

Article



# **Performance Improvement of a Limaçon Gas Expander Using an Inlet Control Valve: Two Case Studies**

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Abstract: Renewable energy-based compact energy-generation systems based on the organic Rankine cycle (ORC) can be employed to meet the ever-growing thirst for affordable and clean energy. The overall performance and effectiveness of ORC systems are constrained by the low efficiency of the gas expander, specifically the positive displacement expander, which is responsible for energy conversion from the working fluid. This low-efficiency scenario can be significantly improved by employing a control valve to regulate and restrict the flow of the working fluid into the expander. A control valve can effectively curve the loss of costly compressed and energized working fluids by allowing them to expand in the expander chamber before discharging through the outlet port. They can thus be used to regulate the amount of energy yield and output power. In this work, two direct drive rotary valves (DDRVs) operated by a stepper motor (SM-DDRV) and rotary solenoid (RS-DDRV) are suggested, and the behavior of the valves is examined. The effect of friction and temperature on the valve response is also studied. Additionally, the effect of inlet control valves on the overall system performance of the limaçon expander is assessed. Thermodynamic properties such as the isentropic efficiency and filling factor are also computed. The effect of leakage due to valve response delay is analyzed at different inlet pressures. The performance indices are compared to the expander performance without any inlet valve. The SM-DDRV setup results in a 14.86% increase in isentropic efficiency and a 220% increase in the filling factor, whereas the RS-DDRV performs moderately with a 2.58% increase in isentropic efficiency and an 80% increase in the filling factor compared to a ported expander. The SM-DDRV provides better performance indices compared to the RS-DDRV and without valve setups. However, the performance of the limaçon expander with the SM-DDRV is sensitive to the inlet pressure and degrades at higher pressure. Overall, the valves proposed in this work present key insights into improving the performance characteristics of gas expanders of ORC systems.

**Keywords:** limaçon gas expander; rotary machine; energy conversion; direct drive rotary valve; control valve; system efficiency; mathematical modeling; system response

# 1. Introduction

In this technologically advanced 21st century, we have to endure the rising energy crisis on a global scale. This energy crisis has a twofold aspect. On the one hand, the ever-increased large-scale production of energy in developed countries has escalated the greenhouse effect; on the other hand, the lack of energy and insufficient development in the energy sector in the underdeveloped and isolated sections of the world has widened the gap between these groups. Aiming for net zero emissions while, at the same time, ensuring affordable energy for all will be one of the biggest challenges that we need to address now and in the coming future.

High carbon emissions due to energy production using fossil fuels are primarily contributing to climate change and global warming. The low efficiency of energy conversion



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**Copyright:** © 2024 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). processes largely used currently is mainly contributing to this issue. Currently, traditional power-generation methods using fossil fuels have system efficiencies ranging from 30% to 40%. This limitation is attributed primarily to the sub-optimal performance of turbines, boilers, and condensers [1]. However, the adoption of a combined cycle configuration can raise the efficiency to 55% [1]. Nonetheless, approximately 45% of the usable energy that is wasted in the form of heat and remains unrealized could potentially be converted into electricity through appropriate means.

The utilization of low-grade or waste heat can be one approach to curtail energy loss and improve efficiency. Waste heat recovery (WHR) systems can be applied for this purpose where the low-grade sources can be transformed into usable power using the ORC process. Traditional ORCs consist of the following four basic components; expander, condenser, pump, and evaporator, as shown in Figure 1. The ORC cycle starts in the evaporator by heating the organic working fluid in an isobaric process. The heated fluid/vapor is then passed through the expander, where it expands in an isentropic process. During this expansion process, the thermal potential energy converts to mechanical energy and rotates a rotor inside the expander chamber. The expanded fluid is later condensed in the condenser in an isobaric process. Finally, the compressor compresses the fluid to its normal state in an isentropic process, and the cycle continues.

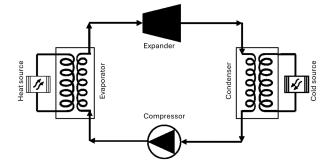


Figure 1. Basic arrangement of an ORC process.

The physical and thermal characteristics of the fluid play a crucial role in the overall performance and efficiency of an ORC system [2]. In the ORC, water, the traditional working fluid in the conventional Rankine cycle, is replaced by organic refrigerants having lower boiling temperatures and pressures [3,4]. Therefore, less heat is required to obtain the compressed vapors for these fluids compared to water [5]. This modification permits the ORC system to operate with low-grade heat. Various working fluids have been used such as R2341a [6], R245fa [7], pentane [8], hexamethyldisloxane [9], R134 [10], cyclopentane [11], and n-heptane [12]. In this study, R245fa is selected as the primary working fluid as it is safer for the atmosphere having a low global warming potential (GWP) of 1030 and a zero ozone depletion potential (ODP) [13].

Accordingly, the ORC–WHR combined systems can be employed in renewable power generators, such as solar or bio-thermal plants, which produce heat in the range of 90–500 °C [14]. ORC–WHR systems can function within that low-temperature range. Operating temperatures as low as 73.3 °C have been reported, presenting a significant advantage for improving the efficiency of various thermal cycles [15,16]. Moreover, ORC–WHR systems prove to be more economical when compared to other similar technologies. For instance, they exhibit approximately 75% and 200% lower investment, operation, and maintenance costs than conventional gasification systems [17]. Recently, numerous ORC-based small-scale systems have been developed, which are suited for domestic applications. These are perfect for residential use as they operate at lower and safer operating temperatures and pressures. For instance, a 1 kWe ORC-based off-grid solar thermal power plant has been built in rural Lesotho, which produced cleaner affordable energy compared to diesel engines (<\$0.18/kWh) [18]. Furthermore, these systems can be used as combined heat and power (CHP) systems, providing both electricity and heat using various renewable energy

sources [19–21]. However, the poor efficiency of expanders currently constrains their more widespread use.

