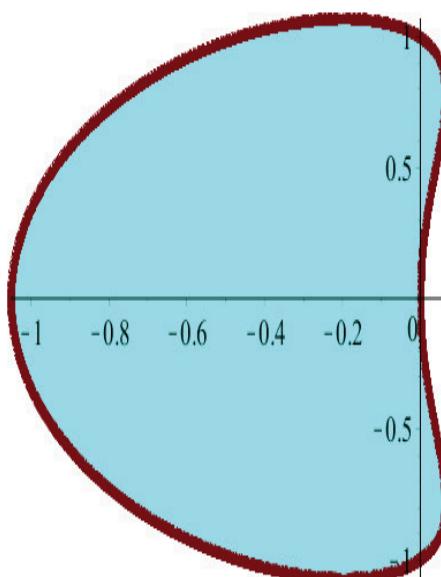
Stability Region of the Amplification Fitted Adams-Bashforth Method of Algebraic Order 4 with Phase-Lag of Order 4 (Algorithm II)- Case
 $v=30$

Figure 2. Cont.

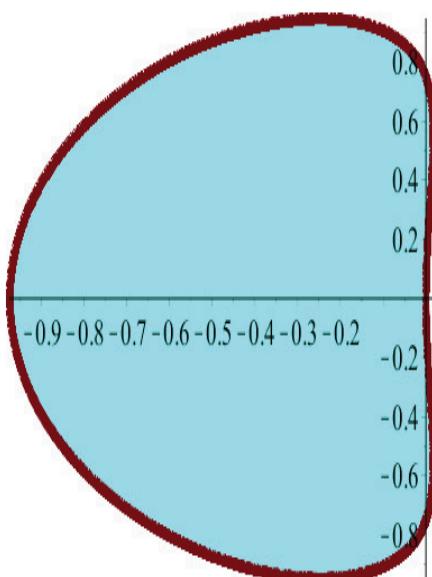


Stability Region of the Amplification Fitted
Adams-Basforth Method of Algebraic Order 4
with Phase-Lag of Order 4 (Algorithm II)- Case
 $v=1000$

Figure 2. Stability Region of the Amplification-Fitted Adams–Bashforth Method of Algebraic Order Four with Phase-Lag of Order Four (Algorithm II). The axes of the stability region are H and θ .

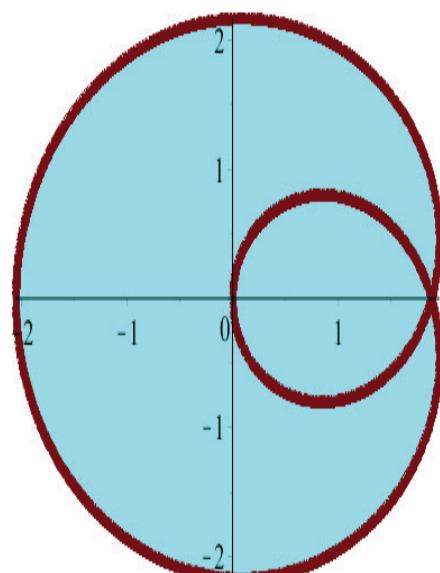


Stability Region of the Amplification Fitted
Adams-Basforth Method of Algebraic Order 3
with Phase-Lag of Order 6 (Algorithm III)- Case
 $v=1$



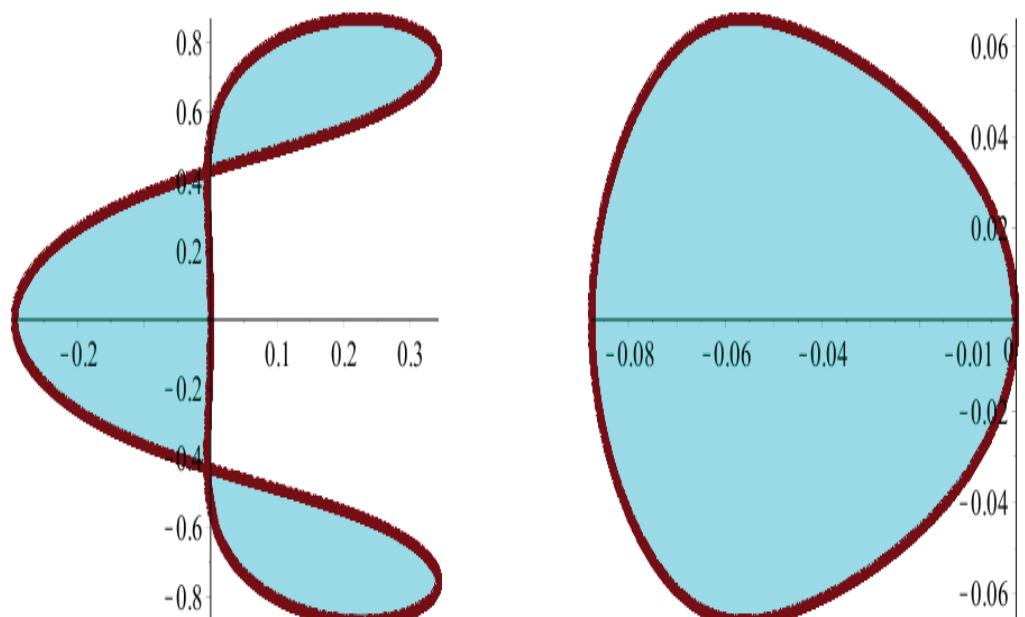
Stability Region of the Amplification Fitted
Adams-Basforth Method of Algebraic Order 3
with Phase-Lag of Order 6 (Algorithm III)- Case
 $v=30$

Figure 3. Cont.



Stability Region of the Amplification Fitted
Adams-Basforth Method of Algebraic Order 3
with Phase-Lag of Order Six (Algorithm III)- Case
 $v=1000$

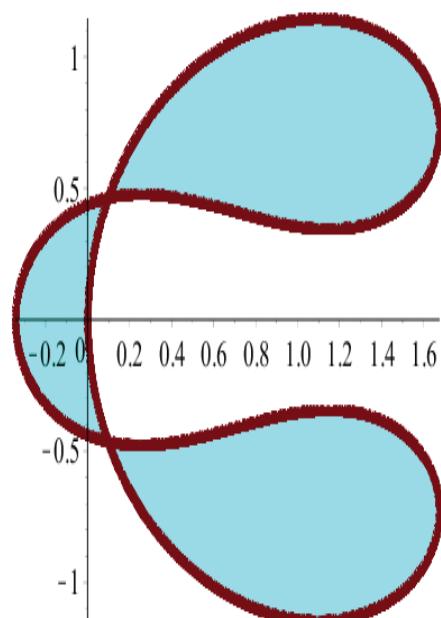
Figure 3. Stability Region of the Amplification-Fitted Adams–Bashforth Method of Algebraic Order Three with Phase-Lag of Order Six (Algorithm III). The axes of the stability region are H and θ .



Stability Region of the Amplification Fitted
Adams-Basforth Method of Algebraic Order 4
(Algorithm IV) - Case=1

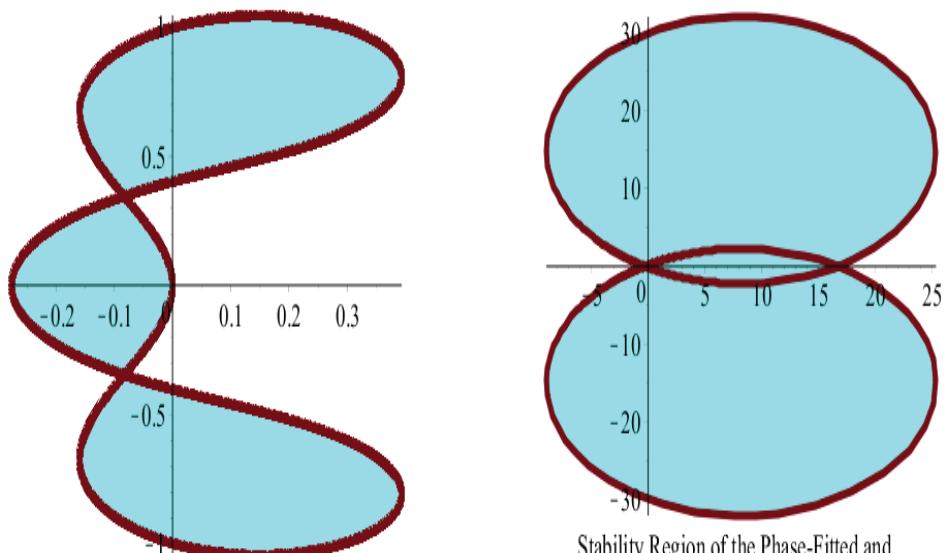
Stability Region of the Amplification Fitted
Adams-Basforth Method of Algebraic Order 4
(Algorithm IV) - Case=10

Figure 4. *Cont.*



Stability Region of the Amplification Fitted Adams-Basforth Method of Algebraic Order 4 (Algorithm IV) - Case=1000

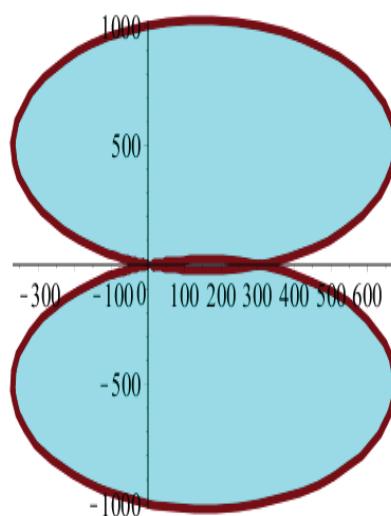
Figure 4. Stability Region of the Amplification-Fitted Adams–Bashforth Method of Algebraic Order Four (Algorithm IV). The axes of the stability region are H and θ .



Stability Region of the Phase-Fitted and Amplification-Fitted Adams-Basforth Method of Algebraic Order 4 (Algorithm V) - Case with $v=1$

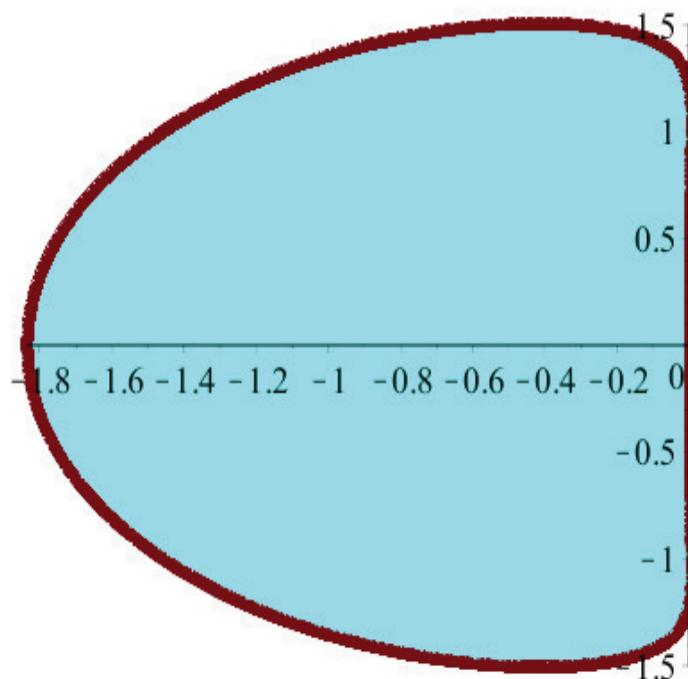
Stability Region of the Phase-Fitted and Amplification-Fitted Adams-Basforth Method of Algebraic Order 4 (Algorithm V) - Case with $v=30$

Figure 5. Cont.



Stability Region of the Phase-Fitted and
Amplification-Fitted Adams-Bashforth Method
of Algebraic Order 4 (Algorithm V) - Case with
 $v=1000$

Figure 5. Stability Region of the Phase-Fitted and Amplification-Fitted Adams–Bashforth Method of Algebraic Order Four (Algorithm V). The axes of the stability region are H and θ .



Stability Region of Adams-Moulton 5th algebraic
order method (Algorithm VI)

Figure 6. Stability Region of Adams–Moulton Fifth Algebraic Order Method (Algorithm VI). The axes of the stability region are H and θ .

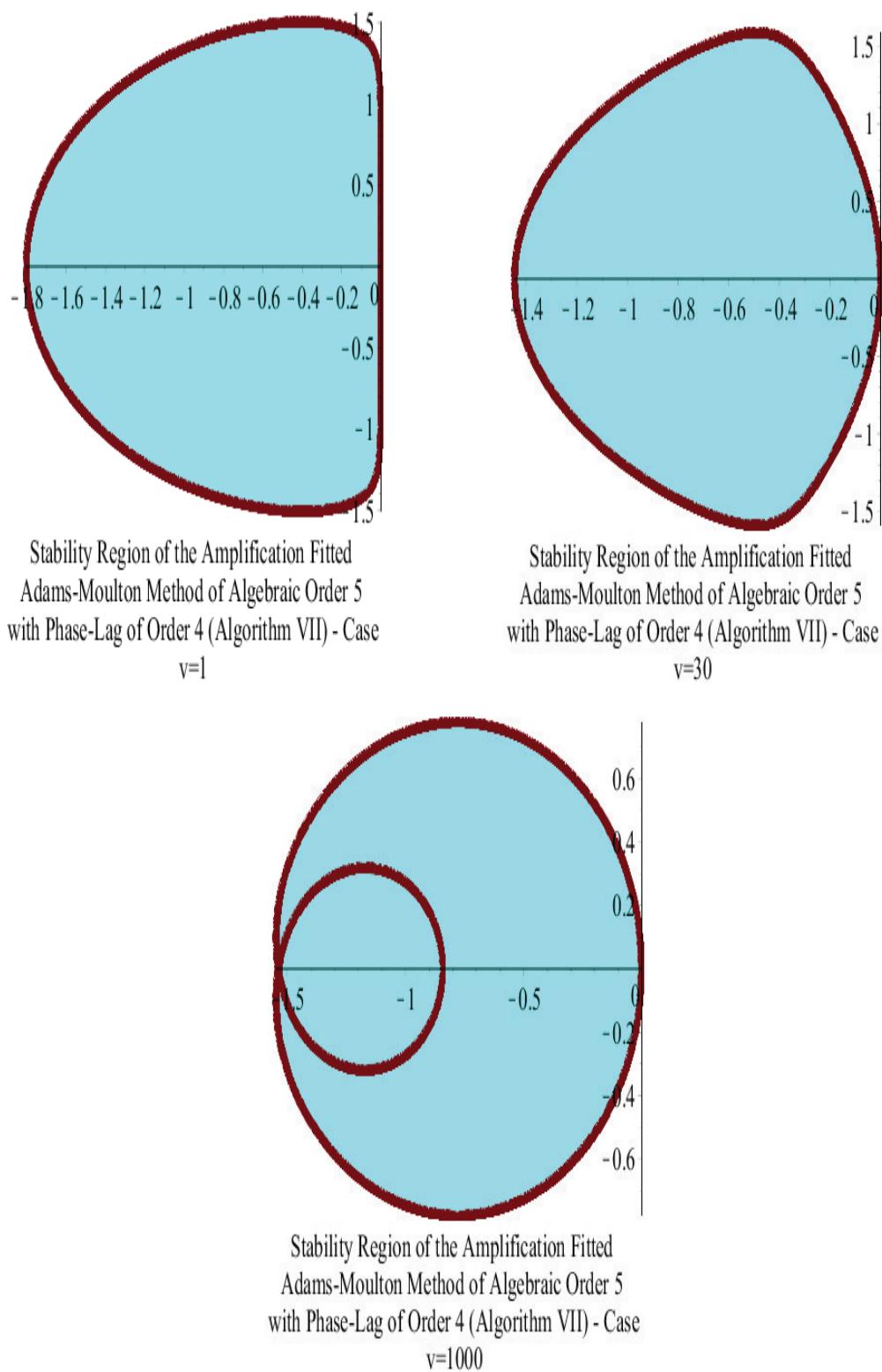


Figure 7. Stability Region of the Amplification-Fitted Adams–Moulton Method of Algebraic Order Five with Phase-Lag of Order Four (Algorithm VII). The axes of the stability region are H and θ .

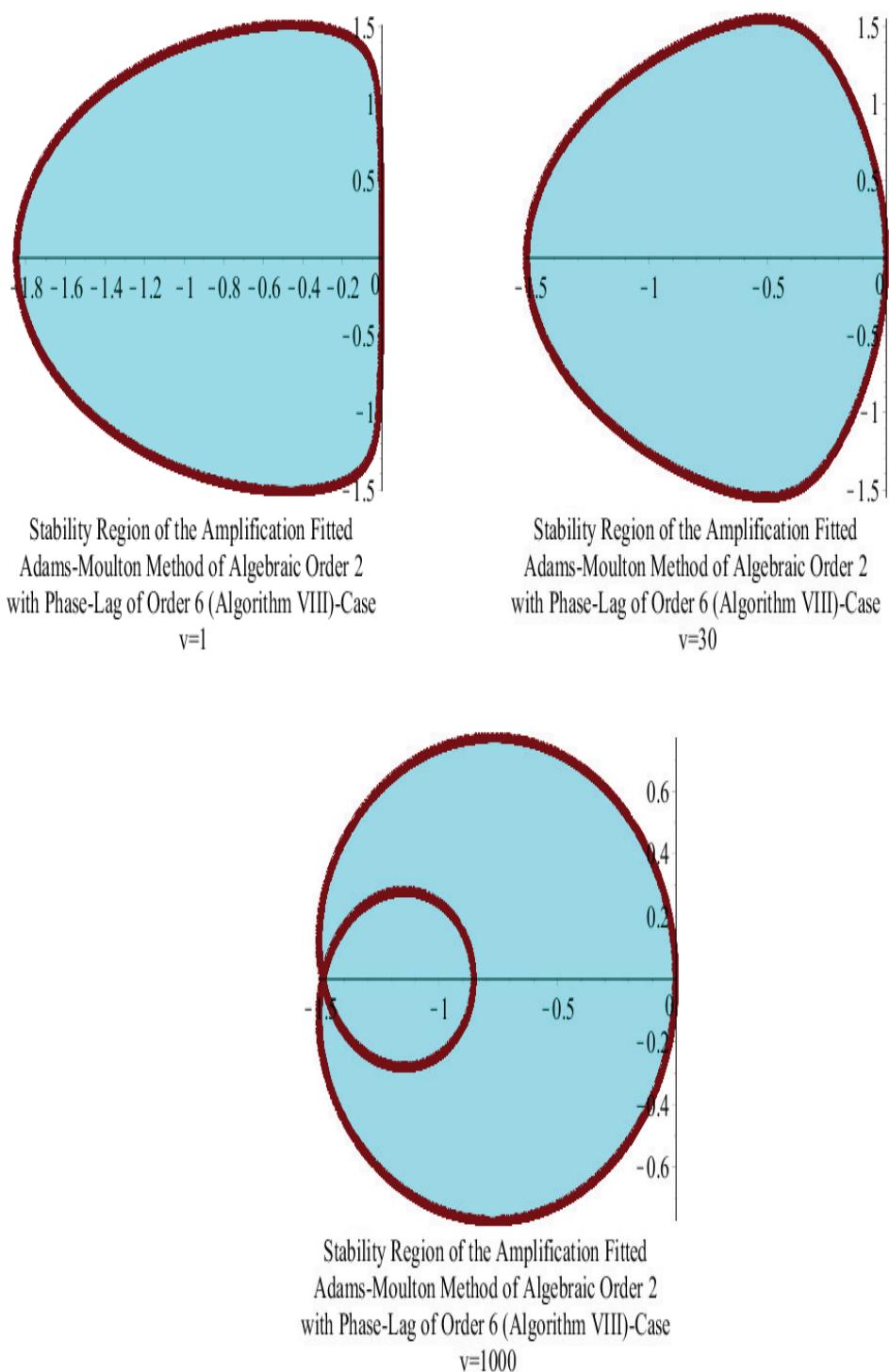


Figure 8. Stability Region of the Amplification-Fitted Adams–Moulton Method of Algebraic Order Two with Phase-Lag of Order Six (Algorithm VIII). The axes of the stability region are H and θ .

