

## Supplementary Videos

### Supplementary Video S1

Simulated current-induced motion of isolated skyrmions and their nucleation in wedge-shaped samples of cubic helimagnets with the elevation angle  $\alpha = 10^\circ$ . The movie frames are taken with the interval  $\Delta t = 1 \text{ ns}$ . Skyrmions nucleated at the sharp boundary also “climb” the hill. The current density,  $j_y = 1 \times 10^{13} \text{ J/m}^2$ .

### Supplementary Video S2

Skyrmion annihilation at the sharp boundary of the wedge for  $\alpha = 25^\circ$ . Since the energy minimum of the edge-skyrmion interaction ceases to exist for such large angles, the edge states do not hamper the skyrmion collapse. The current is zero.

### Supplementary Video S3

Return of isolated skyrmions into the energy minimum of the edge-skyrmion interaction near the sharp boundary after the current is switched off.  $\alpha = 15^\circ$ . In this case, the edge states are able to keep the skyrmions within the nanostructure.

### Supplementary Video S4

An isolated skyrmion approaches the edge of a racetrack memory with thin-film geometry. The SHE in this case is balanced by the edge-skyrmion repulsion, which leads to the straight skyrmion trajectory.  $j = 5 \times 10^{12} \text{ J/m}^2$ , the thickness of the layer is 0.5. Notice that according to Fig. 2 in Ref. [54], the skyrmion first moves with variable velocity until it reaches a constant value near the edge.

### Supplementary Video S5

For  $j = 1 \times 10^{13} \text{ J/m}^2$ , the skyrmion annihilates at the boundary of the racetrack since the repulsive force from the edge state is unable to keep the skyrmion within the sample.