In general, the efficiency of ORC systems typically ranges between 9.6% and 18.1%, contingent on various aspects such as the type of fluids, the expander efficiency, and the operating temperatures [22]. However, the ORC system's performance is largely influenced by the expander's isentropic efficiency. In recent times, efforts have been made towards the improvement of expander design and operation for better efficiency. A limaçon expander, which has largely been overlooked in the past due to the technological challenges associated with manufacturing the machine housing and rotor, has recently experienced a revival in the scientific community. For instance, Sultan et al. proposed a method of power generation using heat recovery by employing such expanders [23,24]. However, advances in manufacturing, control, and optimization technologies have now made it more feasible to realize such small-scale power units. Sultan's work also demonstrated that using a camoperated inlet control valve could potentially boost the expander's overall performance by at least 24%. Nevertheless, cam-operated valves have fixed characteristics and lack the flexibility to adapt to varying loads or operating conditions.

In this paper, we propose replacing the cam-operated valve with a control valve. Two different valves actuated by a stepper motor and rotary solenoid, respectively, have been introduced. Both actuators offer a flexible configuration and adaptive controllability, allowing programmable valve operation at specified expander rotor angles. Here, we will assess the performance of both the valves at the inlet of the Limaçon expander to explore their potential in WHR applications. The selection of two different actuators is motivated by their high precision controllability and affordability. A stepper motor is a DC synchronous motor producing discrete precise movements rather than continuous rotation. In contrast, a rotary solenoid is a modification of the linear solenoid, which can produce a fixed angular rotation. Both of the actuators can be operated by applying a sequence of electric pulses producing rotational motion in either the clockwise or anti-clockwise direction. The amplitude and sequence of the input control signal can be controlled to control the movement and speed of the actuators. For this work, a hybrid stepper motor is chosen as they offer comparatively better performance [25]. Both actuators are modeled mathematically and simulated in conjunction with a thermodynamic model of the limaçon gas expander. This study provides a theoretical basis for the usability of control valves for better performance and controllability of a limaçon gas expander and, consequently, the overall ORC system. The study can be progressed further in the future with necessary experiments and the possible implementation of ORC plants with limaçon expanders and the proposed valves.

# 2. Outline of the Proposed Valve

Figure 2 shows the proposed direct drive rotary valve. The primary part of the valve is the spool, which is actuated by either stepper motors or rotary solenoids. The valve is fit inside the inlet port of the Limaçon expander. The valve is designed in such a way that it is considered normal open (NO) in the absence of any actuation voltage. This is accomplished with the aid of mechanical springs holding the valve at an open position against the differential fluid pressure. This is to ensure that the expander remains operational even in the case of any mechanical or electrical issues disrupting the valve operation.

At the initial NO position of  $\theta = 0$ , the working fluid at a pressure  $P_{in}$  (kPa) enters the expander chamber through the valve antechamber, which is at pressure  $P_a$  (kPa), as shown in Figure 2a. During closing operation, the actuator generates an electromagnetic torque,  $\tau_e$  (Nm), inducing a clockwise rotation,  $\theta$  (rad), of the valve spool with an angular velocity,  $\omega$  (rad/s), as shown in Figure 2b. Two mechanical stoppers can be placed at opposite ends at  $\theta = 0$  and  $\theta = \theta_{max}$  to ensure stoppage.

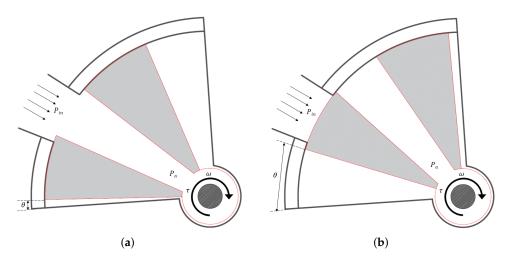


Figure 2. DDRV outline. (a) Open and (b) closed positions.

# 3. Mathematical Modeling

In the following subsections, comprehensive mathematical models of both the stepper motor and rotary solenoid valves will be developed. A thermodynamic model of the limaçon to circular gas expander will also be provided.

#### 3.1. Stepper Motor

Stepper motors are a type of electromechanical transmission device capable of producing high-torque specific and incremental motion in reflection with respect to the digital input electric pulses [26]. The input pulses dictate the direction, angle, and speed of rotation. The stepper motors can be classified into three types according to the construction: variable reluctance, permanent magnet, and hybrid [27]. The variable reluctance type has a toothed rotor and stator, whereas in the permanent magnet type, the toothed rotor is replaced by a permanent magnet [27,28]. The hybrid type is a combination of these two, where the rotor is a toothed-type magnet providing higher torque with smaller steps [28,29]. As such, a hybrid stepper motor should be a good candidate as a valve actuator. Now, the hybrid stepper motor will be modeled mathematically. Matsui, Nakamura, and Kosaka pointed out that the electrical domain of a stepper motor can be represented as a simple R-L circuit and accounting for the rate of change of flux linkages as shown in Figure 3a,b [30]. In this study, a two-phase stepper motor is used, whose voltage–current characteristics in the phases *a* and *b* can be represented as follows:

$$v_a = R_a i_a + \frac{d}{dt} (L_a i_a + \psi_a) \tag{1a}$$

$$v_b = R_b i_b + \frac{d}{dt} (L_b i_b + \psi_b) \tag{1b}$$

where v (volt) is the input voltage, i (amp) is the resulting current through the windings, and R (Ohm), L (H), and  $\psi$  (H) are the winding resistance, inductance, and mutual flux linkage, respectively. Substituting the expressions for mutual flux linkages  $\psi_a$  and  $\psi_b$  provided by [31] into Equations (1a) and (1b) gives:

$$v_a = R_m i_a + L_m \frac{di_a}{dt} - p \psi_m \omega sin(p\theta)$$
(2a)

$$v_b = R_m i_b + L_m \frac{di_b}{dt} + p \psi_m \omega_m cos(p\theta)$$
<sup>(2b)</sup>

where  $R_m = R_a = R_b$  is the phase resistance,  $L_m = L_a = L_b$  is the phase self-inductance, p is the number of rotor pole pairs, and  $\psi_m$  is the maximum flux linkage. The voltage and

currents in the orthogonal static axis a - b can be transformed into a rotating d - q axis using the Park transformation as given in Equation (3) for the ease of analysis [32]. The equivalent circuit for the d - q rotary axis is shown in Figure 4.