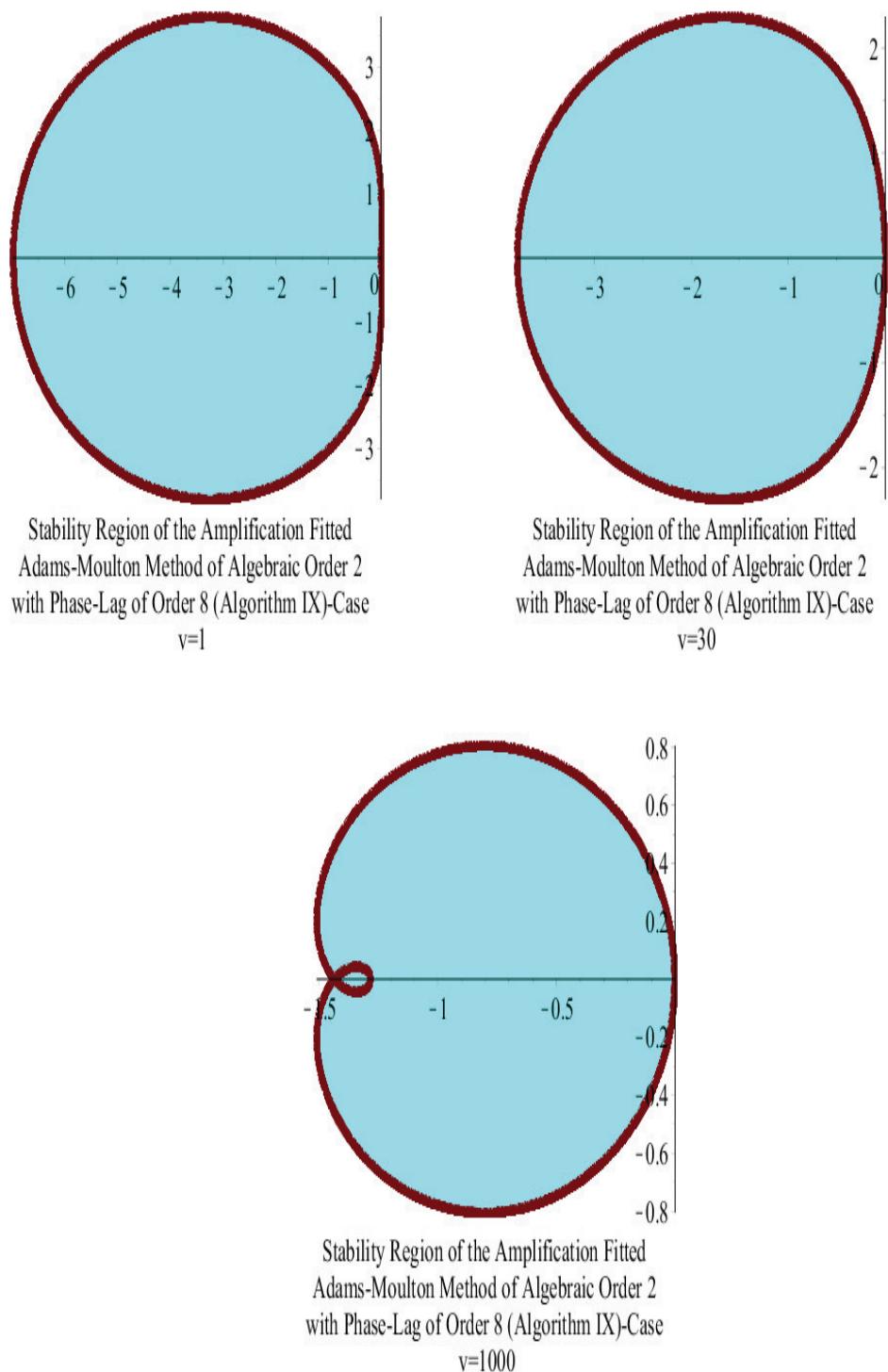


Figure 9. Stability Region of the Amplification-Fitted Adams–Moulton Method of Algebraic Order Two with Phase-Lag of Order Eight (Algorithm IX). The axes of the stability region are H and θ .

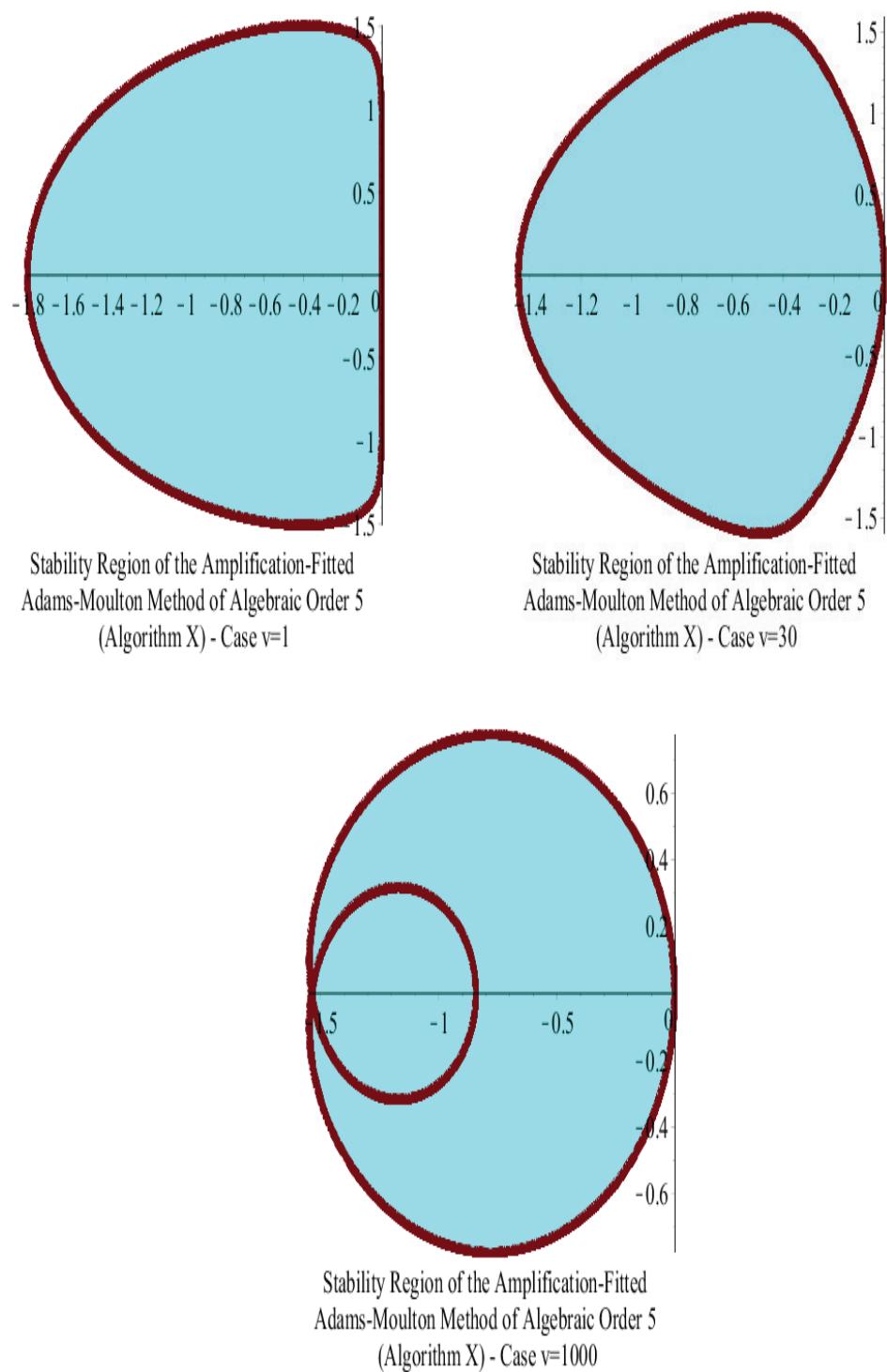


Figure 10. Stability Region of the Amplification-Fitted Adams–Moulton Method of Algebraic Order Five (Algorithm X). The axes of the stability region are H and θ .

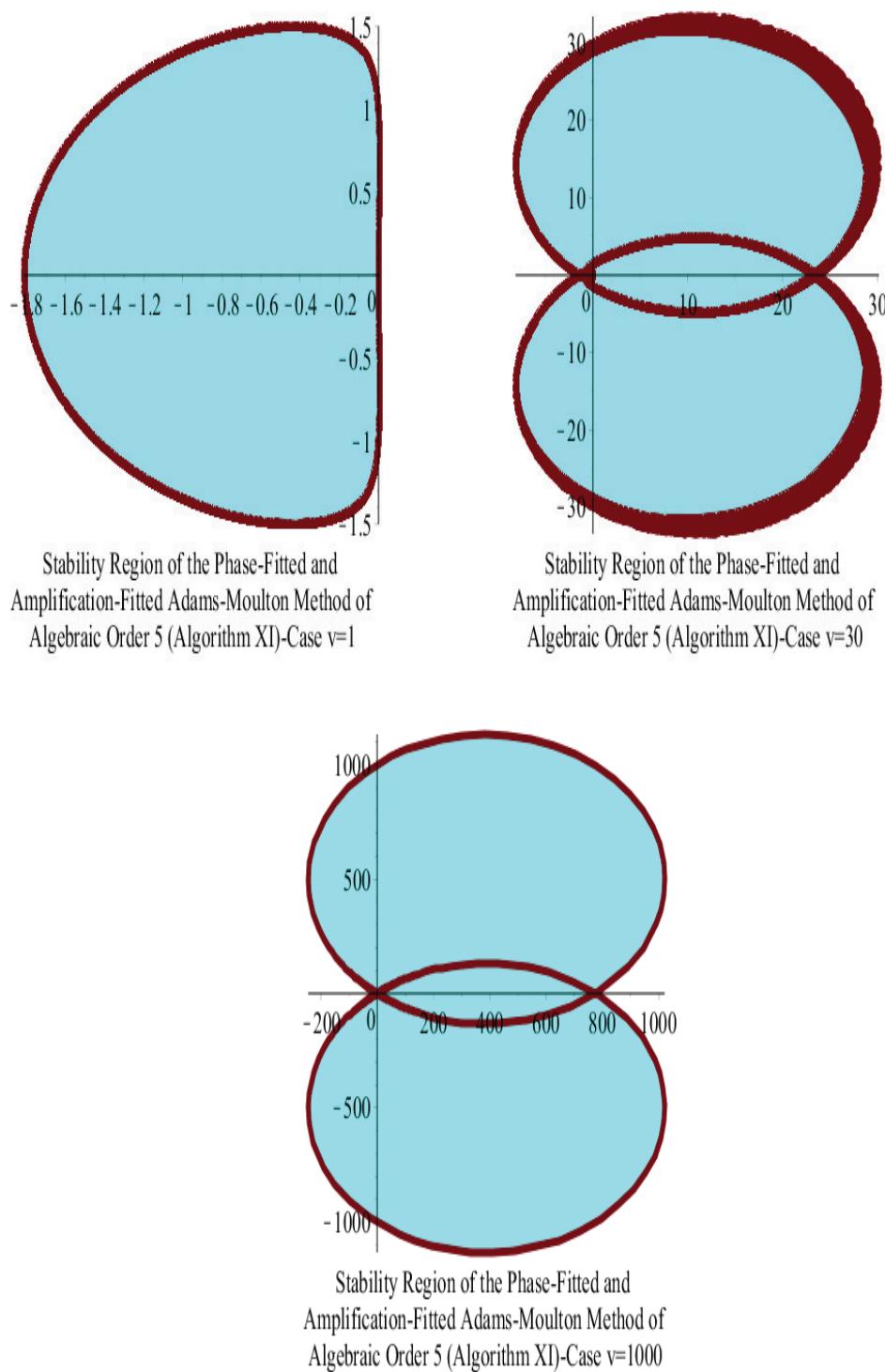


Figure 11. Stability Region of the Phase-Fitted and Amplification-Fitted Adams–Moulton Method of Algebraic Order Five (Algorithm XI). The axes of the stability region are H and θ .

16. Numerical Results

In this section, we will investigate the efficiency of the methods developed in Sections 4–14 comparing them with very well-known methods in the literature. The comparison will take place for well-known problems in the literature.

The newly developed methods are applied in the form of predictor–corrector. More specifically, for each problem, and for the initial four steps, we use a high-order Runge–

Kutta method. Then, we apply the Adams–Bashforth methods developed above as a predictor, and finally, we apply the Adams–Moulton methods developed above as a corrector.

16.1. Problem of Stiefel and Bettis

Stiefel and Bettis [116] investigated the following nearly periodic orbit issue, which we take into consideration.

$$\begin{aligned} s_1''(x) &= -s_1(x) + 0.001 \cos(x), & s_1(0) = 1, & s_1'(0) = 0, \\ s_2''(x) &= -s_2(x) + 0.001 \sin(x), & s_2(0) = 0, & s_2'(0) = 0.9995. \end{aligned} \quad (95)$$

Here is the exact solution:

$$\begin{aligned} s_1(x) &= \cos(x) + 0.0005x \sin(x), \\ s_2(x) &= \sin(x) - 0.0005x \cos(x). \end{aligned} \quad (96)$$

We apply the parameter $\omega = 1$ to this problem.

For values of $0 \leq x \leq 100,000$, the following numerical approaches have been used to solve Equation (95):

- The Classical Adams–Bashforth–Moulton Algorithm of the fifth order (Algorithms (22)–(54)), which is denoted as **Numer. Algor. I**;
- The amplification-fitted Adams–Bashforth–Moulton Algorithm of the fifth order (algorithms (45)–(76)), which is denoted as **Numer. Algor. II**;
- The Runge–Kutta–Dormand–Prince fourth-order method [48], which is denoted as **Numer. Algor. III**;
- The Runge–Kutta–Dormand–Prince fifth-order method [48], which is denoted as **Numer. Algor. IV**;
- The Runge–Kutta–Fehlberg fourth-order method [117], which is denoted as **Numer. Algor. V**;
- The Runge–Kutta–Fehlberg fifth-order method [117], which is denoted as **Numer. Algor. VI**;
- The Runge–Kutta–Cash–Karp fifth-order method [118], which is denoted as **Numer. Algor. VII**;
- The amplification-fitted Adams–Bashforth–Moulton Algorithm of the fifth algebraic order with phase-lag of order four (Algorithms (34)–(60)), which is denoted as **Numer. Algor. VIII**;
- The amplification-fitted Adams–Bashforth–Moulton Algorithm of the second algebraic order with phase-lag of order six (Algorithms (37)–(A31)), which is denoted as **Numer. Algor. IX**;
- The amplification-fitted Adams–Bashforth–Moulton Algorithm of the second algebraic order with phase-lag of order eight (Algorithms (37)–(66)), which is denoted as **Numer. Algor. X**;
- The amplification-fitted and phase-fitted Adams–Bashforth–Moulton Algorithm of the fifth order (Algorithms (53)–(85)), which is denoted as **Numer. Algor. XI**.

We show the greatest absolute error of the solutions obtained by each of the numerical approaches outlined earlier in Figure 6, which pertains to the Stiefel and Bettis problem [116].

The following may be seen in Figure 12:

- Numer. Algor. VII is more efficient than Numer. Algor. IV;
- Numer. Algor. V is more efficient than Numer. Algor. VII;
- Numer. Algor. VI is more efficient than Numer. Algor. V;
- Numer. Algor. III is more efficient than Numer. Algor. VI for the most step sizes but for small step sizes has approximately the same efficiency as Numer. Algor. VI;
- Numer. Algor. I is more efficient than Numer. Algor. VI;
- Numer. Algor. II and Numer. Algor. VIII are more efficient than Numer. Algor. I;

- Numer. Algor. IX has mixed behavior. For big step sizes, it has approximately the same efficiency as Numer. Algor. II and Numer. Algor. VIII. For middle step sizes, it is more efficient than Numer. Algor. III but less efficient than Numer. Algor. I. For small step sizes, it has approximately the same efficiency as Numer. Algor. II and Numer. Algor. VI;
- Numer. Algor. X has mixed behavior. For big step sizes, it is more efficient than Numer. Algor. II. For middle step sizes, it is more efficient than Numer. Algor. III but is less efficient than Numer. Algor. I. For small step sizes, it has approximately the same efficiency as Numer. Algor. III;
- Numer. Algor. XI gives the most efficient results.

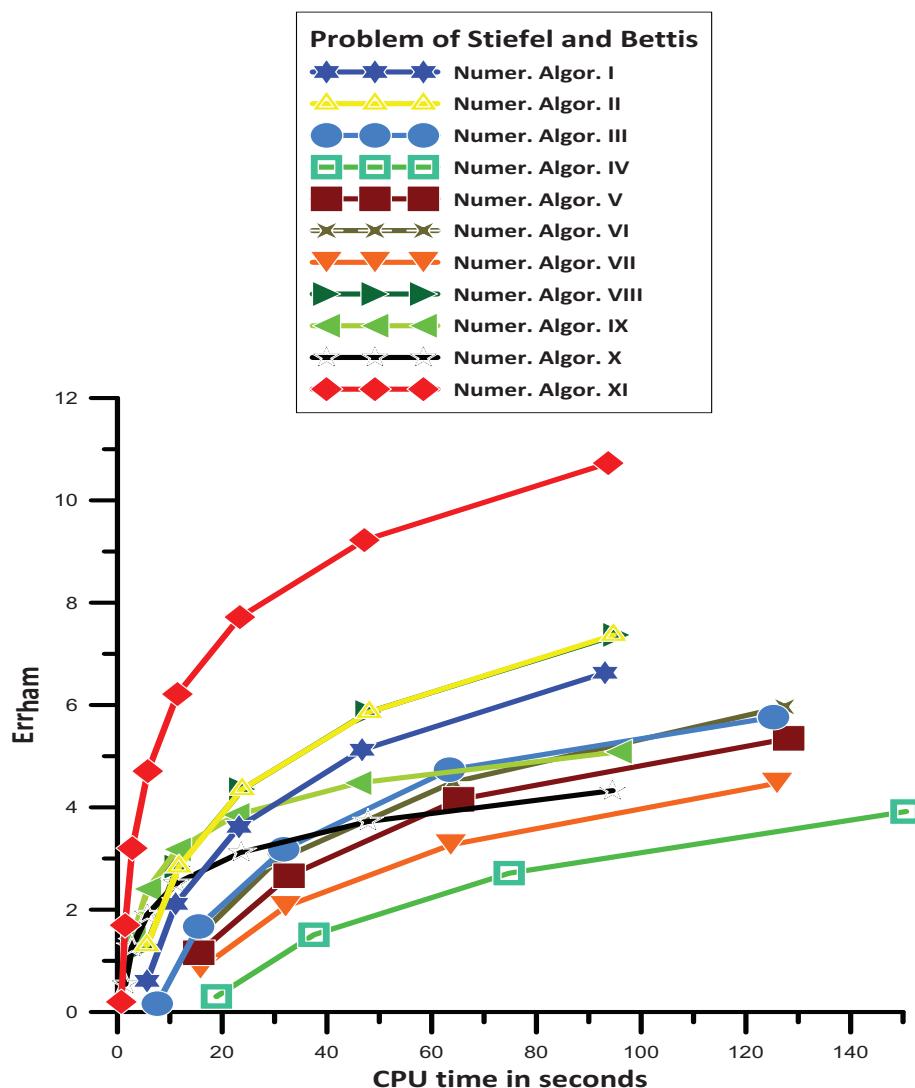


Figure 12. Numerical results for the problem of Stiefel and Bettis [116].

16.2. Problem of Franco et al. [119]

The inhomogeneous linear problem that Franco et al. [119] examined is taken into consideration here:

$$\begin{aligned} s_1''(x) &= -\frac{1}{2} (\mu^2 + 1) s_1(x) - \frac{1}{2} (\mu^2 - 1) s_2(x), \quad s_1(0) = 1, \quad s_1'(0) = 1, \\ s_2''(x) &= -\frac{1}{2} (\mu^2 - 1) s_1(x) - \frac{1}{2} (\mu^2 + 1) s_2(x), \quad s_2(0) = -1, \quad s_2'(0) = -1. \end{aligned} \quad (97)$$

The exact solution is

$$\begin{aligned}s_1(x) &= \cos(x) + \sin(x), \\ s_2(x) &= -\cos(x) - \sin(x).\end{aligned}\quad (98)$$

where $\mu = 10^4$. For this problem, we use $\omega = 1$.

For $0 \leq x \leq 100,000$, the numerical solution of the system of Equation (97) has been found using the techniques outlined in Section 16.1.