$$\begin{bmatrix} v_d \\ v_q \end{bmatrix} = T_p \begin{bmatrix} v_a \\ v_b \end{bmatrix} and \begin{bmatrix} i_d \\ i_q \end{bmatrix} = T_p \begin{bmatrix} i_a \\ i_b \end{bmatrix}$$
(3)

where  $[T_p] = \begin{bmatrix} cos(p\theta) & sin(p\theta) \\ -sin(p\theta) & cos(p\theta) \end{bmatrix}$  is the Park transformation matrix. Performing this transformation on Equations (2a) and (2b) gives the following expression:

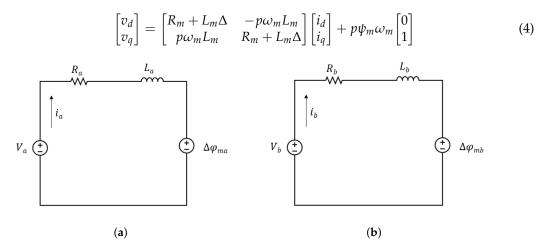


Figure 3. Stepper motor equivalent circuit (a) phase a and (b) phase b.

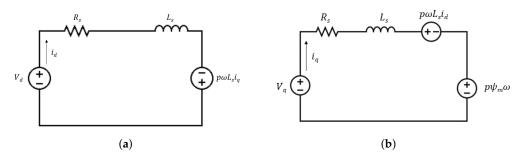


Figure 4. Stepper motor equivalent circuit (a) d axis and (b) q axis.

The expressions for the electrical power input to the motor can be utilized to obtain the electromagnetic torque as shown below:

$$Pow_{elec} = v_d i_d + v_q i_q$$

$$= R_m (i_d^2 + i_q^2) + \frac{1}{2} L_m \frac{d}{dt} (i_d^2 + i_q^2) + p \psi_m \omega_m i_q$$

$$= Pow_{cu} + Pow_{ind} + Pow_{mech}$$
(5)

where  $Pow_{cu}$  is the motor copper loss wasted as heat of the conductors,  $Pow_{ind}$  is the magnetic stored energy, and  $Pow_{mech}$  is the mechanical power component producing torque. The electromagnetic torque can, thus, be derived from  $Pow_{mech}$  as:

$$\tau_m = \frac{Pow_{mech}}{\omega_m} = p\psi_m i_q \tag{6}$$

# 3.2. Rotary Solenoid

The rotary solenoid actuator is an electromechanical device comparable to a common linear solenoid actuator, where the former generates torque, unlike the linear push–pull

force of the latter. In this manuscript, two rotary solenoids, as shown in Figure 5, are used in tandem to operate the valve in both directions. The construction of the rotary solenoid is shown in Figure 6. The linear motion of the plunger is converted to rotary motion by using a ball bearing. The movement of the ball along the races during actuation is shown in Figure 6b. The races are designed in such a way as to produce rapid motion at the start of the stroke and eventually slow down at the stroke end. This is performed to balance the force along the stroke as the downward force pulling the plunger is smaller at the start due to the higher air gap and increases as the air gap becomes smaller at the end of the stroke. When the solenoid coil is energized, a flow of current through it produces a magnetic field around the coil, which creates an induction action in the armature or plunger, producing force.

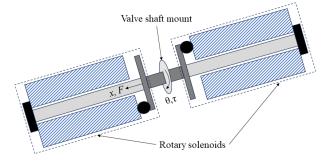
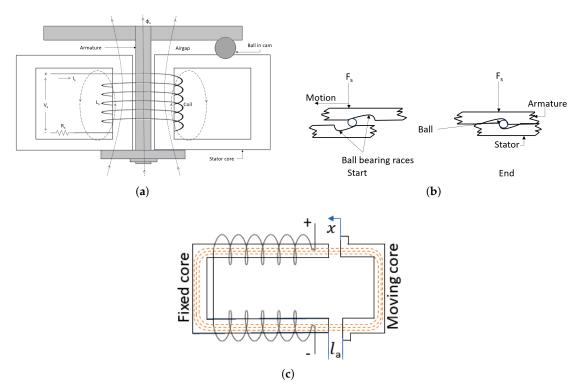


Figure 5. Arrangement of solenoids.



**Figure 6.** Rotary solenoid. (**a**) Schematic diagram of rotary solenoid, (**b**) armature rotary action, and (**c**) solenoid magnetic circuit.

Considering the solenoid magnetic circuit of Figure 6c with  $N_s$  number of turns, the electromotive force due to the flow of current  $i_s$  can be written as:

$$N_s i_s = H_c l_c + 2H_a l_a \approx H_c l_e \tag{7}$$

where  $H_c$  and  $H_a$  are the magnetic field intensity of the core and air gap, respectively;  $l_c$  and  $l_a$  are the core and air gap length, respectively. As the voltage is applied, the moving

$$L_{s} = \frac{N_{s}\phi}{i_{s}} = \frac{N_{s}^{2}\phi}{H_{c}l_{e}} = \frac{N_{s}^{2}\mu_{c}A_{e}H_{c}}{H_{c}l_{e}} = \frac{N_{s}^{2}\mu_{c}A_{e}}{l_{e}}$$

$$\Rightarrow L_{s} = \frac{N_{s}^{2}\mu_{c}A_{e}}{x_{s} + 2\mu_{r}(S - x_{s})}$$
(8)

The rate of change of inductance is given by:

$$\dot{L_s} = \frac{2\mu_r N_s^2 \mu_c A_e \dot{x}_s}{[l_c + 2\mu_r (l_a - x_s)]^2} = \frac{(1 - 2\mu_r) L_s^2 \dot{x}_s}{N_s^2 \mu_c A_e}$$
(9)

The voltage–current relationship can be represented as follows:

$$v_{s} = R_{s}i_{s} + L_{s}i_{s} + i_{s}L_{s}$$

$$= R_{s}i_{s} + L_{s}\dot{i}_{s} + i_{s}\frac{2\mu_{r}L_{s}^{2}\dot{x}_{s}}{N_{s}^{2}\mu_{c}A_{e}}$$

$$\Rightarrow \dot{i}_{s} = \frac{1}{L_{s}}[v_{s} - R_{s}i_{s}] - i_{s}\frac{(1 - 2\mu_{r})L_{s}\dot{x}_{s}}{N_{s}^{2}\mu_{c}A_{e}}$$
(10)

The axial force  $F_s$  and resultant torque  $\tau_s$  is given by Bolden [33] as follows:

$$F_s = \frac{1}{2\mu_0 A_e} sin^2(b) (\frac{L_s i_s}{N_s})^2$$
(11)

$$\tau_{s} = \frac{1}{4\mu_{0}A_{e}}r_{s}sin(2b)(\frac{L_{s}i_{s}}{N_{s}})^{2}$$
(12)

where *b* is the angle traversed by the ball bearing along the race of radius  $r_s$ . *b* is given by [33] as  $arctan(1.07 \times 10^{-3} x_s^{-0.7})$ .

## 3.3. Valve Dynamics

The value is a rotating mechanical system, which can be described by the torquefriction relations according to Newton's second law of rotation as follows:

$$J\dot{\omega} = \tau_{ext} - \tau_{fric} - \tau_f - \tau_l \tag{13}$$

where  $\tau_{ext}$  (Nm) is the external electromagnetic torque generated by the actuators denoted by  $\tau_m$  and  $\tau_s$  for the motor and solenoid, respectively, J (kg · m<sup>2</sup>) is the total system inertia including the actuator and spool,  $\tau_{fric}$  (Nm) is the frictional torque,  $\tau_f$  (Nm) is the torque corresponding to steady state fluid flow, and  $\tau_l$  (Nm) is the torque due to connected load. For the rotary solenoid valve, another equation should be added to account for the linear motion of the plunger as below:

$$F_s + F_p = M\ddot{x_s} + B\dot{x_s} + kx_s \tag{14}$$

where *M* is the total mass of the solenoid actuator and associated valve,  $F_p$  is the differential pressure force at the valve spool, and *k* is the spring constant (N/m) The flow torque  $\tau_f$  is due to the flow of fluid through the valve port. When fluid passes through the valve port, it introduces a flow force on the rotating valve spool. The axial component of this force

produces a torque on the spool, which obstructs the movement of the spool. The expression of flow torque is given by Okhotnikov et al. [34] as below:

$$\tau_f = 2C_d C_v (P_{in} - P_a) A(\theta) rsin\alpha_j \tag{15}$$

where  $C_d$  is the discharge coefficient,  $C_v$  is the discharge velocity coefficient,  $A(\theta)$  is the dynamic area of the orifice, r is the radius of the rotating spool from the center of rotation, and  $\alpha_j$  is the angle between the radial flow direction and the radius of the rotating spool denoted as the jet angle. Figure 7 shows the variation of the passage area,  $A(\theta)$ , which can be expressed as a function of  $\theta$  as shown below:

$$A(\theta) = D_{valve} D_{orifice} (\theta_{max} - \theta)$$
(16)

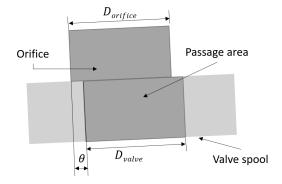


Figure 7. Dynamic passage area.

Substituting (16) into (15) gives the expression of flow torque as follows:

$$\tau_f = 2rC_d C_v D_{valve} D_{orifice} \Delta \theta (P_{in} - P_a) sin\alpha_j \tag{17}$$

where  $\Delta \theta = (\theta_{max} - \theta)$ . The flow rate through the orifice as given by Watton [35] is shown below:

$$q_l = C_d A(\theta) \sqrt{\frac{1}{\rho} (P_{in} - \frac{\theta}{|\theta|} P_a)}$$
(18)

Due to factors such as delay in the current buildup, inertia, and friction, the closing or opening of the valve will never be instantaneous. This delay in response will constitute a leakage flow through the orifice. This leakage is significant during the closing operation as the valve is intended to block all fluid flow into the expander. The leakage can be calculated as follows:

$$q_{leak} = \int_{t_1}^{t_2} q_l dt \tag{19}$$

where  $t_1$  (s) is the time at which the input voltage to the actuators is given and  $t_2$  (s) is when the valve has finished its closing or opening cycle. Apart from the delay, different friction losses will also contribute to the overall energy conversion. The total energy loss per expander cycle due to frictions in the valve rotating parts can be calculated as follows:

$$F_{loss} = \int_0^{t_c} (\tau_{fric} + \tau_f) dt \tag{20}$$

where  $t_c$  (s) is the time in a cycle. Besides friction, the operating temperature plays a vital role in the performance of the actuators. A rise in temperature will result in a rise in the coil resistance of both actuators, which can be calculated using the following well-known equation:

$$R_e = R_i (1 + \alpha \Delta T) \tag{21}$$

where  $R_i$  and  $R_e$  are the initial and effective resistance in  $\Omega$ ,  $\alpha$  (\°C) is the temperature coefficient of resistivity, and  $\Delta T$  (°C) is the rise in temperature. The increase in resistivity of the copper coil will reduce the current flowing through it and, consequently, reduce the angular displacement.