The following may be seen in Figure 13:

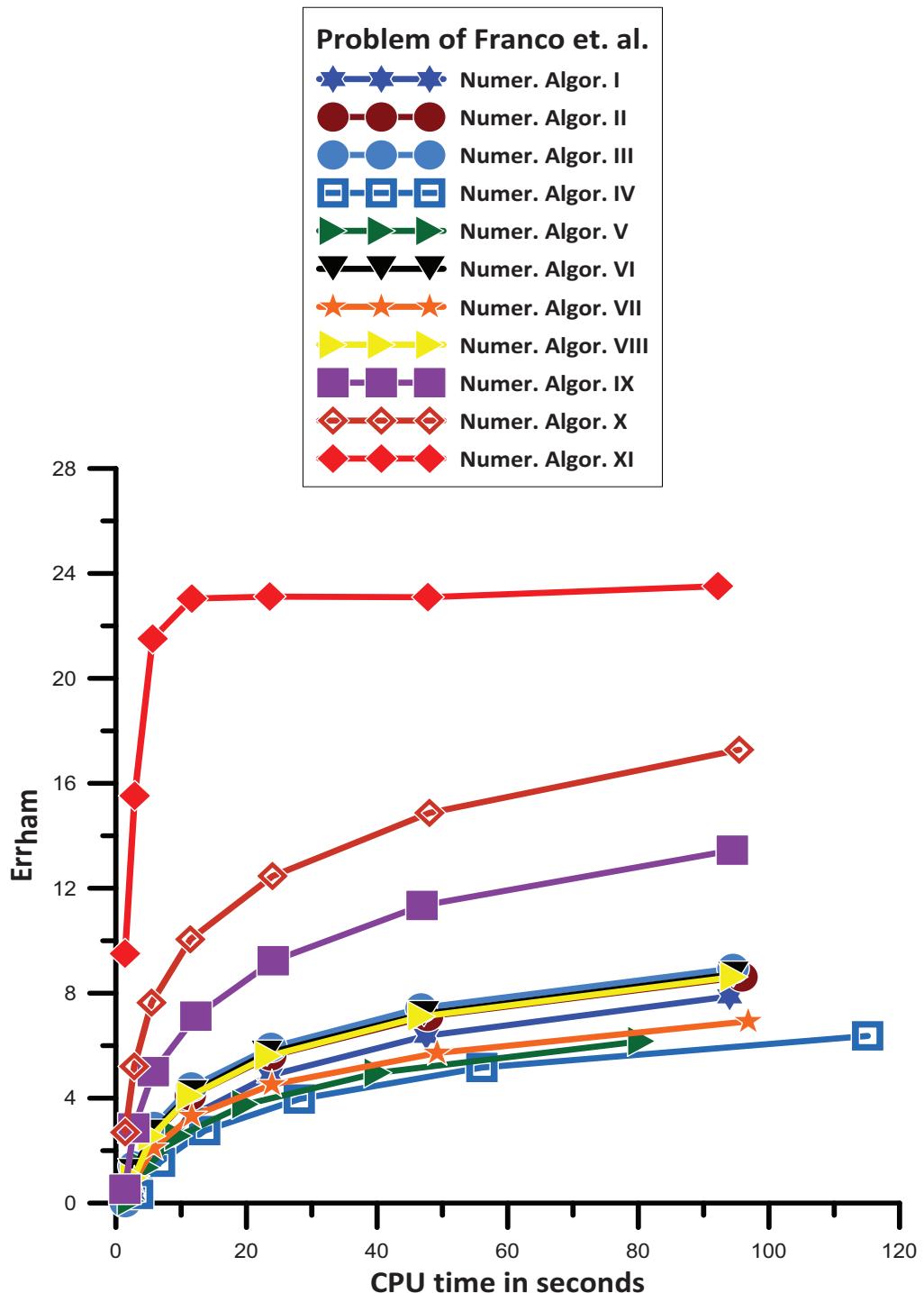


Figure 13. Numerical results for the problem of Franco et al. [119].

- Numer. Algor. V is more efficient than Numer. Algor. IV;
- Numer. Algor. VII is more efficient than Numer. Algor. V;
- Numer. Algor. I is more efficient than Numer. Algor. VII;
- Numer. Algor. VIII is more efficient than Numer. Algor. I;
- Numer. Algor. VIII has approximately the same efficiency as Numer. Algor. VI, Numer. Algor. III, and Numer. Algor. II;
- Numer. Algor. IX is more efficient than Numer. Algor. VIII;
- Numer. Algor. X is more efficient than Numer. Algor. IX;
- Numer. Algor. XI gives the most efficient results.

16.3. Problem of Franco and Palacios [120]

The problem that Franco and Palacios [120] investigated is taken into account here:

$$\begin{aligned} s_1''(x) &= -s_1(x) + \varepsilon \cos(\vartheta x), \quad s_1(0) = 1, \quad s_1'(0) = 0, \\ s_2''(x) &= -s_2(x) + \varepsilon \sin(\vartheta x), \quad s_2(0) = 0, \quad s_2'(0) = 1. \end{aligned} \quad (99)$$

The exact solution is

$$\begin{aligned} s_1(x) &= \frac{1 - \varepsilon - \vartheta^2}{1 - \vartheta^2} \cos(x) + \frac{\varepsilon}{1 - \vartheta^2} \cos(\vartheta x), \\ s_2(x) &= \frac{1 - \varepsilon \vartheta - \vartheta^2}{1 - \vartheta^2} \sin(x) + \frac{\varepsilon}{1 - \vartheta^2} \sin(\vartheta x). \end{aligned} \quad (100)$$

where $\varepsilon = 0.001$ and $\vartheta = 0.01$. For this problem, we use $\omega = \max(1, |\vartheta|)$.

Using the techniques outlined in Section 16.1, the numerical solution to the system of Equation (99) has been obtained for $0 \leq x \leq 100,000$.

The following may be seen in Figure 14:

- Numer. Algor. V is more efficient than Numer. Algor. IV;
- Numer. Algor. VII is more efficient than Numer. Algor. V;
- Numer. Algor. VI is more efficient than Numer. Algor. VII;
- Numer. Algor. VI has approximately the same efficiency as Numer. Algor. III;
- Numer. Algor. I is more efficient than Numer. Algor. VI;
- Numer. Algor. VIII is more efficient than Numer. Algor. II;
- Numer. Algor. IX is more efficient than Numer. Algor. VIII;
- Numer. Algor. X is more efficient than Numer. Algor. IX;
- Numer. Algor. XI gives the most efficient results.

16.4. A Nonlinear Orbital Problem [121]

The nonlinear orbital problem that Simos investigated in [121] is taken into consideration here:

$$\begin{aligned} s_1''(x) &= -s^2 s_1(x) + \frac{2 s_1(x) s_2(x) - \sin(2 s x)}{(s_1(x)^2 + s_2(x)^2)^{\frac{3}{2}}}, \quad s_1(0) = 1, \quad s_1'(0) = 0, \\ s_2''(x) &= -s^2 s_2(x) + \frac{s_1(x)^2 - s_2(x)^2 - \cos(2 s x)}{(s_1(x)^2 + s_2(x)^2)^{\frac{3}{2}}}, \quad s_2(0) = 0, \quad s_2'(0) = s. \end{aligned} \quad (101)$$

The exact solution is

$$s_1(x) = \cos(s x), \quad s_2(x) = \sin(s x), \quad (102)$$

where $s = 10$. For this problem, we use $\omega = 10$.

The numerical solution of the system of Equation (101) has been achieved for $0 \leq x \leq 100,000$ by using the techniques outlined in Section 16.1.

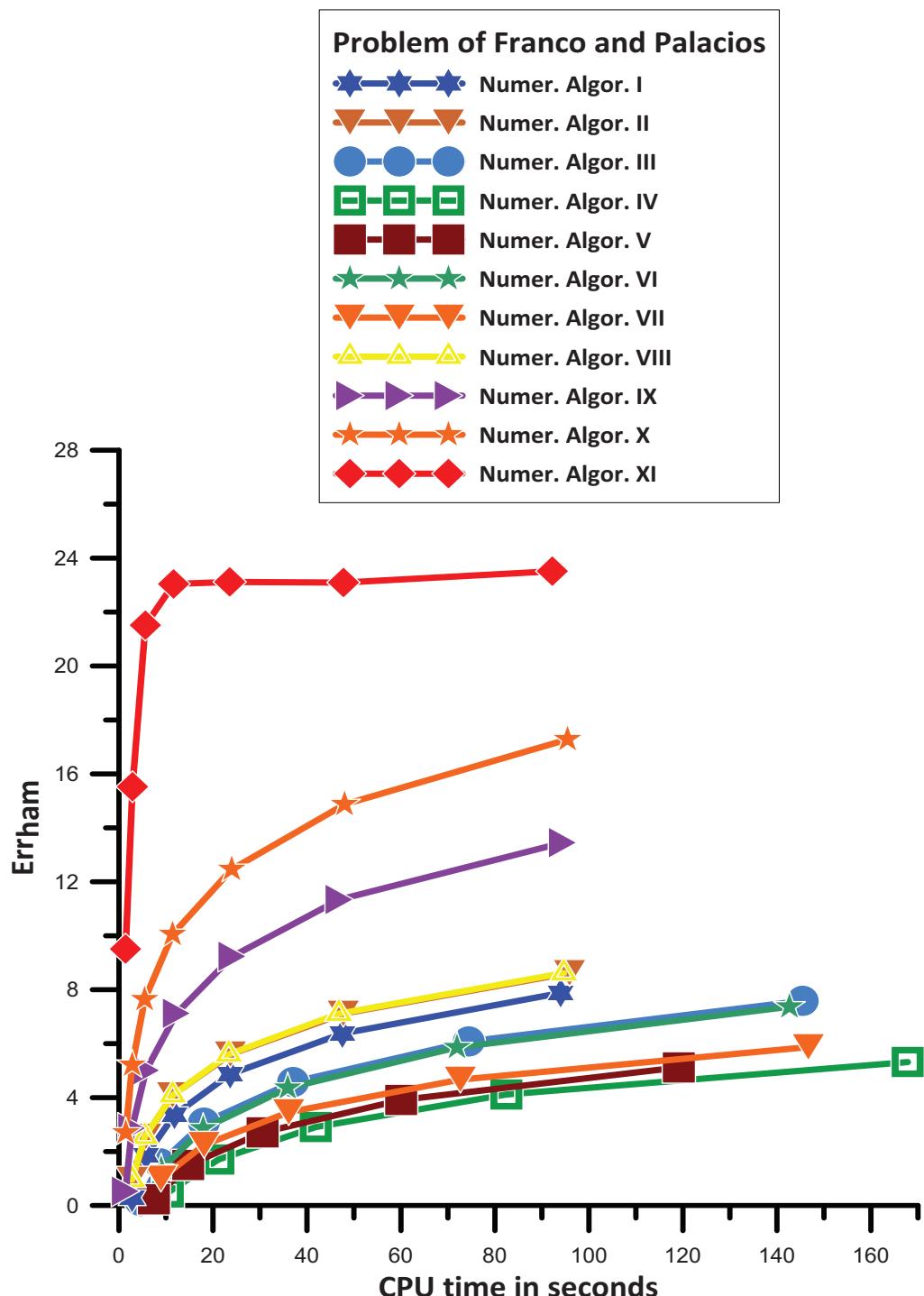


Figure 14. Numerical results for the problem of Franco and Palacios [120].

The following may be seen in Figure 15:

- Numer. Algor. IV has approximately the same efficiency as Numer. Algor. V;
- Numer. Algor. VII is more efficient than Numer. Algor. V;
- Numer. Algor. VI is more efficient than Numer. Algor. VII;
- Numer. Algor. III has approximately the same efficiency as Numer. Algor. VI;
- Numer. Algor. I is more efficient than Numer. Algor. III;
- Numer. Algor. VIII is more efficient than Numer. Algor. I;
- Numer. Algor. II has approximately the same efficiency as Numer. Algor. VIII;

- Numer. Algor. IX is more efficient than Numer. Algor. VIII;
- Numer. Algor. X is more efficient than Numer. Algor. IX;
- Numer. Algor. XI gives the most efficient results.

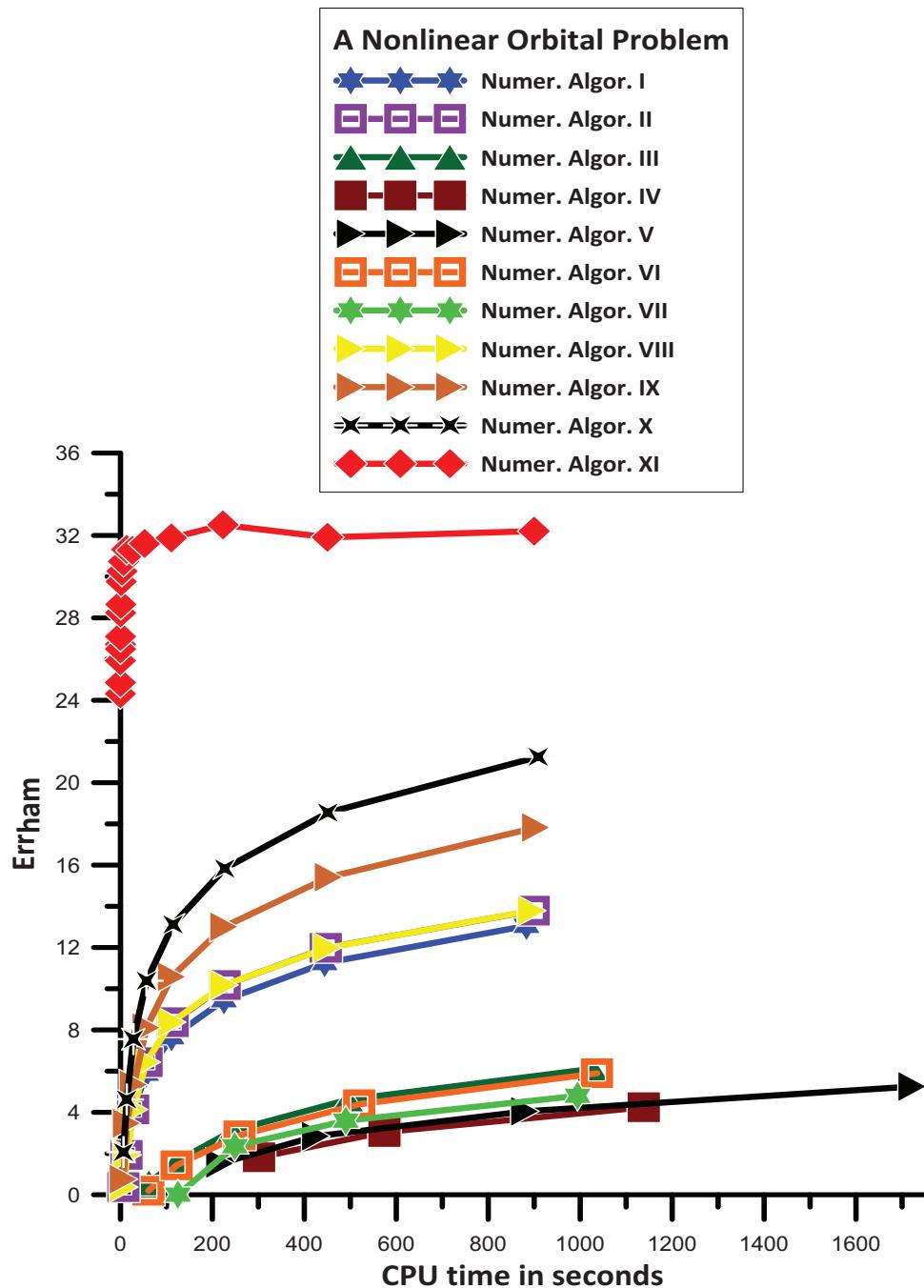


Figure 15. Numerical results for the Nonlinear Orbital problem of [121].

16.5. Nonlinear Problem of Petzold [122]

Petzold [122] investigated the following nonlinear problem, which we consider here:

$$\begin{aligned} s'_1(x) &= \lambda s_2(x), \quad s_1(0) = 1, \\ s'_2(x) &= -\lambda s_1(x) + \frac{\alpha}{\lambda} \sin(\lambda x), \quad s_2(0) = -\frac{\alpha}{2\lambda^2}. \end{aligned} \quad (103)$$

The exact solution is

$$\begin{aligned}s_1(x) &= \left(1 - \frac{\alpha}{2\lambda}x\right) \cos(\lambda x), \\ s_2(x) &= -\left(1 - \frac{\alpha}{2\lambda}x\right) \sin(\lambda x) - \frac{\alpha}{2\lambda^2} \cos(\lambda x),\end{aligned}\quad (104)$$

where $\lambda = 1000$, $\alpha = 100$. For this problem, we use $\omega = 1000$.

The numerical solution to the system of Equation (103) for $0 \leq x \leq 1000$ has been achieved by using the techniques outlined in Section 16.1.

The following may be seen in Figure 16:

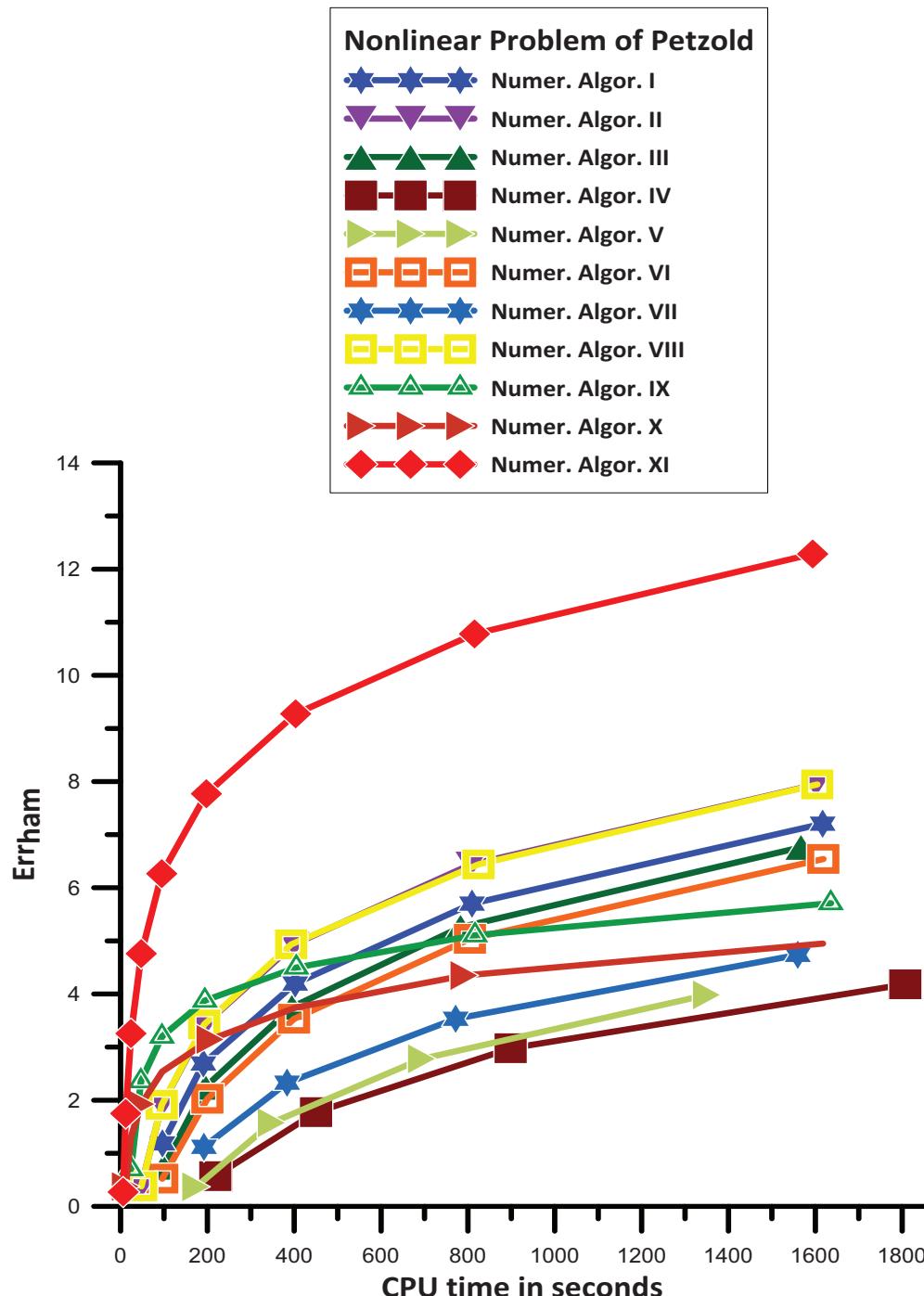


Figure 16. Numerical results for the nonlinear problem of [122].

- Numer. Algor. V is more efficient than Numer. Algor. IV;
- Numer. Algor. VII is more efficient than Numer. Algor. V;
- Numer. Algor. VI is more efficient than Numer. Algor. VII;
- Numer. Algor. III is more efficient than Numer. Algor. VI;
- Numer. Algor. I is more efficient than Numer. Algor. III;
- Numer. Algor. VIII is more efficient than Numer. Algor. I;
- Numer. Algor. II has approximately the same efficiency as Numer. Algor. VIII;
- Numer. Algor. IX has mixed behavior. For big step sizes, it is more efficient than Numer. Algor. VIII. For middle step sizes, it is more efficient than Numer. Algor. VI. For small step sizes, it is more efficient than Numer. Algor. VII;
- Numer. Algor. X has mixed behavior. For big step sizes, it is more efficient than Numer. Algor. VIII but less efficient than Numer. Algor. IX. For middle step sizes, it is more efficient than Numer. Algor. VII but less efficient than Numer. Algor. VI. For small step sizes, it is more efficient than Numer. Algor. VII;
- Numer. Algor. XI gives the most efficient results.

16.6. Two-Body Gravitational Problem

The two-body gravitational issue is under our consideration.

$$\begin{aligned} s_1''(x) &= -\frac{s_1(x)}{\left(s_1(x)^2 + s_2(x)^2\right)^{\frac{3}{2}}}, \quad s_1(0) = 1, \quad s_1'(0) = 0, \\ s_2''(x) &= -\frac{s_2(x)}{\left(s_1(x)^2 + s_2(x)^2\right)^{\frac{3}{2}}}, \quad s_2(0) = 0, \quad s_2'(0) = 1. \end{aligned} \quad (105)$$

The exact solution is

$$\begin{aligned} s_1(x) &= \cos(x), \\ s_2(x) &= \sin(x). \end{aligned} \quad (106)$$

For this problem, we use $\omega = \frac{1}{\left(s_1(x)^2 + s_2(x)^2\right)^{\frac{3}{4}}}.$

Using the techniques outlined in Section 16.1, the numerical solution to the system of Equation (105) has been obtained for $0 \leq x \leq 100,000$.

The following may be seen in Figure 17:

- Numer. Algor. I has approximately the same efficiency as Numer. Algor. VII;
- Numer. Algor. VI is more efficient than Numer. Algor. I;
- Numer. Algor. VIII is more efficient than Numer. Algor. VI;
- Numer. Algor. II has approximately the same efficiency as Numer. Algor. VIII, Numer. Algor. III, and Numer. Algor. V;
- Numer. Algor. IV is more efficient than Numer. Algor. II;
- Numer. Algor. IX is more efficient than Numer. Algor. IV;
- Numer. Algor. X is more efficient than Numer. Algor. IX;
- Numer. Algor. XI gives the most efficient results.

16.7. Perturbed Two-Body Gravitational Problem

16.7.1. Case $\mu = 0.1$

Here, we take into account the perturbed two-body Kepler's plane problem.

$$\begin{aligned}
s_1''(x) &= -\frac{s_1(x)}{\left(s_1(x)^2 + s_2(x)^2\right)^{\frac{3}{2}}} - \mu(\mu+2)\frac{s_1(x)}{\left(s_1(x)^2 + s_2(x)^2\right)^{\frac{5}{2}}}, \\
s_1(0) &= 1, \quad s_1'(0) = 0, \\
s_2''(x) &= -\frac{s_2(x)}{\left(s_1(x)^2 + s_2(x)^2\right)^{\frac{3}{2}}} - \mu(\mu+2)\frac{s_2(x)}{\left(s_1(x)^2 + s_2(x)^2\right)^{\frac{5}{2}}}, \\
s_2(0) &= 0, \quad s_2'(0) = 1 + \mu.
\end{aligned} \tag{107}$$

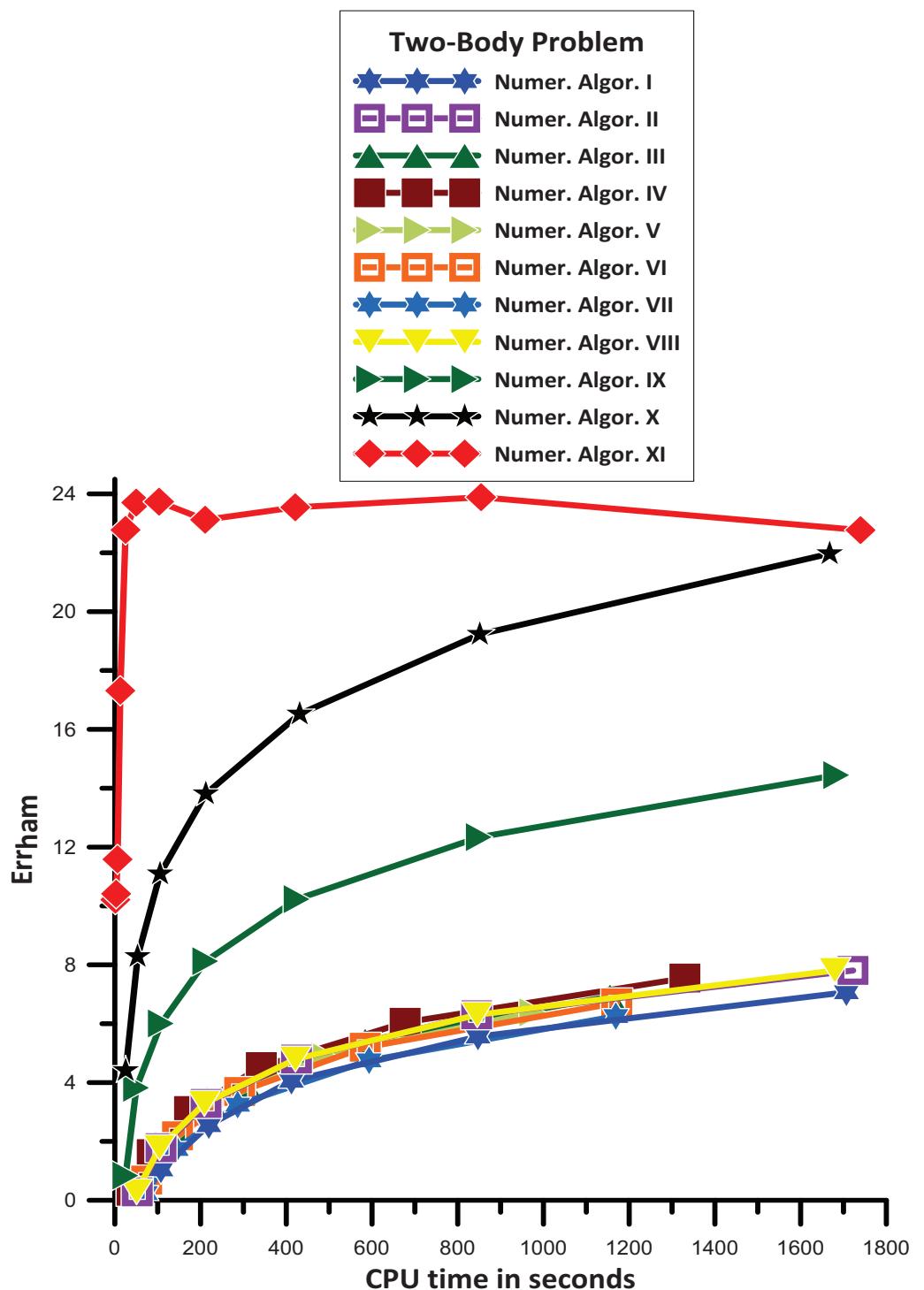


Figure 17. Numerical results for two-body gravitational problem (Kepler's plane problem).

The exact solution is

$$\begin{aligned}s_1(x) &= \cos(x + \mu x), \\ s_2(x) &= \sin(x + \mu x).\end{aligned}\quad (108)$$

For this problem, we use $\omega = \frac{\sqrt{1+\mu(\mu+2)}}{(s_1(x)^2+s_2(x)^2)^{\frac{3}{4}}}$.

Numerical solutions have been found for $0 \leq x \leq 100,000$ using $\mu = 0.1$ and the techniques described in Section 16.1 for the system of Equation (107).

The following may be seen in Figure 18:

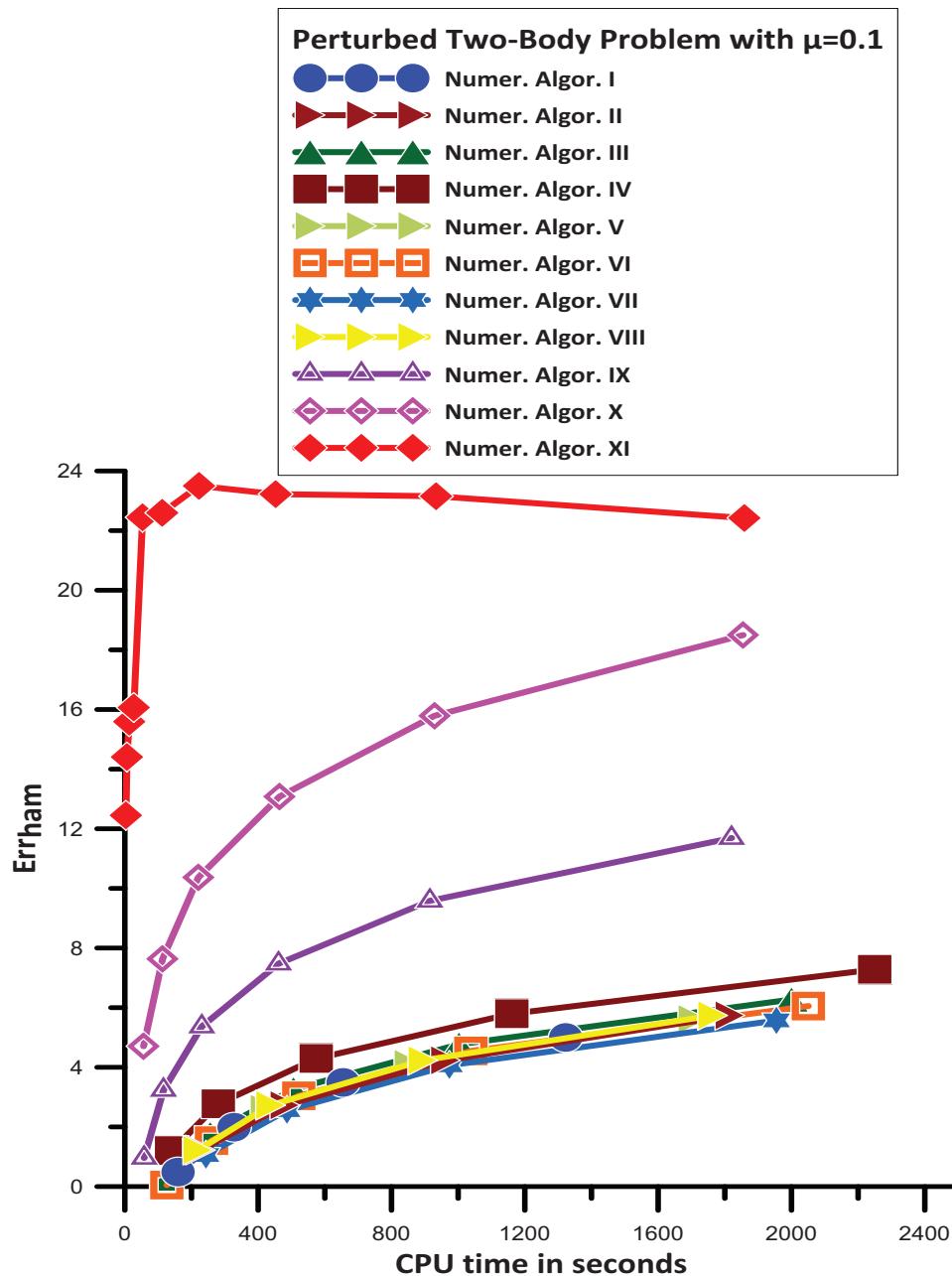


Figure 18. Numerical results for perturbed two-body gravitational problem (perturbed Kepler's problem) with $\mu = 0.1$.

- Numer. Algor. I is more efficient than Numer. Algor. VII;
- Numer. Algor. I, Numer. Algor. II, Numer. Algor. V, Numer. Algor. VI, and Numer. Algor. VIII, have approximately the same efficiency;
- Numer. Algor. III is more efficient than Numer. Algor. I;
- Numer. Algor. IV is more efficient than Numer. Algor. III;
- Numer. Algor. IX is more efficient than Numer. Algor. IV;
- Numer. Algor. X is more efficient than Numer. Algor. IX;
- Numer. Algor. XI gives the most efficient results.

16.7.2. Case $\mu = 0.4$

Numerical solutions have been found for $0 \leq x \leq 100,000$ using $\mu = 0.1$ and the techniques described in Section 16.1 for the system of Equation (107).

The following may be seen in Figure 19:

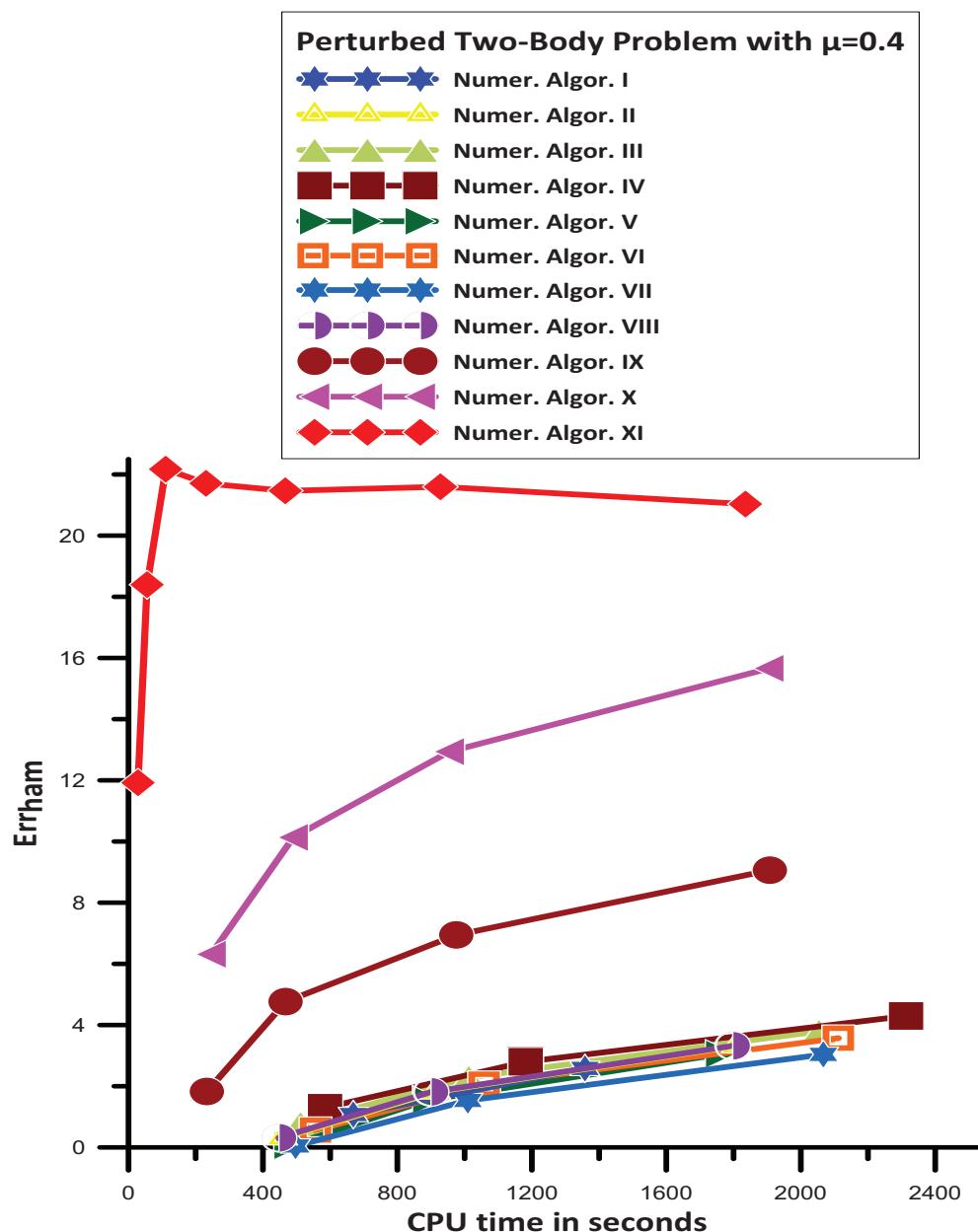


Figure 19. Numerical results for perturbed two-body gravitational problem (perturbed Kepler's problem) with $\mu = 0.4$.

- Numer. Algor. I is more efficient than Numer. Algor. VII;
- Numer. Algor. I, Numer. Algor. II, Numer. Algor. V, Numer. Algor. VI, and Numer. Algor. VIII, have approximately the same efficiency;
- Numer. Algor. III is more efficient than Numer. Algor. I;
- Numer. Algor. IV is more efficient than Numer. Algor. III;
- Numer. Algor. IX is more efficient than Numer. Algor. IV;
- Numer. Algor. X is more efficient than Numer. Algor. IX;
- Numer. Algor. XI gives the most efficient results.

Based on the numerical examples provided above, we may deduce:

- Results for all problems are most efficiently produced by the phase-fitted and amplification-fitted approach (Numer. Algor. XI);
- Results for the majority of problems are second-best when using the amplification-fitted Adams–Bashforth–Moulton Algorithm of second algebraic order with a phase-lag of order eight (Numer. Algor. X);
- Results for the majority of problems are third-best when using the amplification-fitted Adams–Bashforth–Moulton Algorithm of second algebraic order with a phase-lag of order six (Numer. Algor. IX).

In light of the above, it is clear that the strategies offered in this work that provide the best results are:

- The strategy that disregards the algebraic order of the procedure in favor of minimizing the phase-lag;
- Strategies that concentrate on eliminating phase-lag and the amplification factor

The effectiveness of frequency-dependent approaches, such as the recently presented ones, is clearly influenced by the parameter v that is chosen. This option may often be defined directly from the problem's model in many cases. The literature (e.g., [123,124]) has proposed approaches for determining the parameter v in circumstances when this is not simple.

Remark 2. One thing to keep in mind when solving systems of high-order ordinary differential equations using the recently introduced techniques is that there are already established ways to simplify such systems into first-order differential equations. For examples of such methods, see [125].

In order to solve systems of partial differential equations using the recently introduced techniques mentioned earlier, it is important to note that there are already established methods, see [126], that can reduce such a system to a system of first-order differential equations.

17. Conclusions

For implicit multistep approaches to first-order initial-value problems, we presented here the theory of phase-lag and amplification-error analysis. Several strategies for the construction of efficient predictor–corrector methods were offered, based on the theory described above and that developed in [115] for explicit methods. Our development efforts focused on the following strategies:

- Strategies for reducing the phase-lag;
- A strategy for the construction of an amplification-fitted method;
- A strategy for the construction of a phase-fitted method.

We created a number of multistep predictor–corrector approaches by using the aforementioned strategies. We based on the fourth algebraic order the Adams–Bashforth explicit method and on the fifth algebraic order Adams–Moulton implicit method.

The effectiveness of the aforementioned strategies was evaluated by applying them to many problems involving oscillating solutions.

It is worth mentioning that the idea put forward in this work and [115] is novel in the literature in relation to the development of:

- multistep methods for first-order initial-value problems with minimal phase-lag;
- phase-fitted and amplification-fitted multistep methods for first-order initial-value problems.

The same theory can be applied to all categories of multistep methods for first-order initial-value problem with oscillating solutions.

We also note that the methods presented in this paper can be applied to any problem with oscillating solution.

The computations were carried out using a 64-bit quadruple-precision arithmetic data type-compatible personal computer that conformed to the IEEE Standard 754, using a *FORTRAN* package.

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Data Availability Statement: Data are contained within the article.

Conflicts of Interest: The authors declare no conflicts of interest.

Appendix A. Development of Algorithm III Eliminating the Amplification Factor

We obtain the following result when we use the straightforward approach for computing the amplification factor (28):

$$AF = \frac{\sin(4v) - \sin(3v) - K_0(vs.)v \cos(3v) - K_1(vs.)vs. \cos(2v) - K_2(vs.)v \cos(vs.) - vK_3(vs.)}{-9v^2K_0(vs.) - 4v^2K_1(vs.) - v^2K_2(vs.)}, \quad (A1)$$

where AF denotes the amplification factor.