#### 3.4. System Equations

In order to analyze the stepper motor-operated valve, Equations (4), (6), and (13) are reformulated as follows:

$$v_{d} = v_{a}cos(p\theta) + v_{b}sin(p\theta)$$

$$v_{q} = -v_{a}sin(p\theta) + v_{b}cos(p\theta)$$

$$i_{d} = \frac{1}{L_{m}}(v_{d} - R_{m}i_{d} + p\omega L_{m}i_{q})$$

$$i_{q} = \frac{1}{L_{m}}(v_{q} - R_{m}i_{q} - p\omega L_{m}i_{d} - p\psi_{m}\omega)$$

$$\tau_{m} = p\psi_{m}\omega_{m}$$

$$\theta_{m}^{i} = \omega_{m}$$

$$\omega_{m} = \frac{1}{I}(\tau_{m} - B\omega_{m} - \tau_{f} - \tau_{I})$$
(22)

The set of equations in (22) can be solved numerically to simulate the operation of the stepper motor valve. Similarly, the set of equations in (23) can be solved for a rotary solenoid valve.

$$\begin{split} \dot{L}_{s} &= \frac{(1-2\mu_{r})L_{s}^{2}\dot{x}_{s}}{N_{s}^{2}\mu_{c}A_{e}} \\ \dot{i}_{s} &= \frac{1}{L_{s}}[V_{s} - R_{s}i_{s}] - i_{s}\frac{(1-2\mu_{r})L_{s}\dot{x}_{s}}{N_{s}^{2}\mu_{c}A_{e}} \\ F_{s} &= \frac{1}{2\mu_{0}A_{e}}sin^{2}(b)(\frac{L_{s}i_{s}}{N_{s}})^{2} \\ \tau_{s} &= \frac{1}{4\mu_{0}A_{e}}rsin(2b)(\frac{L_{s}i_{s}}{N_{s}})^{2} \\ \ddot{x}_{s} &= \frac{1}{M}(F_{s} + F_{p} - B\dot{x}_{s} - kx_{s}) \\ \dot{\theta}_{s} &= \omega_{s} \\ \dot{\omega}_{s} &= \frac{1}{I}(\tau_{s} - B\omega_{s} - \tau_{f} - \tau_{I}) \end{split}$$
(23)

# 3.5. Limaçon Positive Displacement Gas Expander

A positive displacement gas expander is an energy converter that converts the potential energy of a compressed fluid into usable mechanical energy by allowing the fluid to expand in an enclosed chamber. They are specially designed to effectively function with a low shaft speed and low-flow-rate working fluids [23]. This is particularly useful in implementing ORC systems with low-grade heat sources. These expanders can be broadly classified as rotary, piston, and reciprocating according to their construction [36]. Limaçon gas expanders are a type of rotary positive displacement expanders where the machines' housing and rotor are constructed based on the Limaçon curves of the form  $u = b + a \cos \phi$  or  $u = a + b \sin \phi$  in polar coordinates  $(u, \phi)$ . Figure 8 shows a limaçon expander consisting of a dual-lobe rotor, housing, inlet–outlet ports, and an inlet control valve. The rotor chord  $p_1p_t$  with its center at *m* slides and rotates following the path constituting the base circle. This rotation of the rotor is due to the expansion of the compressed gas as it flows into the chamber through the inlet. An inlet control valve is used to regulate this flow of compressed and costly fluid. Although limaçon-based positive displacement machines

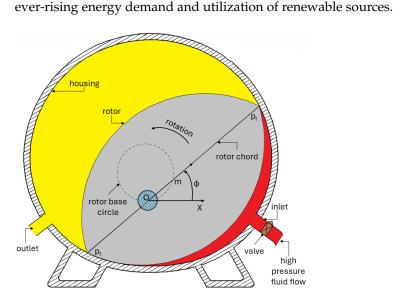


Figure 8. Limaçon expander with inlet valve.

As the gas expands in each cycle of operation, its volume, v (m<sup>3</sup>), and density,  $\rho$  (kg/m<sup>3</sup>), change, which can be described by the following nonlinear differential equations:

of manufacturing techniques, this technology can be an innovative tool to address the

$$v(t)\frac{d\rho(t)}{dt} = \frac{dm_i(t)}{dt} - \frac{dm_o(t)}{dt} - \rho(t)\frac{dv(t)}{dt} \text{ for } 0 \le t \le t_c$$
(24)

$$\rho(t)v(t)T(t)\frac{dS(t)}{dt} = (h_i(t) - h_c(t))\frac{dm_i(t)}{dt} - (h_c(t) - h_o(t))\frac{dm_o(t)}{dt} \text{ for } 0 \le t \le t_c \quad (25)$$

where  $m_i$  (kg) and  $m_o$  (kg) are the masses of the fluid at the inlet and outlet, respectively;  $h_i$ ,  $h_c$ , and  $h_o$  are the enthalpies in J/kg at the inlet, chamber, and outlet, respectively; T(K) and S are the temperature and entropy inside the chamber, respectively; and  $t_c$  (s) is the time in a cycle. Equations (24) and (25) can be solved iteratively until the following condition is met:

$$((\rho(t_c) - \rho(0))^2 + (S(t_c) - S(0))^2)^{\frac{1}{2}} \le \epsilon$$
(26)

where  $\epsilon$  is a small error value. The rate of fluid mass flowing into the chamber is a function of the dynamic passage area,  $A(\theta)$  (m<sup>2</sup>), as defined in (16); the pressure difference across the valve,  $\Delta P$  (kPa); and the fluid density inside the chamber,  $\rho$  (kg/m<sup>3</sup>). This can be represented by the following equation:

$$\frac{dm_i(t)}{dt} = f(A(\theta), \Delta P(t), \rho(t)) = f(\theta(t), \Delta P(t), \rho(t))$$
(27)

where  $\theta(t)$  is the rotation of the DDRV. The converted energy due to the expansion of the fluid inside the chamber can be expressed as:

$$E = \int_0^{t_c} P(t) \frac{dv(t)}{dt} dt$$
(28)

where *P* (kPa) is the chamber pressure. The isentropic efficiency,  $\eta$ , can, thus, be obtained from the ratio of the energy obtained to the maximum obtainable energy as follows:

$$\eta = \frac{E}{(h_i - h_o)\rho_i v_i} \tag{29}$$

where  $\rho_i$  (kg/m<sup>3</sup>) is the fluid density at the inlet and  $v_i$  (m<sup>3</sup>) is the chamber volume when the inlet valve is closed to stop fluid flow. Another important performance index is the filling factor,  $\zeta$ , which is the ratio of actual mass admitted into the chamber to the mass that could be incorporated into the chamber volume. This is given as follows:

$$\zeta = \frac{m_i}{\rho_i v_i} \tag{30}$$

As seen in Equation (27), the rate of mass flow into the expander depends on the valve rotation,  $\theta(t)$ , along with the pressure and density of the fluid. Therefore, the expander model should be simulated in conjunction with the DDRV model.

#### 4. Simulation of Actuators

The mathematical models of the proposed DDRVs as described in Equations (22) and (23) are solved analytically to simulate the valve dynamic response. Later, the DDRVs are tested against the thermodynamic model of the limaçon gas expander as described in Section 3.5. For the ease of analysis and clarity, the simulation of the valve–expander system is conducted in terms of the expander rotor angle to study the effect of the proposed DDRVs on various performance indices of the expander.

The introduction of the inlet valve to the gas expander enables it to selectively permit or obstruct the fluid into the chamber. The aim is to control the valve to operate at a certain predefined position of the rotor in its trajectory, as shown in Figure 9. The angular positions of the expander rotor at which the valve closes and opens are referred to as the cutoff and pass angle, respectively.

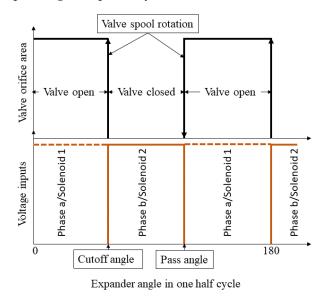


Figure 9. Valve operating sequence.

#### 4.1. Dynamic Response of the Valve

To simulate the operation of the valve in conjunction with the expander rotor angle  $\theta_{ex}$ , the sets of differential equations in (22) and (23) for the stepper motor and rotary solenoid actuators, respectively, are reformulated as shown in (31) and (32). These two sets of differential equations are solved in an iterative manner using Matlab (https://www.mathworks.com/products/matlab.html) with the different parameters provided in Tables 1–3. It is necessary to note the dissimilarity between the parameters of the SM-DDRV and RS-DDRV. This is due to the difference in the architecture and operating mechanism between the two actuators. The values are chosen to produce similar angular movements of the valve spool. A later study on the optimization of these valve parameters can yield a better speed of response.

$$\begin{split} \dot{i_d} &= \frac{1}{L_s \omega_{ex}} (V_d - R_s i_d + p \omega_m L_s i_q) \\ \dot{i_q} &= \frac{1}{L_s \omega_{ex}} (V_q - R_s i_q - p \omega_m L_s i_d - p \psi_m \omega) \\ \dot{\theta_m} &= \frac{\omega_m}{\omega_{ex}} \\ \dot{\omega_m} &= \frac{1}{J \omega_{ex}} (\tau_e - B \omega_m - \tau_f - \tau_l) \end{split}$$
(31)

where  $\omega_{ex}$  denotes the angular velocity of the expander rotor.

$$\begin{split} \dot{L}_{s} &= \frac{\omega_{ex}(1-2\mu_{r})L_{s}^{2}\dot{x}_{s}}{N_{s}^{2}\mu_{c}A_{e}} \\ \dot{i}_{s} &= \frac{1}{\omega_{ex}L_{s}}[V_{s}-R_{s}i_{s}] - i_{s}\frac{(1-2\mu_{r})L_{s}\dot{x}_{s}}{N_{s}^{2}\mu_{c}A_{e}} \\ \ddot{x}_{s} &= \frac{1}{M\omega_{ex}^{2}}(F_{s}+F_{p}-B\omega_{ex}\dot{x}_{s}-kx_{s}) \\ \dot{\theta}_{s} &= \frac{\omega_{s}}{\omega_{ex}} \\ \dot{\omega}_{s} &= \frac{1}{J\omega_{ex}}(\tau_{s}-B\omega_{s}-\tau_{f}-\tau_{l}) \end{split}$$
(32)

Table 1. Stepper motor data.

Parameters	Value	
Number of phases	2	
Phase voltages ( $V_a$ and $V_b$ )	20 V	
Step angle	$10^{\circ}$	
Winding self-inductance $(L_s)$	1 mH	
Winding resistance $(R_s)$	1.2 Ω	
Maximum flux linkage ( $\psi_m$ )	0.04 V s	

Table 2. Rotary solenoid data.

Parameters	Value
Applied voltage $(V_s)$	20 V
Number of turns $(N_s)$	500
Coil resistance $(R_s)$	$7.8\Omega$
Spring constant (k)	2000 N/m
Effective air gap $(A_e)$	$1.115  imes 10^{-3} \mathrm{m^2}$
Radius of ball bearing races (r)	$1.59  imes 10^{-2}  \mathrm{m}$

Both the stepper motor and rotary solenoid actuators are actuated by a voltage input, resulting in a current flow through their windings, as shown in Figure 10a,b. This flow of current produces mechanical movement. The various stages of response due to this current are shown in Figures 11–13.

Table 3. Valve data.