Assuming that the amplification factor must be eliminated, or that $AF = 0$, we derive:

$$K_0(vs.)(vs.) = \frac{-K_1(vs.)v \cos(2v) - K_2(vs.)v \cos(vs.) - vs.K_3(vs.) + \sin(4v) - \sin(3v)}{v \cos(3v)}. \quad (A2)$$

Procedure for Minimizing the Phase-Lag

We can obtain the phase-lag by plugging the values of $K_0(vs.)(vs.)$ that is previously provided into the direct formula for calculating it (19):

$$PhErr = \frac{v NUMRT_1(vs.)}{\Xi_3(vs.)}, \quad (A3)$$

where

$$\begin{aligned} NUMRT_1(vs.) &= 4 \sin(vs.) (\cos(vs.))^2 v K_3(vs.) + 2 \sin(vs.) \cos(vs.) v K_2(vs.) \\ &\quad + K_1(vs.) v \sin(vs.) - \sin(vs.) v K_3(vs.) - \cos(vs.) + 1, \\ \Xi_3(vs.) &= 8 (\cos(vs.))^3 v K_1(vs.) + 4 (\cos(vs.))^3 v K_2(vs.) \\ &\quad + 24 \sin(vs.) (\cos(vs.))^3 - 28 (\cos(vs.))^3 vs. \\ &\quad - 6 v K_1(vs.) (\cos(vs.))^2 - 12 \sin(vs.) (\cos(vs.))^2 \\ &\quad - 6 \cos(vs.) v K_1(vs.) - 6 K_2(vs.) v \cos(vs.) \\ &\quad - 12 \sin(vs.) \cos(vs.) + 21 \cos(vs.) v \\ &\quad + 3 K_1(vs.) vs. - 3 v K_3(vs.) + 3 \sin(vs.), \end{aligned} \quad (A4)$$

and $PhErr$ denotes the phase-lag.

By applying the Taylor series expansion to Formula (A4), we are able to retrieve:

$$\begin{aligned}
 PhErr &= \frac{\left(3K_3(vs.) + 2K_2(vs.) + K_1(vs.) + \frac{1}{2}\right)}{-K_1(vs.) - 2K_2(vs.) - 4 - 3K_3(vs.)} v^2 \\
 &+ \frac{\Xi_4(vs.)}{-K_1(vs.) - 2K_2(vs.) - 4 - 3K_3(vs.)} v^4 \\
 &+ \frac{\Xi_5(vs.)}{-K_1(vs.) - 2K_2(vs.) - 4 - 3K_3(vs.)} v^6 \\
 &+ \frac{\Xi_6(vs.)}{-K_1(vs.) - 2K_2(vs.) - 4 - 3K_3(vs.)} v^8 \\
 &+ \frac{\Xi_7(vs.)}{-K_1(vs.) - 2K_2(vs.) - 4 - 3K_3(vs.)} v^{10} + \dots,
 \end{aligned} \tag{A5}$$

where

$$\begin{aligned}
 \Xi_4(vs.) &= -\frac{9}{2}K_3(vs.) - \frac{4}{3}K_2(vs.) - \frac{1}{6}K_1(vs.) - \frac{1}{24} \\
 &+ \frac{1}{2} \frac{(6K_3(vs.) + 4K_2(vs.) + 2K_1(vs.) + 1)(-3K_1(vs.) - 3K_2(vs.) + 13)}{K_1(vs.) + 2K_2(vs.) + 4 + 3K_3(vs.)}, \\
 \Xi_5(vs.) &= \frac{81K_3(vs.)}{40} + \frac{4K_2(vs.)}{15} + \frac{K_1(vs.)}{120} + \frac{1}{720} \\
 &+ \frac{1}{2} \frac{6K_3(vs.) + 4K_2(vs.) + 2K_1(vs.) + 1}{K_1(vs.) + 2K_2(vs.) + 4 + 3K_3(vs.)} \left(\frac{19K_1(vs.)}{4} + \frac{13K_2(vs.)}{4} - \frac{41}{10}\right) \\
 &- \frac{1}{24} \frac{\Psi_1(vs.) (-3K_1(vs.) - 3K_2(vs.) + 13)}{(K_1(vs.) + 2K_2(vs.) + 4 + 3K_3(vs.))}, \\
 \Psi_1(vs.) &= 76K_1(vs.)^2 + 256K_1(vs.)K_2(vs.) + 336K_1(vs.)K_3(vs.) \\
 &+ 208K_2(vs.)^2 + 528K_2(vs.)K_3(vs.) - 259K_1(vs.)K_3(vs.) \\
 &+ 324K_3(vs.)^2 - 458K_2(vs.) - 501K_3(vs.) - 152, \\
 \Xi_6(vs.) &= -\frac{243K_3(vs.)}{560} - \frac{8K_2(vs.)}{315} - \frac{K_1(vs.)}{5040} - \frac{1}{40320} \\
 &+ \frac{1}{2} \frac{6K_3(vs.) + 4K_2(vs.) + 2K_1(vs.) + 1}{K_1(vs.) + 2K_2(vs.) + 4 + 3K_3(vs.)} \Psi_2(vs.) \\
 &- \frac{1}{24} \frac{\Psi_3(vs.)\Psi_4(vs.)}{(K_1(vs.) + 2K_2(vs.) + 4 + 3K_3(vs.))^2} \\
 &+ \frac{\Psi_5(vs.) (-3K_1(vs.) - 3K_2(vs.) + 13)}{720(K_1(vs.) + 2K_2(vs.) + 4 + 3K_3(vs.))}, \\
 \Psi_2(vs.) &= -\frac{211K_1(vs.)}{120} - \frac{121K_2(vs.)}{120} - \frac{229}{168}, \\
 \Psi_3(vs.) &= 76K_1(vs.)^2 + 256K_1(vs.)K_2(vs.) + 336K_1(vs.)K_3(vs.) \\
 &+ 208K_2(vs.)^2 + 528K_2(vs.)K_3(vs.) + 324K_3(vs.)^2 \\
 &- 259K_1(vs.) - 458K_2(vs.) - 501K_3(vs.) - 152, \\
 \Psi_4(vs.) &= \frac{19K_1(vs.)}{4} + \frac{13K_2(vs.)}{4} - \frac{41}{10},
 \end{aligned}$$

$$\begin{aligned}
\Psi_5(vs.) &= 10266 K_1(vs.)^3 + 46116 K_1(vs.)^2 K_2(vs.) + 52254 K_1(vs.)^2 K_3(vs.) \\
&+ 65592 K_1(vs.) K_2(vs.)^2 + 139896 K_1(vs.) K_2(vs.) K_3(vs.) + 68742 K_1(vs.) K_3(vs.)^2 \\
&+ 28848 K_2(vs.)^3 + 83736 K_2(vs.)^2 K_3(vs.) + 69444 K_2(vs.) K_3(vs.)^2 \\
&+ 13122 K_3(vs.)^3 - 40463 K_1(vs.)^2 - 133232 K_1(vs.) K_2(vs.) \\
&- 135858 K_1(vs.) K_3(vs.) - 110012 K_2(vs.)^2 - 226896 K_2(vs.) K_3(vs.) \\
&- 117927 K_3(vs.)^2 + 80990 K_1(vs.) + 146140 K_2(vs.) + 178890 K_3(vs.) + 53392, \\
\Xi_7(vs.) &= \frac{243 K_3(vs.)}{4480} + \frac{4 K_2(vs.)}{2835} + \frac{K_1(vs.)}{362880} + \frac{1}{3628800} \\
&+ \frac{1}{2} \frac{6 K_3(vs.) + 4 K_2(vs.) + 2 K_1(vs.) + 1}{K_1(vs.) + 2 K_2(vs.) + 4 + 3 K_3(vs.)} \Psi_6(vs.) \\
&- \frac{1}{24} \frac{\Psi_7(vs.)}{(K_1(vs.) + 2 K_2(vs.) + 4 + 3 K_3(vs.))^2} \Psi_8(vs.) \\
&+ \frac{1}{720} \frac{\Psi_9(vs.)}{(K_1(vs.) + 2 K_2(vs.) + 4 + 3 K_3(vs.))^3} \Psi_{10}(vs.) \\
&- \frac{1}{40320} \frac{1}{(K_1(vs.) + 2 K_2(vs.) + 4 + 3 K_3(vs.))^4} \Psi_{11}(vs.) (-3 K_1(vs.) - 3 K_2(vs.) + 13), \\
\Psi_6(vs.) &= \frac{2059 K_1(vs.)}{6720} + \frac{1093 K_2(vs.)}{6720} + \frac{2617}{3024}, \\
\Psi_7(vs.) &= 76 K_1(vs.)^2 + 256 K_1(vs.) K_2(vs.) + 336 K_1(vs.) K_3(vs.) \\
&+ 208 K_2(vs.)^2 + 528 K_2(vs.) K_3(vs.) + 324 K_3(vs.)^2 \\
&- 259 K_1(vs.) - 458 K_2(vs.) - 501 K_3(vs.) - 152, \\
\Psi_8(vs.) &= -\frac{211 K_1(vs.)}{120} - \frac{121 K_2(vs.)}{120} - \frac{229}{168}, \\
\Psi_9(vs.) &= 10266 K_1(vs.)^3 + 46116 K_1(vs.)^2 K_2(vs.) + 52254 K_1(vs.)^2 K_3(vs.) \\
&+ 65592 K_1(vs.) K_2(vs.)^2 + 139896 K_1(vs.) K_2(vs.) K_3(vs.) + 68742 K_1(vs.) K_3(vs.)^2 \\
&+ 28848 K_2(vs.)^3 + 83736 K_2(vs.)^2 K_3(vs.) + 69444 K_2(vs.) K_3(vs.)^2 \\
&+ 13122 K_3(vs.)^3 - 40463 K_1(vs.)^2 - 133232 K_1(vs.) K_2(vs.) \\
&- 135858 K_1(vs.) K_3(vs.) - 110012 K_2(vs.)^2 - 226896 K_2(vs.) K_3(vs.) \\
&- 117927 K_3(vs.)^2 + 80990 K_1(vs.) + 146140 K_2(vs.) \\
&+ 178890 K_3(vs.) + 53392, \\
\Psi_{10}(vs.) &= \frac{19 K_1(vs.)}{4} + \frac{13 K_2(vs.)}{4} - \frac{41}{10},
\end{aligned}$$

$$\begin{aligned}
\Psi_{11}(vs.) &= 2402072 K_1(vs.)^4 + 13610080 K_1(vs.)^3 K_2(vs.) \\
&+ 27841152 K_1(vs.)^2 K_2(vs.)^2 + 54097920 K_1(vs.)^2 K_2(vs.) K_3(vs.) \\
&+ 24249888 K_1(vs.)^2 K_3(vs.)^2 + 13935024 K_1(vs.)^3 K_3(vs.) \\
&+ 24184192 K_1(vs.) K_2(vs.)^3 + 65984832 K_1(vs.) K_2(vs.)^2 K_3(vs.) \\
&+ 53883360 K_1(vs.) K_2(vs.) K_3(vs.)^2 - 72514521 K_1(vs.) K_3(vs.)^2 \\
&+ 12347856 K_1(vs.) K_3(vs.)^3 + 7451264 K_2(vs.)^4 + 24880896 K_2(vs.)^3 K_3(vs.) \\
&+ 26733888 K_2(vs.)^2 K_3(vs.)^2 + 9581760 K_2(vs.) K_3(vs.)^3 + 472392 K_3(vs.)^4 \\
&- 13778147 K_1(vs.)^3 - 61836294 K_1(vs.)^2 K_2(vs.) - 59902335 K_1(vs.)^2 K_3(vs.) \\
&- 89481732 K_1(vs.) K_2(vs.)^2 - 165890772 K_1(vs.) K_2(vs.) K_3(vs.) \\
&- 40936264 K_2(vs.)^3 - 104978844 K_2(vs.)^2 K_3(vs.) - 78163110 K_2(vs.) K_3(vs.)^2 \\
&- 12874437 K_3(vs.)^3 + 35158660 K_1(vs.)^2 + 116855920 K_1(vs.) K_2(vs.) \\
&+ 116770488 K_1(vs.) K_3(vs.) - 44992832 K_1(vs.) \\
&+ 97149520 K_2(vs.)^2 + 193493136 K_2(vs.) K_3(vs.) + 94000644 K_3(vs.)^2
\end{aligned} \tag{A6}$$

$$- 83437696 K_2(vs.) - 108869952 K_3(vs.) - 34241728 + \dots \tag{A7}$$

Requiring the phase-lag to be minimized, we obtain the system of equations mentioned below:

$$\begin{aligned} \frac{\left(3K_3(vs.) + 2K_2(vs.) + K_1(vs.) + \frac{1}{2}\right)}{-K_1(vs.) - 2K_2(vs.) - 4 - 3K_3(vs.)} &= 0, \\ \frac{\Xi_4(vs.)}{-K_1(vs.) - 2K_2(vs.) - 4 - 3K_3(vs.)} &= 0, \\ \frac{\Xi_5(vs.)}{-K_1(vs.) - 2K_2(vs.) - 4 - 3K_3(vs.)} &= 0. \end{aligned} \quad (\text{A8})$$

We obtain the following result after solving the system of Equations (A8):

$$\begin{aligned} K_1(vs.) &= -\frac{179}{288}, \\ K_2(vs.) &= \frac{13}{180}, \\ K_3(vs.) &= -\frac{11}{1440}. \end{aligned} \quad (\text{A9})$$

Appendix B. Development of Algorithm VII

Eliminating the Amplification Factor

We obtain the following result when we use the straightforward approach for computing the amplification factor (21):

$$AF = \frac{\Psi_{14}(vs.)}{-v^2 Q_3(vs.) - 16v^2 Q_0(vs.) - \frac{793}{120}v^2 - 1}. \quad (\text{A10})$$

where AF denotes the amplification factor, and

$$\begin{aligned} \Psi_{14}(vs.) &= \sin(4v) - \sin(3v) - Q_0(vs.)v \cos(4v) - \frac{323}{360}v \cos(3v) \\ &+ \frac{11}{30}v \cos(2v) - Q_3(vs.)v \cos(vs.) - vQ_4(vs.). \end{aligned} \quad (\text{A11})$$

Assuming that the amplification factor must be eliminated, or that $AF = 0$, we derive:

$$Q_0(vs.) = \frac{\Psi_{15}(vs.)}{360v \cos(4v)}, \quad (\text{A12})$$

where

$$\begin{aligned} \Psi_{15}(vs.) &= -360Q_3(vs.)v \cos(vs.) - 323v \cos(3v) + 132v \cos(2v) \\ &- 360vQ_4(vs.) + 360 \sin(4v) - 360 \sin(3v). \end{aligned} \quad (\text{A13})$$

Procedure for Minimizing the Phase-Lag

We can obtain the phase-lag by plugging the value of $Q_0(vs.)$ that was previously provided into the direct formula for calculating it (19):

$$PhErr = \frac{v\Psi_{16}(vs.)}{\Psi_{17}(vs.)}, \quad (\text{A14})$$

where

$$\begin{aligned}
\Psi_{16}(vs.) &= 2880(\cos(vs.))^3 \sin(vs.) v Q_4(vs.) + 1440(\cos(vs.))^2 \sin(vs.) v Q_3(vs.) \\
&- 1440 \cos(vs.) \sin(vs.) v Q_4(vs.) - 264 \cos(vs.) \sin(vs.) vs. + 323 v \sin(vs.) \\
&- 360 Q_3(vs.) v \sin(vs.) + 360 \cos(vs.) - 360, \\
\Psi_{17}(vs.) &= 2880(\cos(vs.))^4 v Q_3(vs.) - 14520(\cos(vs.))^4 v + 11520(\cos(vs.))^3 \sin(vs.) \\
&- 5168 v (\cos(vs.))^3 - 2880(\cos(vs.))^2 v Q_3(vs.) - 5760(\cos(vs.))^2 \sin(vs.) \\
&+ 15576(\cos(vs.))^2 v - 1440 Q_3(vs.) v \cos(vs.) - 5760 \cos(vs.) \sin(vs.) \\
&+ 3876 v \cos(vs.) + 360 v Q_3(vs.) - 1440 v Q_4(vs.) + 1440 \sin(vs.) - 2343 v,
\end{aligned} \tag{A15}$$

and $PhErr$ denotes the phase-lag.

By applying the Taylor series expansion to Formula (A15), we are able to retrieve:

$$\begin{aligned}
PhErr &= \frac{(1440 Q_4(vs.) + 1080 Q_3(vs.) - 121)}{-1080 Q_3(vs.) - 1139 - 1440 Q_4(vs.)} v^2 \\
&+ \frac{1}{-1080 Q_3(vs.) - 1139 - 1440 Q_4(vs.)} v^4 \Psi_{18}(vs.) \\
&+ \frac{1}{-1080 Q_3(vs.) - 1139 - 1440 Q_4(vs.)} v^6 \Psi_{19}(vs.) + \dots,
\end{aligned} \tag{A16}$$

where

$$\begin{aligned}
\Psi_{18}(vs.) &= -3840 Q_4(vs.) - 1620 Q_3(vs.) + \frac{823}{6} \\
&+ \frac{(1440 Q_4(vs.) + 1080 Q_3(vs.) - 121)(-2160 Q_3(vs.) + 10398)}{1080 Q_3(vs.) + 1139 + 1440 Q_4(vs.)}, \\
\Psi_{19}(vs.) &= 3072 Q_4(vs.) + 729 Q_3(vs.) - \frac{3961}{120} \\
&+ \frac{(1440 Q_4(vs.) + 1080 Q_3(vs.) - 121)(3780 Q_3(vs.) - \frac{27993}{2})}{1080 Q_3(vs.) + 1139 + 1440 Q_4(vs.)} \\
&- \frac{1}{6} \frac{\Psi_{20}(vs.) (-2160 Q_3(vs.) + 10398)}{(1080 Q_3(vs.) + 1139 + 1440 Q_4(vs.))^2}, \\
\Psi_{20}(vs.) &= 24494400 Q_3(vs.)^2 + 57542400 Q_{1(vs.)} Q_4(vs.) \\
&+ 33177600 Q_4(vs.)^2 - 58764960 Q_3(vs.) \\
&- 64781280 Q_4(vs.) + 6611551.
\end{aligned} \tag{A17}$$

Requiring the phase-lag to be minimized, we obtain the system of equations mentioned below:

$$\begin{aligned}
\frac{(1440 Q_4(vs.) + 1080 Q_3(vs.) - 121)}{-1080 Q_3(vs.) - 1139 - 1440 Q_4(vs.)} &= 0, \\
\frac{1}{-1080 Q_3(vs.) - 1139 - 1440 Q_4(vs.)} v^4 \Psi_{18}(vs.) &= 0.
\end{aligned} \tag{A18}$$

The solution of the above system of equations is given by:

$$Q_3(vs.) = \frac{53}{360}, \quad Q_4(vs.) = -\frac{19}{720}. \tag{A19}$$

Appendix C. Development of Algorithm VIII Eliminating the Amplification Factor

We obtain the following result when we use the straightforward approach for computing the amplification factor (21):

$$AF = \frac{\Psi_{21}(vs.)}{-v^2 Q_3(vs.) - 4 v^2 Q_2(vs.) - 16 v^2 Q_0(vs.) - \frac{323}{40} v^2 - 1}, \quad (\text{A20})$$

where AF denotes the amplification factor, and

$$\begin{aligned} \Psi_{21}(vs.) &= \sin(4v) - \sin(3v) - Q_0(vs.)v \cos(4v) - \frac{323}{360}v \cos(3v) \\ &\quad - Q_2(vs.)v \cos(2v) - Q_3(vs.)v \cos(v) - v Q_4(vs.). \end{aligned} \quad (\text{A21})$$

Assuming that the amplification factor must be eliminated, or that $AF = 0$, we derive:

$$Q_0(vs.) = \frac{\Psi_{22}(vs.)}{360v \cos(4v)}, \quad (\text{A22})$$

where

$$\begin{aligned} \Psi_{22}(vs.) &= -360 Q_2(vs.)v \cos(2v) - 360 Q_3(vs.)v \cos(v) - 323v \cos(3v) \\ &\quad - 360v Q_4(vs.) + 360 \sin(4v) - 360 \sin(3v). \end{aligned} \quad (\text{A23})$$

Procedure for Minimizing the Phase-Lag

We can obtain the phase-lag by plugging the value of $Q_0(vs.)$ that was previously provided into the direct formula for calculating it (19):

$$PhErr = \frac{v \Psi_{23}(vs.)}{\Psi_{24}(vs.)}, \quad (\text{A24})$$

where

$$\begin{aligned} \Psi_{23}(vs.) &= 2880 \sin(vs.) (\cos(vs.))^3 v Q_4(vs.) + 1440 \sin(vs.) (\cos(vs.))^2 v Q_3(vs.) \\ &\quad + 720 \sin(vs.) \cos(vs.) v Q_2(vs.) - 1440 \sin(vs.) \cos(vs.) v Q_4(vs.) \\ &\quad - 360 Q_3(vs.) v \sin(vs.) + 323 v \sin(vs.) + 360 \cos(vs.) - 360, \\ \Psi_{24}(vs.) &= 5760 (\cos(vs.))^4 v Q_2(vs.) + 2880 (\cos(vs.))^4 v Q_3(vs.) \\ &\quad - 12408 (\cos(vs.))^4 v + 11520 \sin(vs.) (\cos(vs.))^3 - 5168 v (\cos(vs.))^3 \\ &\quad - 8640 (\cos(vs.))^2 v Q_2(vs.) - 2880 (\cos(vs.))^2 v Q_3(vs.) - 5760 \sin(vs.) (\cos(vs.))^2 \\ &\quad + 12408 (\cos(vs.))^2 v - 1440 Q_3(vs.) v \cos(vs.) - 5760 \sin(vs.) \cos(vs.) \\ &\quad + 3876 v \cos(vs.) + 2160 v Q_2(vs.) + 360 v Q_3(vs.) \\ &\quad - 1440 v Q_4(vs.) + 1440 \sin(vs.) - 1551 v, \end{aligned} \quad (\text{A25})$$

and $PhErr$ denotes the phase-lag.