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Parameters	Value	
Total mass ( <i>M</i> )	$5.25  imes 10^{-2}  \mathrm{kg}$	
Total moment of inertia $(J)$	$5.47  imes 10^{-6}  \mathrm{kgm}^2$	
Total friction ( <i>B</i> )	$1 \times 10^{-3}  \text{kgm/s}$	
Supply pressure $(P_{in})$	1000 kPa	
Supply temperature $(T_{in})$	120 °C	
Valve ante-chamber pressure $(P_a)$	600 kPa	
Expander speed	800 rpm	
Diameter of orifice $(D_{orifice})$	25 mm	
Diameter of valve shaft $(D_{valve})$	15 mm	
Discharge velocity coefficient ( $C_v$ )	0.98	
Discharge coefficient ( $C_d$ )	0.65	
Jet angle $(\alpha_i)$	69°	
Cutoff angle ( $\theta_{cutoff}$ )	$90^{\circ}$	
Pass angle ( $\theta_{pass}$ )	$180^{\circ}$	

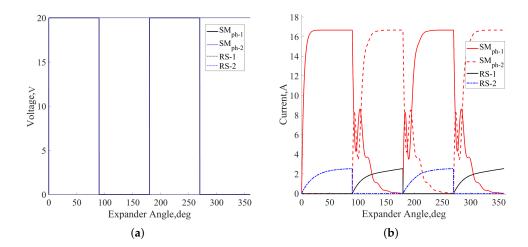


Figure 10. Input (a) phase voltage and (b) phase current.

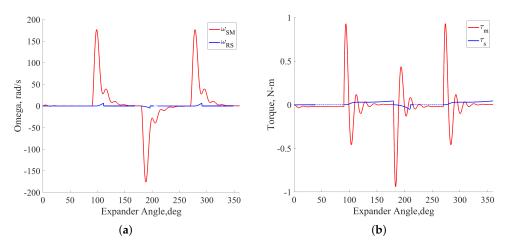


Figure 11. Valve response (a) angular velocity and (b) torque produced.

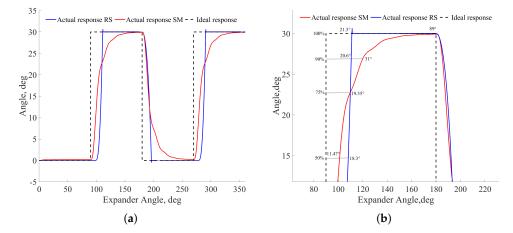


Figure 12. (a) Valve displacement response and (b) delay.

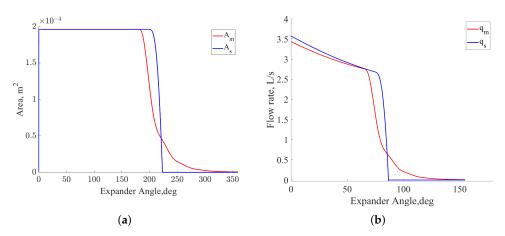


Figure 13. Variation of (a) valve area and (b) flow rate.

The angular velocity and torque produced in the rotary solenoid actuator are much less compared to the stepper motor actuator, as shown in Figure 11a,b. It can be seen that the stepper motor actuator produces higher torque and angular velocity compared to the rotary solenoid. Thus, the stepper motor actuator is faster in response to the same input voltage. The actuation speed of the rotary solenoid could be increased by increasing the input voltage and its configuration and architecture.

A comparison of the angular displacements for both actuators to the ideal characteristics is shown in Figure 12. Ideally, the valve should operate at the designated cutoff angle  $(90^{\circ})$  and pass angle  $(180^{\circ})$  to block and allow the fluid flow, respectively. But, as seen here, both actuators have some delay in responding to the voltage input. For better understanding and ease of depiction, a closing operation of the valve is shown in Figure 12b. It is seen that the valve spool could not attain its maximum rotation instantaneously, but rather, continues rotation throughout the closing cycle. Initially, at the half-closed position, the delay of the SM-DDRV is smaller at 11.47° compared to that of the RS-DDRV (18.3°). At the 75%-closed position, both actuators have a delay of 19.35°, but the SM-DDRV slows down afterwards towards the fully closed position, while the delay increases roughly to the closing duty cycle of 90°. The RS-DDRV, on the contrary, maintains a constant delay, and at full-closed position, the delay is 21.3°. The delay could be reduced by optimizing one or more of the parameters, including the duty cycle, step angle, and valve orifice diameter. The valve area and the resulting fluid flow rate vary accordingly, as shown in Figure 13. Both the valve area and flow rate variation are faster initially for the SM-DDRV compared to the RS-DDRV, but falter at the later stage.

Effect of Temperature and Friction

The effect of the temperature rise on the valve is studied to find the temperature limit of the actuators. Figure 14 shows the effect of temperature on the peak currents and maximum stroke. As seen in Figure 14a, the current in both actuators decreases as the temperature rises, increasing resistance. At high temperatures, the current will decrease to a value that is unable to produce any effective motion. This limiting temperature is different for the two actuators, as seen in Figure 14b, where 75 °C and 115 °C are the limiting temperatures for the RS-DDRV and SM-DDRV, respectively. However, proper thermal insulation and future optimization could be adopted to increase the temperature limit.

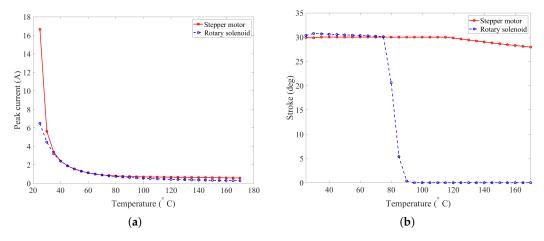


Figure 14. Change in (a) peak current and (b) stroke, due to temperature rise.

Besides temperature, friction also affects the overall system performance. The energy loss due to rotational and fluid friction at different inlet pressures is shown in Figure 15.

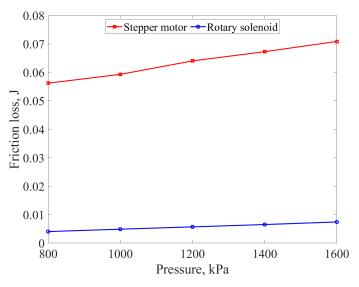


Figure 15. Frictional loss.