By applying the Taylor series expansion to Formula (A25), we are able to retrieve:

$$\begin{aligned}
PhErr &= \frac{(1440 Q_4(vs.) + 1080 Q_3(vs.) + 720 Q_2(vs.) + 143)}{-720 Q_2(vs.) - 1080 Q_3(vs.) - 1403 - 1440 Q_4(vs.)} v^2 \\
&+ \frac{1}{-720 Q_2(vs.) - 1080 Q_3(vs.) - 1403 - 1440 Q_4(vs.)} \Psi_{25}(vs.) v^4 \\
&+ \frac{1}{-720 Q_2(vs.) - 1080 Q_3(vs.) - 1403 - 1440 Q_4(vs.)} \Psi_{27}(vs.) v^6 \\
&+ \frac{1}{-720 Q_2(vs.) - 1080 Q_3(vs.) - 1403 - 1440 Q_4(vs.)} \Psi_{29}(vs.) v^8 + \dots,
\end{aligned} \tag{A26}$$

where

$$\begin{aligned}
\Psi_{25}(vs.) &= -3840 Q_4(vs.) - 1620 Q_3(vs.) - 480 Q_2(vs.) \\
&- \frac{233}{6} + \frac{\Psi_{26}(vs.) (-2880 Q_2(vs.) - 2160 Q_3(vs.) + 9342)}{720 Q_2(vs.) + 1080 Q_3(vs.) + 1403 + 1440 Q_4(vs.)}, \\
\Psi_{26}(vs.) &= (1440 Q_4(vs.) + 1080 Q_3(vs.) + 720 Q_2(vs.) + 143), \\
\Psi_{27}(vs.) &= 3072 Q_4(vs.) + 729 Q_3(vs.) + 96 Q_2(vs.) \\
&+ \frac{263}{120} + \frac{1440 Q_4(vs.) + 1080 Q_3(vs.) + 720 Q_2(vs.) + 143}{720 Q_2(vs.) + 1080 Q_3(vs.) + 1403 + 1440 Q_4(vs.)} (6720 Q_2(vs.) \\
&+ 3780 Q_3(vs.) - \frac{23065}{2}) - \frac{1}{6} \frac{\Psi_{28}(vs.) (-2880 Q_2(vs.) - 2160 Q_3(vs.) + 9342)}{(720 Q_2(vs.) + 1080 Q_3(vs.) + 1403 + 1440 Q_4(vs.))^2}, \\
\Psi_{28}(vs.) &= 14515200 Q_2(vs.)^2 + 38102400 Q_2(vs.) Q_3(vs.) + 45619200 Q_2(vs.) Q_4(vs.) \\
&+ 24494400 Q_3(vs.)^2 + 57542400 Q_3(vs.) Q_4(vs.) + 33177600 Q_4(vs.)^2 \\
&- 33678000 Q_2(vs.) - 44794080 Q_3(vs.) - 48054240 Q_4(vs.) - 7688537, \\
\Psi_{29}(vs.) &= -\frac{8192 Q_4(vs.)}{7} - \frac{2187 Q_3(vs.)}{14} - \frac{64 Q_2(vs.)}{7} \\
&- \frac{139}{2520} + \frac{1440 Q_4(vs.) + 1080 Q_3(vs.) + 720 Q_2(vs.) + 143}{720 Q_2(vs.) + 1080 Q_3(vs.) + 1403 + 1440 Q_4(vs.)} \Psi_{30}(vs.) \\
&- \frac{1}{6} \frac{\Psi_{31}(vs.)}{(720 Q_2(vs.) + 1080 Q_3(vs.) + 1403 + 1440 Q_4(vs.))^2} \Psi_{32}(vs.) \\
&+ \frac{\Psi_{33}(vs.) (-2880 Q_2(vs.) - 2160 Q_3(vs.) + 9342)}{120 (720 Q_2(vs.) + 1080 Q_3(vs.) + 1403 + 1440 Q_4(vs.))^3}, \\
\Psi_{30}(vs.) &= -3968 Q_2(vs.) - 2046 Q_3(vs.) + \frac{2551639}{420}, \\
\Psi_{31}(vs.) &= 14515200 Q_2(vs.)^2 + 38102400 Q_2(vs.) Q_3(vs.) + 45619200 Q_2(vs.) Q_4(vs.) \\
&+ 24494400 Q_3(vs.)^2 + 57542400 Q_3(vs.) Q_4(vs.) + 33177600 Q_4(vs.)^2 \\
&- 33678000 Q_2(vs.) - 44794080 Q_3(vs.) - 48054240 Q_4(vs.) - 7688537, \\
\Psi_{32}(vs.) &= 6720 Q_2(vs.) + 3780 Q_3(vs.) - \frac{23065}{2},
\end{aligned}$$

$$\begin{aligned}
\Psi_{33}(vs.) &= 1260085248000 Q_2(vs.)^3 + 4374279936000 Q_2(vs.)^2 Q_3(vs.) \\
&+ 4852410624000 Q_2(vs.) Q_3(vs.)^2 + 9524542464000 Q_2(vs.) Q_3(vs.) Q_4(vs.) \\
&+ 4371480576000 Q_2(vs.) Q_4(vs.)^2 + 1689273792000 Q_3(vs.)^3 \\
&+ 4598788608000 Q_3(vs.)^2 Q_4(vs.) + 4514807808000 Q_2(vs.)^2 Q_4(vs.) \\
&+ 3701873664000 Q_3(vs.) Q_4(vs.)^2 + 764411904000 Q_4(vs.)^3 - 4448256134400 Q_2(vs.)^2 \\
&- 11242923379200 Q_2(vs.) Q_3(vs.) - 11573966361600 Q_2(vs.) Q_4(vs.) \\
&- 7103038910400 Q_3(vs.)^2 - 14650375910400 Q_3(vs.) Q_4(vs.) \\
&- 7578470937600 Q_4(vs.)^2 + 4494085662240 Q_2(vs.) + 1159393471547 \\
&+ 5990509706760 Q_3(vs.) + 6624257647680 Q_4(vs.). \tag{A27}
\end{aligned}$$

Requiring the phase-lag to be minimized, we obtain the system of equations mentioned below:

$$\begin{aligned}
\frac{(1440 Q_4(vs.) + 1080 Q_3(vs.) + 720 Q_2(vs.) + 143)}{-720 Q_2(vs.) - 1080 Q_3(vs.) - 1403 - 1440 Q_4(vs.)} &= 0, \\
\frac{1}{-720 Q_2(vs.) - 1080 Q_3(vs.) - 1403 - 1440 Q_4(vs.)} \Psi_{25}(vs.) &= 0, \\
\frac{1}{-720 Q_2(vs.) - 1080 Q_3(vs.) - 1403 - 1440 Q_4(vs.)} \Psi_{27}(vs.) &= 0. \tag{A29}
\end{aligned}$$

The solution of the above system of equations is given by:

$$Q_2(vs.) = -\frac{167}{480}, \quad Q_3(vs.) = \frac{317}{2520}, \quad Q_4(vs.) = -\frac{397}{20160}. \tag{A30}$$

Appendix D. Development of Algorithm IX Eliminating the Amplification Factor

We obtain the following result when we use the straightforward approach for computing the amplification factor (21):

$$AF = \frac{\Psi_{35}(vs.)}{-16 v^2 Q_0(vs.) - 9 v^2 Q_1(vs.) - 4 v^2 Q_2(vs.) - v^2 Q_3(vs.) - 1}, \tag{A31}$$

where AF denotes the amplification factor, and

$$\begin{aligned}
\Psi_{35}(vs.) &= \sin(4v) - \sin(3v) - Q_0(vs.) v \cos(4v) - Q_1(vs.) v \cos(3v) \\
&- Q_2(vs.) v \cos(2v) - Q_3(vs.) v \cos(v) - v Q_4(vs.). \tag{A32}
\end{aligned}$$

Assuming that the amplification factor must be eliminated, or that $AF = 0$, we derive:

$$Q_0(vs.) = \frac{\Psi_{36}(vs.)}{v \cos(4v)}, \tag{A33}$$

where

$$\begin{aligned}
\Psi_{36}(vs.) &= -Q_1(vs.) v \cos(3v) - Q_2(vs.) v \cos(2v) - Q_3(vs.) v \cos(v) \\
&- v Q_4(vs.) + \sin(4v) - \sin(3v). \tag{A34}
\end{aligned}$$

Procedure for Minimizing the Phase-Lag

We can obtain the phase-lag by plugging the value of $Q_0(vs.)$ that was previously provided into the direct formula for calculating it (19):

$$PhErr = \frac{\Psi_{37}(vs.)}{\Psi_{38}(vs.)}, \quad (\text{A35})$$

where

$$\begin{aligned} \Psi_{37}(vs.) &= vs. \left(8 \sin(vs.) (\cos(vs.))^3 v Q_4(vs.) + 4 \sin(vs.) (\cos(vs.))^2 v Q_3(vs.) \right. \\ &\quad + 2 \sin(vs.) \cos(vs.) v Q_2(vs.) - 4 \sin(vs.) \cos(vs.) v Q_4(vs.) \\ &\quad \left. + v Q_1(vs.) \sin(vs.) - Q_3(vs.) v \sin(vs.) + \cos(vs.) - 1 \right), \end{aligned}$$

$$\begin{aligned} \Psi_{38}(vs.) &= 24 (\cos(vs.))^4 v Q_1(vs.) + 16 (\cos(vs.))^4 v Q_2(vs.) \\ &\quad + 8 (\cos(vs.))^4 v Q_3(vs.) - 56 (\cos(vs.))^4 v \\ &\quad - 16 v Q_1(vs.) (\cos(vs.))^3 + 32 \sin(vs.) (\cos(vs.))^3 \\ &\quad - 24 (\cos(vs.))^2 v Q_1(vs.) - 24 (\cos(vs.))^2 v Q_2(vs.) \\ &\quad - 8 (\cos(vs.))^2 v Q_3(vs.) - 16 \sin(vs.) (\cos(vs.))^2 \\ &\quad + 56 (\cos(vs.))^2 v + 12 \cos(vs.) v Q_1(vs.) \\ &\quad - 4 Q_3(vs.) v \cos(vs.) - 16 \sin(vs.) \cos(vs.) + 3 v Q_1(vs.) \\ &\quad + 6 Q_2(vs.) v + Q_3(vs.) v - 4 v Q_4(vs.) + 4 \sin(vs.) - 7 v, \quad (\text{A36}) \end{aligned}$$

and $PhErr$ denotes the phase-lag.

By applying the Taylor series expansion to Formula (A36), we are able to retrieve:

$$\begin{aligned} PhErr &= \frac{4 Q_4(vs.) + 3 Q_3(vs.) + 2 Q_2(vs.) + Q_1(vs.) - \frac{1}{2}}{-Q_1(vs.) - 2 Q_2(vs.) - 3 Q_3(vs.) - 3 - 4 Q_4(vs.)} v^2 \\ &\quad + \frac{\Psi_{39}(vs.)}{-Q_1(vs.) - 2 Q_2(vs.) - 3 Q_3(vs.) - 3 - 4 Q_4(vs.)} v^4 \\ &\quad + \frac{\Psi_{41}(vs.)}{-Q_1(vs.) - 2 Q_2(vs.) - 3 Q_3(vs.) - 3 - 4 Q_4(vs.)} v^6 \\ &\quad + \frac{\Psi_{46}(vs.)}{-Q_1(vs.) - 2 Q_2(vs.) - 3 Q_3(vs.) - 3 - 4 Q_4(vs.)} v^8 \\ &\quad + \frac{\Psi_{52}(vs.)}{-Q_1(vs.) - 2 Q_2(vs.) - 3 Q_3(vs.) - 3 - 4 Q_4(vs.)} v^{10} + \dots, \quad (\text{A37}) \end{aligned}$$

where

$$\begin{aligned}
\Psi_{39}(vs.) &= -\frac{32 Q_4(vs.)}{3} - \frac{9}{2} Q_3(vs.) - \frac{4}{3} Q_2(vs.) - \frac{1}{6} Q_1(vs.) + \frac{1}{24} + \Psi_{40}(vs.), \\
\Psi_{40}(vs.) &= \frac{1}{2} \frac{8 Q_4(vs.) + 6 Q_3(vs.) + 4 Q_2(vs.) + 2 Q_1(vs.) - 1}{Q_1(vs.) + 2 Q_2(vs.) + 3 Q_3(vs.) + 3 + 4 Q_4(vs.)} \\
&\quad \left(-6 Q_1(vs.) - 8 Q_2(vs.) - 6 Q_3(vs.) + \frac{94}{3} \right), \\
\Psi_{41}(vs.) &= \frac{128}{15} Q_4(vs.) + \frac{81}{40} Q_3(vs.) + \frac{4}{15} Q_2(vs.) \\
&+ \frac{Q_1(vs.)}{120} - \frac{1}{720} + \frac{1}{2} \Psi_{42}(vs.) \Psi_{43}(vs.) \\
&- \frac{1}{24} \frac{\Psi_{44}(vs.)}{(Q_1(vs.) + 2 Q_2(vs.) + 3 Q_3(vs.) + 3 + 4 Q_4(vs.))^2} \Psi_{45}(vs.), \\
\Psi_{42}(vs.) &= \frac{8 Q_4(vs.) + 6 Q_3(vs.) + 4 Q_2(vs.) + 2 Q_1(vs.) - 1}{Q_1(vs.) + 2 Q_2(vs.) + 3 Q_3(vs.) + 3 + 4 Q_4(vs.)}, \\
\Psi_{43}(vs.) &= \frac{37}{2} Q_1(vs.) + \frac{56}{3} Q_2(vs.) + \frac{21}{2} Q_3(vs.) - \frac{1459}{30}, \\
\Psi_{44}(vs.) &= 148 Q_1(vs.)^2 + 520 Q_1(vs.) Q_2(vs.) + 696 Q_1(vs.) Q_3(vs.) \\
&+ 848 Q_1(vs.) Q_4(vs.) + 448 Q_2(vs.)^2 + 1176 Q_2(vs.) Q_3(vs.) \\
&+ 1408 Q_2(vs.) Q_4(vs.) + 756 Q_3(vs.)^2 + 1776 Q_3(vs.) Q_4(vs.) \\
&+ 1024 Q_4(vs.)^2 - 813 Q_1(vs.) - 1506 Q_2(vs.) \\
&- 2007 Q_3(vs.) - 2244 Q_4(vs.) + 373, \\
\Psi_{45}(vs.) &= -6 Q_1(vs.) - 8 Q_2(vs.) - 6 Q_3(vs.) + \frac{94}{3}, \\
\Psi_{46}(vs.) &= -6 Q_1(vs.) - 8 Q_2(vs.) - 6 Q_3(vs.) + \frac{94}{3} \\
&- \frac{1024}{315} Q_4(vs.) - \frac{243}{560} Q_3(vs.) - \frac{8}{315} Q_2(vs.) \\
&- \frac{Q_1(vs.)}{5040} + \frac{1}{40320} \\
&+ \frac{1}{2} \frac{8 Q_4(vs.) + 6 Q_3(vs.) + 4 Q_2(vs.) + 2 Q_1(vs.) - 1}{Q_1(vs.) + 2 Q_2(vs.) + 3 Q_3(vs.) + 3 + 4 Q_4(vs.)} \Psi_{47}(vs.) \\
&- \frac{1}{24} \frac{\Psi_{48}(vs.)}{(Q_1(vs.) + 2 Q_2(vs.) + 3 Q_3(vs.) + 3 + 4 Q_4(vs.))^2} \Psi_{49}(vs.) \\
&+ \frac{\Psi_{50}(vs.)}{720 (Q_1(vs.) + 2 Q_2(vs.) + 3 Q_3(vs.) + 3 + 4 Q_4(vs.))^3} \Psi_{51}(vs.), \\
\Psi_{47}(vs.) &= -\frac{781}{60} Q_1(vs.) - \frac{496}{45} Q_2(vs.) - \frac{341}{60} Q_3(vs.) + \frac{11993}{420}, \\
\Psi_{48}(vs.) &= 148 Q_1(vs.)^2 + 520 Q_1(vs.) Q_2(vs.) + 696 Q_1(vs.) Q_3(vs.) \\
&+ 848 Q_1(vs.) Q_4(vs.) + 448 Q_2(vs.)^2 + 1176 Q_2(vs.) Q_3(vs.) \\
&+ 1408 Q_2(vs.) Q_4(vs.) + 756 Q_3(vs.)^2 + 1776 Q_3(vs.) Q_4(vs.) \\
&+ 1024 Q_4(vs.)^2 - 813 Q_1(vs.) - 1506 Q_2(vs.) \\
&- 2007 Q_3(vs.) - 2244 Q_4(vs.) + 373, \\
\Psi_{49}(vs.) &= \frac{37}{2} Q_1(vs.) + \frac{56}{3} Q_2(vs.) + \frac{21}{2} Q_3(vs.) - \frac{1459}{30},
\end{aligned}$$

$$\begin{aligned}
\Psi_{50}(vs.) &= 39966 Q_1(vs.)^3 + 196056 Q_1(vs.)^2 Q_2(vs.) + 240894 Q_1(vs.)^2 Q_3(vs.) \\
&+ 265392 Q_1(vs.)^2 Q_4(vs.) + 313272 Q_1(vs.) Q_2(vs.)^2 + 750096 Q_1(vs.) Q_2(vs.) Q_3(vs.) \\
&+ 803808 Q_1(vs.) Q_2(vs.) Q_4(vs.) + 435402 Q_1(vs.) Q_3(vs.)^2 + 901152 Q_1(vs.) Q_3(vs.) Q_4(vs.) \\
&+ 446688 Q_1(vs.) Q_4(vs.)^2 + 162048 Q_2(vs.)^3 + 562536 Q_2(vs.)^2 Q_3(vs.) \\
&+ 580608 Q_2(vs.)^2 Q_4(vs.) + 624024 Q_2(vs.) Q_3(vs.)^2 + 1224864 Q_2(vs.) Q_3(vs.) Q_4(vs.) \\
&+ 562176 Q_2(vs.) Q_4(vs.)^2 + 217242 Q_3(vs.)^3 + 591408 Q_3(vs.)^2 Q_4(vs.) \\
&+ 476064 Q_3(vs.) Q_4(vs.)^2 + 98304 Q_4(vs.)^3 - 287141 Q_1(vs.)^2 \\
&- 993644 Q_1(vs.) Q_2(vs.) - 1244286 Q_1(vs.) Q_3(vs.) - 1310968 Q_1(vs.) Q_4(vs.) \\
&- 853124 Q_2(vs.)^2 - 2118852 Q_2(vs.) Q_3(vs.) - 2209616 Q_2(vs.) Q_4(vs.) \\
&- 1304109 Q_3(vs.)^2 - 2692584 Q_3(vs.) Q_4(vs.) - 1375376 Q_4(vs.)^2 \\
&+ 723888 Q_1(vs.) + 1311636 Q_2(vs.) + 1692864 Q_3(vs.) \\
&+ 1814472 Q_4(vs.) - 298105, \\
\Psi_{51}(vs.) &= -6 Q_1(vs.) - 8 Q_2(vs.) - 6 Q_3(vs.) + \frac{94}{3}, \\
\Psi_{52}(vs.) &= \frac{2048 Q_4(vs.)}{2835} + \frac{243 Q_3(vs.)}{4480} + \frac{4 Q_2(vs.)}{2835} \\
&+ \frac{Q_1(vs.)}{362880} - \frac{1}{3628800} + \frac{1}{2} \frac{8 Q_4(vs.) + 6 Q_3(vs.) + 4 Q_2(vs.) + 2 Q_1(vs.) - 1}{Q_1(vs.) + 2 Q_2(vs.) + 3 Q_3(vs.) + 3 + 4 Q_4(vs.)} \Psi_{53}(vs.) \\
&- \frac{1}{24} \frac{\Psi_{54}(vs.)}{(Q_1(vs.) + 2 Q_2(vs.) + 3 Q_3(vs.) + 3 + 4 Q_4(vs.))^2} \Psi_{55}(vs.) \\
&+ \frac{1}{720} \frac{\Psi_{56}(vs.)}{(Q_1(vs.) + 2 Q_2(vs.) + 3 Q_3(vs.) + 3 + 4 Q_4(vs.))^3} \Psi_{57}(vs.) \\
&- \frac{1}{120960} \frac{\Psi_{58}(vs.)}{(Q_1(vs.) + 2 Q_2(vs.) + 3 Q_3(vs.) + 3 + 4 Q_4(vs.))^4} \Psi_{59}(vs.), \\
\Psi_{53}(vs.) &= \frac{14197}{3360} Q_1(vs.) + \frac{1016}{315} Q_2(vs.) \\
&+ \frac{5461}{3360} Q_3(vs.) - \frac{789731}{90720}, \\
\Psi_{54}(vs.) &= 148 Q_1(vs.)^2 + 520 Q_1(vs.) Q_2(vs.) + 696 Q_1(vs.) Q_3(vs.) \\
&+ 848 Q_1(vs.) Q_4(vs.) + 448 Q_2(vs.)^2 + 1176 Q_2(vs.) Q_3(vs.) \\
&+ 1408 Q_2(vs.) Q_4(vs.) + 756 Q_3(vs.)^2 + 1776 Q_3(vs.) Q_4(vs.) \\
&+ 1024 Q_4(vs.)^2 - 813 Q_1(vs.) - 1506 Q_2(vs.) \\
&- 2007 Q_3(vs.) - 2244 Q_4(vs.) + 373, \\
\Psi_{55}(vs.) &= -\frac{781}{60} Q_1(vs.) - \frac{496}{45} Q_2(vs.) \\
&- \frac{341}{60} Q_3(vs.) + \frac{11993}{420}, \\
\Psi_{56}(vs.) &= 39966 Q_1(vs.)^3 + 196056 Q_1(vs.)^2 Q_2(vs.) + 240894 Q_1(vs.)^2 Q_3(vs.) \\
&+ 265392 Q_1(vs.)^2 Q_4(vs.) + 313272 Q_1(vs.) Q_2(vs.)^2 + 750096 Q_1(vs.) Q_2(vs.) Q_3(vs.) \\
&+ 803808 Q_1(vs.) Q_2(vs.) Q_4(vs.) + 435402 Q_1(vs.) Q_3(vs.)^2 + 901152 Q_1(vs.) Q_3(vs.) Q_4(vs.) \\
&+ 446688 Q_1(vs.) Q_4(vs.)^2 + 162048 Q_2(vs.)^3 + 562536 Q_2(vs.)^2 Q_3(vs.) \\
&+ 580608 Q_2(vs.)^2 Q_4(vs.) + 624024 Q_2(vs.) Q_3(vs.)^2 + 1224864 Q_2(vs.) Q_3(vs.) Q_4(vs.) \\
&+ 562176 Q_2(vs.) Q_4(vs.)^2 + 217242 Q_3(vs.)^3 + 591408 Q_3(vs.)^2 Q_4(vs.) \\
&+ 476064 Q_3(vs.) Q_4(vs.)^2 + 98304 Q_4(vs.)^3 - 287141 Q_1(vs.)^2 \\
&- 993644 Q_1(vs.) Q_2(vs.) - 1244286 Q_1(vs.) Q_3(vs.) - 1310968 Q_1(vs.) Q_4(vs.) \\
&- 853124 Q_2(vs.)^2 - 2118852 Q_2(vs.) Q_3(vs.) - 2209616 Q_2(vs.) Q_4(vs.) \\
&- 1304109 Q_3(vs.)^2 - 2692584 Q_3(vs.) Q_4(vs.) - 1375376 Q_4(vs.)^2 \\
&+ 723888 Q_1(vs.) + 1311636 Q_2(vs.) + 1692864 Q_3(vs.) \\
&+ 1814472 Q_4(vs.) - 298105, \\
\Psi_{57}(vs.) &= \frac{37}{2} Q_1(vs.) + \frac{56}{3} Q_2(vs.) + \frac{21}{2} Q_3(vs.) - \frac{1459}{30},
\end{aligned}$$

$$\begin{aligned}
\Psi_{58}(vs.) &= 55659768 Q_1(vs.)^4 + 352129872 Q_1(vs.)^3 Q_2(vs.) \\
&+ 412143264 Q_1(vs.)^3 Q_3(vs.) + 1885960656 Q_1(vs.)^2 Q_2(vs.) Q_3(vs.) \\
&+ 421068192 Q_1(vs.)^3 Q_4(vs.) + 821700000 Q_1(vs.)^2 Q_2(vs.)^2 \\
&+ 1879728768 Q_1(vs.)^2 Q_2(vs.) Q_4(vs.) + 1056381264 Q_1(vs.)^2 Q_3(vs.)^2 \\
&+ 2040241824 Q_1(vs.)^2 Q_3(vs.) Q_4(vs.) + 2816335296 Q_1(vs.) Q_2(vs.)^2 Q_3(vs.) \\
&+ 942304896 Q_1(vs.)^2 Q_4(vs.)^2 + 836548032 Q_1(vs.) Q_2(vs.)^3 \\
&+ 2727370368 Q_1(vs.) Q_2(vs.)^2 Q_4(vs.) + 3073982832 Q_1(vs.) Q_2(vs.) Q_3(vs.)^2 \\
&+ 5734985472 Q_1(vs.) Q_2(vs.) Q_3(vs.) Q_4(vs.) + 2893125600 Q_1(vs.) Q_3(vs.)^2 Q_4(vs.) \\
&+ 2532734208 Q_1(vs.) Q_2(vs.) Q_4(vs.)^2 + 1082328480 Q_1(vs.) Q_3(vs.)^3 \\
&+ 2403198720 Q_1(vs.) Q_3(vs.) Q_4(vs.)^2 + 600645120 Q_1(vs.) Q_4(vs.)^3 \\
&+ 312778752 Q_2(vs.)^4 + 2174279328 Q_2(vs.)^2 Q_3(vs.)^2 \\
&+ 1368555840 Q_2(vs.)^3 Q_3(vs.) + 1281146880 Q_2(vs.)^3 Q_4(vs.) \\
&+ 3891547008 Q_2(vs.)^2 Q_3(vs.) Q_4(vs.) + 1625260032 Q_2(vs.)^2 Q_4(vs.)^2 \\
&+ 1477117296 Q_2(vs.) Q_3(vs.)^3 + 640745472 Q_2(vs.) Q_4(vs.)^3 \\
&+ 3747821184 Q_2(vs.) Q_3(vs.)^2 Q_4(vs.) + 2890374912 Q_2(vs.) Q_3(vs.) Q_4(vs.)^2 \\
&+ 358980984 Q_3(vs.)^4 + 1128657888 Q_3(vs.)^3 Q_4(vs.) \\
&+ 1167405696 Q_3(vs.)^2 Q_4(vs.)^2 - 565295331 Q_1(vs.)^3 \\
&+ 419830272 Q_3(vs.) Q_4(vs.)^3 + 25165824 Q_4(vs.)^4 \\
&- 2779573266 Q_1(vs.)^2 Q_2(vs.) - 3305432331 Q_1(vs.)^2 Q_3(vs.) \\
&- 3315314340 Q_1(vs.)^2 Q_4(vs.) - 10309349520 Q_1(vs.) Q_2(vs.) Q_4(vs.) \\
&- 4489981092 Q_1(vs.) Q_2(vs.)^2 - 10500895404 Q_1(vs.) Q_2(vs.) Q_3(vs.) \\
&- 6022349649 Q_1(vs.) Q_3(vs.)^2 - 11541767256 Q_1(vs.) Q_3(vs.) Q_4(vs.) \\
&- 5369568912 Q_1(vs.) Q_4(vs.)^2 - 7780962960 Q_2(vs.)^2 Q_4(vs.) \\
&- 2374354968 Q_2(vs.)^3 - 8152860204 Q_2(vs.)^2 Q_3(vs.) \\
&- 9095875938 Q_2(vs.) Q_3(vs.)^2 - 16783131312 Q_2(vs.) Q_3(vs.) Q_4(vs.) \\
&- 7402899744 Q_2(vs.) Q_4(vs.)^2 - 7043464368 Q_3(vs.) Q_4(vs.)^2 \\
&- 3267603585 Q_3(vs.)^3 - 8612711748 Q_3(vs.)^2 Q_4(vs.) \\
&- 1685883072 Q_4(vs.)^3 + 2129452501 Q_1(vs.)^2 \\
&+ 7315812004 Q_1(vs.) Q_2(vs.) + 6247837684 Q_2(vs.)^2 \\
&+ 8917489566 Q_1(vs.) Q_3(vs.) + 8861348168 Q_1(vs.) Q_4(vs.) \\
&+ 15124089132 Q_2(vs.) Q_3(vs.) + 14867966416 Q_2(vs.) Q_4(vs.) \\
&+ 9077381469 Q_3(vs.)^2 - 3528107889 Q_1(vs.) \\
&+ 17635437144 Q_3(vs.) Q_4(vs.) + 8425636816 Q_4(vs.)^2 \\
&- 6322631682 Q_2(vs.) - 8014817763 Q_3(vs.) \\
&- 8339272836 Q_4(vs.) + 1310487239, \\
\Psi_{59}(vs.) &= -6 Q_1(vs.) - 8 Q_2(vs.) - 6 Q_3(vs.) + \frac{94}{3}. \tag{A38}
\end{aligned}$$

Requiring the phase-lag to be minimized, we obtain the system of equations mentioned below:

$$\begin{aligned}
\frac{4 Q_4(vs.) + 3 Q_3(vs.) + 2 Q_2(vs.) + Q_1(vs.) - \frac{1}{2}}{-Q_1(vs.) - 2 Q_2(vs.) - 3 Q_3(vs.) - 3 - 4 Q_4(vs.)} &= 0, \\
\frac{\Psi_{39}(vs.)}{-Q_1(vs.) - 2 Q_2(vs.) - 3 Q_3(vs.) - 3 - 4 Q_4(vs.)} &= 0, \\
\frac{\Psi_{41}(vs.)}{-Q_1(vs.) - 2 Q_2(vs.) - 3 Q_3(vs.) - 3 - 4 Q_4(vs.)} &= 0, \\
\frac{\Psi_{46}(vs.)}{-Q_1(vs.) - 2 Q_2(vs.) - 3 Q_3(vs.) - 3 - 4 Q_4(vs.)} &= 0. \tag{A39}
\end{aligned}$$

The solution of the above system of equations is given by:

$$\begin{aligned} Q_1(v_s.) &= \frac{5561}{8640}, \quad Q_2(v_s.) = -\frac{163}{1728}, \\ Q_3(v_s.) &= \frac{23}{1344}, \quad Q_4(v_s.) = -\frac{191}{120960}. \end{aligned} \quad (\text{A40})$$

References

- Landau, L.D.; Lifshitz, F.M. *Quantum Mechanics*; Pergamon: New York, NY, USA, 1965.
- Prigogine, I.; Rice, S. (Eds.) *Advances in Chemical Physics Vol. 93: New Methods in Computational Quantum Mechanics*; John Wiley & Sons: Hoboken, NJ, USA, 1997.
- Simos, T.E. Numerical Solution of Ordinary Differential Equations with Periodical Solution. Ph.D. Thesis, National Technical University of Athens, Athens, Greece, 1990. (In Greek)
- Raptis, A.D. Exponential multistep methods for ordinary differential equations. *Bull. Greek Math. Soc.* **1984**, *25*, 113–126.
- Vigo-Aguiar, J.; Ferrandiz, J.M. A general procedure for the adaptation of multistep algorithms to the integration of oscillatory problems. *SIAM J. Numer. Anal.* **1998**, *35*, 1684–1708. [[CrossRef](#)]
- Vigo-Aguiar, J. Mathematical Methods for the Numerical Propagation of Satellite Orbits. Ph.D. Thesis, University of Valladolid, Valladolid, Spain, 1993. (In Spanish)
- Ixaru, L.G. *Numerical Methods for Differential Equations and Applications*; Reidel: Dordrecht, The Netherlands; Boston, MA, USA; Lancaster, UK, 1984.
- Quinlan, G.D.; Tremaine, S. Symmetric multistep methods for the numerical integration of planetary orbits. *Astron. J.* **1990**, *100*, 1694–1700. [[CrossRef](#)]
- Lyche, T. Chebyshevian multistep methods for ordinary differential equations. *Numer. Math.* **1972**, *10*, 65–75. [[CrossRef](#)]
- Raptis, A.D.; Allison, A.C. Exponential-fitting methods for the numerical solution of the Schrödinger equation. *Comput. Phys. Commun.* **1978**, *14*, 1–5. [[CrossRef](#)]
- Konguetsof, A.; Simos, T.E. On the construction of Exponentially-Fitted Methods for the Numerical Solution of the Schrödinger Equation. *J. Comput. Math. Sci. Eng.* **2001**, *1*, 143–165. [[CrossRef](#)]
- Simos, T.E. *Atomic Structure Computations in Chemical Modelling: Applications and Theory*; Hinchliffe, A., Ed.; The Royal Society of Chemistry: London, UK, 2000; pp. 38–142.
- Simos, T.E.; Vigo-Aguiar, J. On the construction of efficient methods for second order IVPs with oscillating solution. *Int. J. Mod. Phys. C* **2001**, *10*, 1453–1476. [[CrossRef](#)]
- Dormand, J.R.; El-Mikkawy, M.E.A.; Prince, P.J. Families of Runge–Kutta–Nyström formulae. *IMA J. Numer. Anal.* **1987**, *7*, 235–250. [[CrossRef](#)]
- Van De Vyver, H. A Symplectic Exponentially Fitted Modified Runge–Kutta–Nyström Method for the Numerical Integration of Orbital Problems. *New Astron.* **2005**, *10*, 261–269. [[CrossRef](#)]
- Van De Vyver, H. On the Generation of P-Stable Exponentially Fitted Runge–Kutta–Nyström Methods By Exponentially Fitted Runge–Kutta Methods. *J. Comput. Appl. Math.* **2006**, *188*, 309–318. [[CrossRef](#)]
- Franco, J.M.; Khiar, Y.; Rández, L. Two new embedded pairs of explicit Runge–Kutta Methods adapted to the numerical solution of oscillatory problems. *Appl. Math. Comput.* **2015**, *252*, 45–57. [[CrossRef](#)]
- Franco, J.M.; Gomez, I. Symplectic explicit Methods of Runge–Kutta–Nyström type for solving perturbed oscillators. *J. Comput. Appl. Math.* **2014**, *260*, 482–493. [[CrossRef](#)]
- Franco, J.M.; Gomez, I. Some procedures for the construction of high-order exponentially fitted Runge–Kutta–Nyström Methods of explicit type. *Comput. Phys. Commun.* **2013**, *184*, 1310–1321. [[CrossRef](#)]
- Calvo, M.; Franco, J.M.; Montijano, J.I.; Rández, L. On some new low storage implementations of time advancing Runge–Kutta Methods. *J. Comput. Appl. Math.* **2012**, *236*, 3665–3675. [[CrossRef](#)]
- Calvo, M.; Franco, J.M.; Montijano, J.I.; Rández, L. Symmetric and symplectic exponentially fitted Runge–Kutta Methods of high order. *Comput. Phys. Commun.* **2010**, *181*, 2044–2056. [[CrossRef](#)]
- Calvo, M.; Franco, J.M.; Montijano, J.I.; Rández, L. On high order symmetric and symplectic trigonometrically fitted Runge–Kutta Methods with an even number of stages. *BIT Numer. Math.* **2010**, *50*, 3–21. [[CrossRef](#)]
- Franco, J.M.; Gomez, I. Accuracy and linear Stability of RKN Methods for solving second-order stiff problems. *Appl. Numer. Math.* **2009**, *59*, 959–975. [[CrossRef](#)]
- Calvo, M.; Franco, J.M.; Montijano, J.I.; Rández, L. Sixth-order symmetric and symplectic exponentially fitted Runge–Kutta Methods of the Gauss type. *J. Comput. Appl. Math.* **2009**, *22*, 387–398. [[CrossRef](#)]
- Calvo, M.; Franco, J.M.; Montijano, J.I.; Rández, L. Structure preservation of exponentially fitted Runge–Kutta Methods. *J. Comput. Appl. Math.* **2008**, *218*, 421–434. [[CrossRef](#)]
- Calvo, M.; Franco, J.M.; Montijano, J.I.; Rández, L. Sixth-order symmetric and symplectic exponentially fitted modified Runge–Kutta Methods of Gauss type. *Comput. Phys. Commun.* **2008**, *178*, 732–744. [[CrossRef](#)]

27. Franco, J.M. Exponentially fitted symplectic integrators of RKN type for solving oscillatory problems. *Comput. Phys. Commun.* **2007**, *177*, 479–492. [\[CrossRef\]](#)
28. Franco, J.M. New Methods for oscillatory systems based on ARKN Methods. *Appl. Numer. Math.* **2006**, *56*, 1040–1105. [\[CrossRef\]](#)
29. Franco, J.M. Runge–Kutta–Nyström Methods adapted to the numerical integration of perturbed oscillators. *Comput. Phys. Commun.* **2002**, *147*, 770–787. [\[CrossRef\]](#)
30. Franco, J.M. Stability of explicit ARKN Methods for perturbed oscillators. *J. Comput. Appl. Math.* **2005**, *173*, 389–396. [\[CrossRef\]](#)
31. Wu, X.Y.; You, X.; Li, J.Y. Note on derivation of order conditions for ARKN Methods for perturbed oscillators. *Comput. Phys. Commun.* **2009**, *180*, 1545–1549. [\[CrossRef\]](#)
32. Tocino, A.; Vigo-Aguiar, J. Symplectic conditions for exponential fitting Runge–Kutta–Nyström Methods. *Math. Comput. Model.* **2005**, *42*, 873–876. [\[CrossRef\]](#)
33. Van de Vyver, H. Comparison of some special optimized fourth-order Runge–Kutta Methods for the numerical solution of the Schrödinger equation. *Comput. Phys. Commun.* **2005**, *166*, 109–122. [\[CrossRef\]](#)
34. Van de Vyver, H. Frequency evaluation for exponentially fitted Runge–Kutta Methods. *J. Comput. Appl. Math.* **2005**, *184*, 442–463. [\[CrossRef\]](#)
35. Vigo-Aguiar, J.; Martín-Vaquero, J.; Ramos, H. Exponential fitting BDF–Runge–Kutta Algorithms. *Comput. Phys. Commun.* **2008**, *178*, 15–34. [\[CrossRef\]](#)
36. Demba, M.A.; Senu, N.; Ramos, H.; Kumam, P.; Watthayu, W. A Phase– and Amplification–Fitted 5(4) Diagonally Implicit Runge–Kutta–Nyström Pair for Oscillatory Systems. *Bull. Iran. Math. Soc.* **2023**, *49*, 24. [\[CrossRef\]](#)
37. Demba, M.A.; Ramos, H.; Watthayu, W.; Ahmed, I. A New Phase- and Amplification-Fitted Sixth-Order Explicit RKN Method to Solve Oscillating Systems. *Thai J. Math.* **2023**, *21*, 219–236.
38. Monovasilis, T.; Kalogiratou, Z. High Order Two-Derivative Runge–Kutta Methods with Optimized Dispersion and Dissipation Error. *Mathematics* **2021**, *9*, 232. [\[CrossRef\]](#)
39. Ahmad, N.A.; Senu, N.; Ibrahim, Z.B.; Othman, M.; Ismail, Z. Higher Order Three Derivative Runge–Kutta Method with Phase–Fitting and Amplification–Fitting Technique for Periodic IVPs. *Malaysian J. Math. Sci.* **2020**, *14*, 403–418.
40. Lee, K.; Alias, M.A.; Senu, N.; Ahmadian, A. On efficient frequency-dependent parameters of explicit two-derivative improved Runge–Kutta–Nyström method with application to two-body problem. *Alex. Eng. J.* **2023**, *72*, 605–620. [\[CrossRef\]](#)
41. Chien, L.K.; Senu, N.; Ahmadian, A.; Ibrahim, S.N.I. Efficient Frequency-Dependent Coefficients of Explicit Improved Two-Derivative Runge–Kutta Type Methods for Solving Third- Order IVPs. *Pertanika J. Sci. Technol.* **2023**, *31*, 843–873. [\[CrossRef\]](#)
42. Demba, M.A.; Ramos, H.; Kumam, P.; Watthayu, W.; Senu, N.; Ahmed, I. A trigonometrically adapted 6(4) explicit Runge–Kutta–Nyström pair to solve oscillating systems. *Math. Methods Appl. Sci.* **2023**, *46*, 560–578. [\[CrossRef\]](#)
43. Chen, B.Z.; Zhai, W.J. Optimal three-stage implicit exponentially-fitted RKN methods for solving second-order ODEs. *Calcolo* **2022**, *59*, 14. [\[CrossRef\]](#)
44. Senu, N.; Ahmad, N.A.; Othman, M.; Ibrahim, Z.B. Numerical study for periodical delay differential equations using Runge–Kutta with trigonometric interpolation. *Comput. Appl. Math.* **2022**, *41*, 25. [\[CrossRef\]](#)
45. Zhai, W.J.; Fu, S.H.; Zhou, T.C.; Xiu, C. Exponentially-fitted and trigonometrically-fitted implicit RKN methods for solving $y'' = f(t, y)$. *J. Appl. Math. Comput.* **2022**, *68*, 1449–1466. [\[CrossRef\]](#)
46. Senu, N.; Lee, K.C.; Ismail, W.F.W.; Ahmadian, A.; Ibrahim, S.N.I.; Laham, M. Improved Runge–Kutta Method with Trigonometrically-Fitting Technique for Solving Oscillatory Problem. *Malaysian J. Math. Sci.* **2021**, *15*, 253–266.
47. Fang, Y.L.; Yang, Y.P.; You, X. An explicit trigonometrically fitted Runge–Kutta method for stiff and oscillatory problems with two frequencies. *Int. J. Comput. Math.* **2020**, *97*, 85–94. [\[CrossRef\]](#)
48. Dormand, J.R.; Prince, P.J. A family of embedded Runge–Kutta formulae. *J. Comput. Appl. Math.* **1980**, *6*, 19–26. . [\[CrossRef\]](#)
49. Kalogiratou, Z.; Monovasilis, T.; Psihogios, G.; Simos, T.E. Runge–Kutta type methods with special properties for the numerical integration of ordinary differential equations. *Phys. Rep.* **2014**, *536*, 75–146. [\[CrossRef\]](#)
50. Shokri, A.; Khalsaraei, M.M. A new family of explicit linear two-step singularly P-stable Obrechkoff methods for the numerical solution of second-order IVPs. *Appl. Math. Comput.* **2020**, *376*, 125116. [\[CrossRef\]](#)
51. Abdulganiy, R.I.; Ramos, H.; Okunuga, S.A.; Majid, Z.A. A trigonometrically fitted intra-step block Falkner method for the direct integration of second-order delay differential equations with oscillatory solutions. *Afr. Mat.* **2023**, *34*, 36. [\[CrossRef\]](#)
52. Salih, M.M.; Ismail, F. Trigonometrically-Fitted Fifth Order Four-Step Predictor-Corrector Method for Solving Linear Ordinary Differential Equations with Oscillatory Solutions. *Malaysian J. Math. Sci.* **2022**, *16*, 739–748. [\[CrossRef\]](#)
53. Godwin, O.J.; Adewale, O.S.; Otuwatomi, O.P. An efficient block solver of trigonometrically fitted method for stiff odes. *Adv. Differ. Equ. Control Process.* **2022**, *28*, 73–98. [\[CrossRef\]](#)
54. Lee, K.C.; Senu, N.; Ahmadian, A.; Ibrahim, S.N.I. High-order exponentially fitted and trigonometrically fitted explicit two-derivative Runge–Kutta-type methods for solving third-order oscillatory problems. *Math. Sci.* **2022**, *16*, 281–297. [\[CrossRef\]](#)
55. Obaidat, S.; Butt, R. A new implicit symmetric method of sixth algebraic order with vanished phase-lag and its first derivative for solving Schrodinger's equation. *Open Math.* **2021**, *19*, 225–237. [\[CrossRef\]](#)
56. Shokri, A.; Neta, B.; Khalsaraei, M.M.; Rashidi, M.M.; Mohammad-Sedighi, H. A Singularly P-Stable Multi-Derivative Predictor Method for the Numerical Solution of Second-Order Ordinary Differential Equations. *Mathematics* **2021**, *9*, 806. [\[CrossRef\]](#)
57. Fang, Y.L.; Huang, T.; You, X.; Zheng, J.; Wang, B. Two-frequency trigonometrically-fitted and symmetric linear multi-step methods for second-order oscillators. *J. Comput. Appl. Math.* **2021**, *392*, 113312. [\[CrossRef\]](#)

58. Chun, C.; Neta, B. Trigonometrically-Fitted Methods: A Review. *Mathematics* **2019**, *7*, 1197. [[CrossRef](#)]
59. Anastassi, Z.A.; Simos, T.E. Numerical multistep methods for the efficient solution of quantum mechanics and related problems. *Phys. Rep.* **2009**, *482–483*, 1–240. [[CrossRef](#)]
60. Chawla, M.M.; Rao, P.S. A Noumerov-Type Method with Minimal Phase-Lag for the Integration of 2nd Order Periodic Initial-Value Problems. *J. Comput. Appl. Math.* **1984**, *11*, 277–281. [[CrossRef](#)]
61. Thomas, R.M. Phase properties of high order almost P-stable formulae. *BIT* **1984**, *24*, 225–238. [[CrossRef](#)]
62. Chawla, M.M.; Rao, P.S.; Neta, B. 2-Step 4Th-Order P-Stable Methods with Phase-Lag of Order 6 for $Y'' = F(T, Y)$. *J. Comput. Appl. Math.* **1986**, *16*, 233–236. [[CrossRef](#)]
63. Chawla, M.M.; Rao, P.S. An Explicit 6Th-Order Method with Phase-Lag of Order 8 for $Y'' = F(T, Y)$. *J. Comput. Appl. Math.* **1987**, *17*, 365–368. [[CrossRef](#)]
64. JColeman, J.P. Numerical-Methods for $Y'' = F(X, Y)$ Via Rational-Approximations for the Cosine. *IMA J. Numer. Anal.* **1989**, *9*, 145–165. [[CrossRef](#)]
65. Coleman, J.P.; Ixaru, L.G. P-Stability and Exponential-Fitting Methods for $Y'' = F(X, Y)$. *IMA J. Numer. Anal.* **1996**, *16*, 179–199. [[CrossRef](#)]
66. Coleman, J.P.; Duxbury, S.C. Mixed Collocation Methods for $Y'' = F(X, Y)$. *J. Comput. Appl. Math.* **2000**, *126*, 47–75. [[CrossRef](#)]
67. Ixaru, L.G.; Berceanu, S. Coleman Method Maximally Adapted to the Schrödinger-Equation. *Comput. Phys. Commun.* **1987**, *44*, 11–20. [[CrossRef](#)]
68. Ixaru, L.G.; Rizea, M. Numerov Method Maximally Adapted to the Schrödinger-Equation. *J. Comput. Phys.* **1987**, *73*, 306–324. [[CrossRef](#)]
69. Ixaru, L.G.; Vanden Berghe, G.; De Meyer, H.; Van Daele, M. Four-Step Exponential-Fitted Methods for Nonlinear Physical Problems. *Comput. Phys. Commun.* **1997**, *100*, 56–70. [[CrossRef](#)]
70. Ixaru, L.G.; Rizea, M. Four Step Methods for $Y'' = F(X, Y)$. *J. Comput. Appl. Math.* **1997**, *79*, 87–99. [[CrossRef](#)]
71. Van Daele, M.; Vanden Berghe, G.; De Meyer, H.; Ixaru, L.G. Exponential-Fitted Four-Step Methods for $Y'' = F(X, Y)$. *Int. J. Comput. Math.* **1998**, *66*, 299–309. [[CrossRef](#)]
72. Ixaru, L.G.; Paternoster, B. A Conditionally P-Stable Fourth-Order Exponential-Fitting Method for $Y'' = F(X, Y)$. *J. Comput. Appl. Math.* **1999**, *106*, 87–98. [[CrossRef](#)]
73. Ixaru, L.G. Numerical operations on oscillatory functions. *Comput. Chem.* **2001**, *25*, 39–53. [[CrossRef](#)] [[PubMed](#)]
74. Ixaru, L.G.; Vanden Berghe, G.; De Meyer, H. Exponentially Fitted Variable Two-Step BDF Algorithm for First Order Odes. *Comput. Phys. Commun.* **2003**, *150*, 116–128. [[CrossRef](#)]
75. Ixaru, L.G.; Rizea, M. Comparison of some four-Step Methods for the numerical solution of the Schrödinger equation. *Comput. Phys. Commun.* **1985**, *38*, 329–337. [[CrossRef](#)]
76. Ixaru, L.G.; Rizea, M. A Numerov-like scheme for the numerical solution of the Schrödinger equation in the deep continuum spectrum of energies. *Comput. Phys. Commun.* **1980**, *19*, 23–27. [[CrossRef](#)]
77. Avdelas, G.; Simos, T.E. A generator of high-order embedded P-stable method for the numerical solution of the Schrödinger equation. *J. Comput. Appl. Math.* **1996**, *72*, 345–358. [[CrossRef](#)]
78. Chawla, M.M.; Rao, P.S. An Noumerov-type Method with minimal phase-lag for the integration of second order periodic initial-value problems II Explicit Method. *J. Comput. Appl. Math.* **1986**, *15*, 329–337. [[CrossRef](#)]
79. Franco, J.M. Runge–Kutta methods adapted to the numerical integration of oscillatory problems. *Appl. Numer. Math.* **2004**, *50*, 427–443. [[CrossRef](#)]
80. Rizea, M. Exponential fitting Method for the time-dependent Schrödinger equation. *J. Math. Chem.* **2010**, *48*, 55–65. [[CrossRef](#)]
81. Ixaru, L.G.; Rizea, M.; Vanden Berghe, G.; De Meyer, H. Weights of the Exponential Fitting Multistep Algorithms for First-Order Odes. *J. Comput. Appl. Math.* **2001**, *132*, 83–93. [[CrossRef](#)]
82. Raptis, A.D.; Cash, J.R. Exponential and Bessel Fitting Methods for the Numerical-Solution of the Schrödinger-Equation. *Comput. Phys. Commun.* **1987**, *44*, 95–103. [[CrossRef](#)]
83. Simos, T.E. Predictor-corrector phase-fitted methods for $y''=f(x,y)$ and an application to the Schrödinger equation. *Int. J. Quantum Chem.* **1995**, *53*, 473–483. [[CrossRef](#)]
84. Raptis, A.D. Exponentially-Fitted Solutions of the Eigenvalue Shrödinger Equation with Automatic Error Control. *Comput. Phys. Commun.* **1983**, *28*, 427–431. [[CrossRef](#)]
85. Raptis, A.D. 2-Step Methods for the Numerical-Solution of the Schrödinger-Equation. *Comput. Phys. Commun.* **1982**, *28*, 373–378. [[CrossRef](#)]
86. Raptis, A.D. On the Numerical-Solution of the Schrödinger-Equation. *Comput. Phys. Commun.* **1981**, *24*, 1–4. [[CrossRef](#)]
87. Raptis, A.D. Exponential-Fitting Methods for the Numerical-Integration of the 4Th-Order Differential-Equation $Y^{iv} + F.Y = G$. *Computing* **1980**, *24*, 241–250. [[CrossRef](#)]
88. Van Daele, M.; Vanden Berghe, G. P-stable exponentially-fitted Obrechkoff Methods of arbitrary order for second-order differential equations. *Numer. Algorithms* **2007**, *46*, 333–350. [[CrossRef](#)]
89. Fang, Y.; Wu, X. A Trigonometrically Fitted Explicit Numerov-Type Method for Second-Order Initial Value Problems with Oscillating Solutions. *Appl. Numer. Math.* **2008**, *58*, 341–351. [[CrossRef](#)]
90. Vanden Berghe, G.; Van Daele, M. Exponentially-fitted Obrechkoff Methods for second-order differential equations. *Appl. Numer. Math.* **2009**, *59*, 815–829. [[CrossRef](#)]

91. Hollevoet, D.; Van Daele, M.; Vanden Berghe, G. the Optimal Exponentially-Fitted Numerov Method for Solving Two-Point Boundary Value Problems. *J. Comput. Appl. Math.* **2009**, *230*, 260–269. [[CrossRef](#)]
92. Franco, J.M.; Rández, L. Explicit exponentially fitted two-Step hybrid Methods of high order for second-order oscillatory IVPs. *Appl. Math. Comput.* **2016**, *273*, 493–505. [[CrossRef](#)]
93. Franco, J.M.; Gomez, I.; Rández, L. Optimization of explicit two-Step hybrid Methods for solving orbital and oscillatory problems. *Comput. Phys. Commun.* **2014**, *185*, 2527–2537. [[CrossRef](#)]
94. Franco, J.M.; Gomez, I. Trigonometrically fitted nonlinear two-Step Methods for solving second order oscillatory IVPs. *Appl. Math. Comput.* **2014**, *232*, 643–657. [[CrossRef](#)]
95. Konguetsof, A. A generator of families of two-Step numerical Methods with free parameters and minimal phase-lag. *J. Math. Chem.* **2017**, *55*, 1808–1832. [[CrossRef](#)]
96. Konguetsof, A. A hybrid Method with phase-lag and derivatives equal to zero for the numerical integration of the Schrödinger equation. *J. Math. Chem.* **2011**, *49*, 1330–1356. [[CrossRef](#)]
97. Van de Vyver, H. A phase-fitted and amplification-fitted explicit two-Step hybrid Method for second-order periodic initial value problems. *Int. J. Mod. Phys. C* **2006**, *17*, 663–675. [[CrossRef](#)]
98. Van de Vyver, H. An explicit Numerov-type Method for second-order differential equations with oscillating solutions. *Comput. Math. Appl.* **2007**, *53*, 1339–1348. [[CrossRef](#)]
99. Fang, Y.; Wu, X. A trigonometrically fitted explicit hybrid Method for the numerical integration of orbital problems. *Appl. Math. Comput.* **2007**, *189*, 178–185. [[CrossRef](#)]
100. Van de Vyver, H. Phase-fitted and amplification-fitted two-Step hybrid Methods for $y'' = f(x, y)$. *J. Comput. Appl. Math.* **2007**, *209*, 33–53. [[CrossRef](#)]
101. Van de Vyver, H. Efficient one-Step Methods for the Schrödinger equation. *MATCH-Commun. Math. Comput. Chem.* **2008**, *60*, 711–732.
102. Martín-Vaquero, J.; Vigo-Aguiar, J. Exponential fitted Gauss, Radau and Lobatto Methods of low order. *Numer. Algorithms* **2008**, *48*, 327–346.
103. Konguetsof, A. A new two-Step hybrid Method for the numerical solution of the Schrödinger equation. *J. Math. Chem.* **2010**, *47*, 871–890. [[CrossRef](#)]
104. Hendi, F.A. P-Stable Higher Derivative Methods with Minimal Phase-Lag for Solving Second Order Differential Equations. *J. Appl. Math.* **2011**, *2011*, 407151. [[CrossRef](#)]
105. Wang, Z.; Zhao, D.; Dai, Y.; Wu, D. An improved trigonometrically fitted P-stable Obrechkoff Method for periodic initial-value problems. *Proc. R. Soc. A-Math. Phys. Eng. Sci.* **2005**, *461*, 1639–1658. [[CrossRef](#)]
106. Van Daele, M.; Vanden Berghe, G.; De Meyer, H. Properties and Implementation of R-Adams Methods Based On Mixed-Type Interpolation. *Comput. Phys. Commun.* **1995**, *30*, 37–54. [[CrossRef](#)]
107. Wang, Z. Trigonometrically-fitted Method with the Fourier frequency spectrum for undamped duffing equation. *Comput. Phys. Commun.* **2006**, *174*, 109–118. [[CrossRef](#)]
108. Wang, Z. Trigonometrically-fitted Method for a periodic initial value problem with two frequencies. *Comput. Phys. Commun.* **2006**, *175*, 241–249. [[CrossRef](#)]
109. Tang, C.; Yan, H.; Zhang, H.; Li, W. The various order explicit multistep exponential fitting for systems of ordinary differential equations. *J. Comput. Appl. Math.* **2004**, *169*, 171–182. [[CrossRef](#)]
110. Tang, C.; Yan, H.; Zhang, H.; Chen, Z.; Liu, M.; Zhang, G. The arbitrary order implicit multistep schemes of exponential fitting and their applications. *J. Comput. Appl. Math.* **2005**, *173*, 155–168. [[CrossRef](#)]
111. Coleman, J.P.; Ixaru, L.G. Truncation Errors in exponential fitting for oscillatory problems. *SIAM J. Numer. Anal.* **2006**, *44*, 1441–1465. [[CrossRef](#)]
112. Paternoster, B. Present state-of-the-art in exponential fitting. A contribution dedicated to Liviu Ixaru on their 70th birthday. *Comput. Phys. Commun.* **2012**, *183*, 2499–2512. [[CrossRef](#)]
113. Wang, Z. Obrechkoff one-Step Method fitted with Fourier spectrum for undamped Duffing equation. *Comput. Phys. Commun.* **2006**, *175*, 692–699. [[CrossRef](#)]
114. Wang, C.; Wang, Z. A P-stable eighteenth-order six-Step Method for periodic initial value problems. *Int. J. Mod. Phys. C* **2007**, *18*, 419–431. [[CrossRef](#)]
115. Simos, T.E. A New Methodology for the Development of Efficient Multistep Methods for First-Order IVPs with Oscillating Solutions. *Mathematics* **2024**, *12*, 504. [[CrossRef](#)]
116. Stiefel, E.; Bettis, D.G. Stabilization of Cowell’s method. *Numer. Math.* **1969**, *13*, 154–175. [[CrossRef](#)]
117. Fehlberg, E. Classical fifth-, Sixth-, Seventh-, and Eighth-Order Runge–Kutta Formulas with Step-size Control. NASA Technical Report 287. 1968. Available online: <https://ntrs.nasa.gov/api/citations/19680027281/downloads/19680027281.pdf> (accessed on 30 March 2024)
118. Cash, J.R.; Karp, A.H. A variable order Runge–Kutta method for initial value problems with rapidly varying right-hand sides. *ACM Trans. Math. Softw.* **1990**, *16*, 201–222. [[CrossRef](#)]
119. Franco, J.; Gómez, I.; Rández, L. Four-stage symplectic and P-stable SDIRKN methods with dispersion of high order. *Numer. Algorithms* **2001**, *26*, 347–363. [[CrossRef](#)]
120. Franco, J.M.; Palacios, M. High-order P-stable multistep methods. *J. Comput. Appl. Math.* **1990**, *30*, 1–10. [[CrossRef](#)]

121. Simos, T.E. New Open Modified Newton Cotes Type Formulae as Multilayer Symplectic Integrators. *Appl. Math. Modell.* **2013**, *37*, 1983–1991. [[CrossRef](#)]
122. Petzold, L.R. An efficient numerical method for highly oscillatory ordinary differential equations. *SIAM J. Numer. Anal.* **1981**, *18*, 455–479. [[CrossRef](#)]
123. Ramos, H.; Vigo-Aguiar, J. On the frequency choice in trigonometrically fitted methods. *Appl. Math. Lett.* **2010**, *11*, 1378–1381. [[CrossRef](#)]
124. Ixaru, L.G.; Vanden Berghe, G.; De Meyer, H. Frequency evaluation in exponential fitting multistep algorithms for ODEs. *J. Comput. Appl. Math.* **2002**, *140*, 423–434. [[CrossRef](#)]
125. Boyce, W.E.; DiPrima, R.C.; Meade, D.B. *Elementary Differential Equations and Boundary Value Problems*, 11th ed.; John Wiley & Sons: Hoboken, NJ, USA, 2017.
126. Lawrence, C. *Evans, Partial Differential Equations*, 2nd ed.; American Mathematical Society: Providence, RI, USA, 2010; Chapter 3, p. 91135.

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