It is seen that frictional losses increase with pressure, and the SM-DDRV has higher frictional losses than the RS-DDRV. However, the losses incurred are negligible compared to the total energy converted, as will be seen in the next section.

# 4.2. Application of Proposed DDRVs to the Limaçon Expander

The performance of the SM-DDRV and RS-DDRV as the inlet control valves of a limaçon expander is validated. The thermodynamic model of the gas expander with the design parameters as stipulated in Table 4 is utilized for this study. The simulation of the expander for a full operating cycle with and without the proposed DDRVs is performed, and the resulting thermodynamic properties including the temperature, pressure, entropy, and volume in the expander chamber are calculated.

Table 4.	Data	for the	L2C	expander.
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Parameters	Value
Half of rotor chord length ( <i>L</i> )	46.4 mm
Base circle radius $(r)$	7.98 mm
Limaçon aspect ratio ( $b = \frac{r}{T}$ )	0.171
Housing–rotor clearance $(\tilde{C})$	0.71 mm
Clearance ratio ( $C_L = \frac{C}{L}$ )	0.0153
Design coefficient $(a)^{T}$	1.73
Depth of rotor housing $(H)$	60.38 mm
Fluid type	R245fa
Outlet pressure $(P_{out})$	100 kPa
Inlet port start angle	$-24.9^{\circ}$
Inlet port end angle	$-5.9^{\circ}$
Inlet port length	13.35 mm
Outlet port start angle	$140^{\circ}$
Outlet port end angle	175°
Outlet port length	21.47 mm

As seen from the pressure–volume (PV) diagram of Figure 16, the area of the PV diagram for SM-DDRV is comparable to the without valve case, but has a major distinction from that of the RS-DDRV, which is much less. The pressure builds up initially and is sustained for some time, but drops around 155 degree due to exposure to the discharge port, as shown in Figure 17a. The pressure build-up is higher for the expander without a valve, accounting for the constant flow through the inlet port. The pressure buildup in the RS-DDRV is much less compared to the SM-DDRV and port operation. The temperature and entropy changes also take place accordingly, as shown in Figure 17b,c. Both the entropy and temperature are high for the RS-DDRV compared to the other two cases, indicating lower energy conversion for this case. The various performance indices including mass flow, filling factor, isentropic efficiency, energy, and power are shown in Figure 18. In the case of the expander without a valve, the energy (83.91 J) and power (1.12 kW) output is highest compared to the other two cases. This is because the fluid flow is not restricted at the inlet port, which is reflected in a higher mass flow (3.17 gm). In the SM-DDRV, the energy (80.72 J) and indicated power (1.08 kW) output are similar to the case of the expander without a valve, whereas the RS-DDRV produces less energy (39.83 J) and power (0.53 kW) compared to the other two cases, which are accountable for its low mass flow rate of 1.35 gm. In terms of isentropic efficiency and the filling factor, which are two crucial performance indicators of the expansion process, the SM-DDRV has the upper hand with an isentropic efficiency of 63.14% and a filling factor of 0.64. The RS-DDRV also shows improvements in these two indices with an isentropic efficiency of 56.39% and a filling factor of 0.36. The values of these two indices are low in the ported without the value expander operation with an isentropic efficiency of 54.97% and a filling factor of just 0.21.

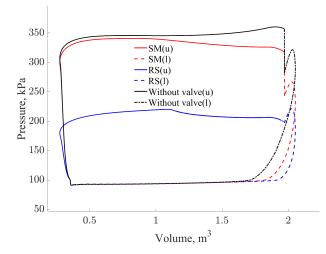


Figure 16. PV diagram.

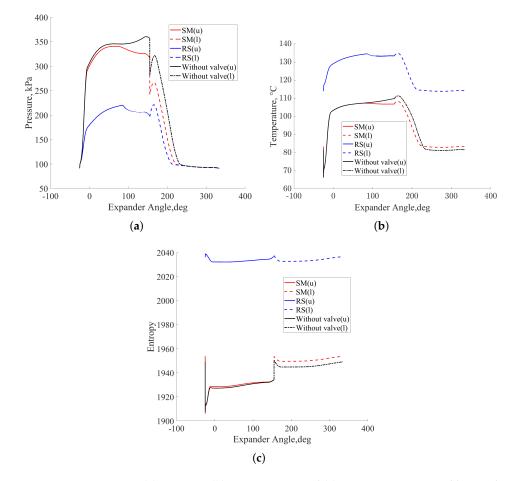
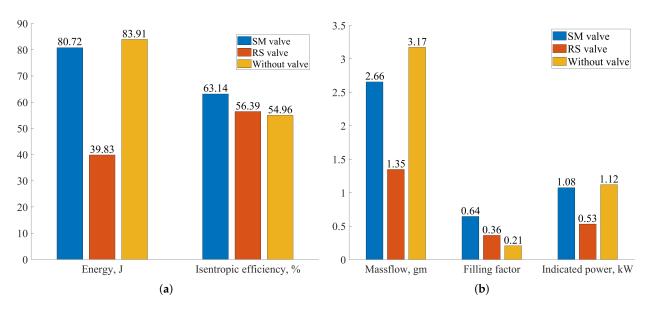


Figure 17. (a) Pressure, (b) temperature, and (c) entropy in upper and lower chamber.



**Figure 18.** Expander outputs: (**a**) energy, and isentropic efficiency; (**b**) mass flow, filling factor, and indicated power.

#### 4.3. Sensitivity to Inlet Pressure

In the above analysis, the fluid pressure at the inlet is considered constant. However, it is now required to check how the valves react to a changing inlet pressure. The pressure is varied from 800 kPa to 1600 kPa, and the sensitivity of the valves is analyzed. Figure 19a depicts the gas leakage through the orifice due to the delay in valve response while closing.

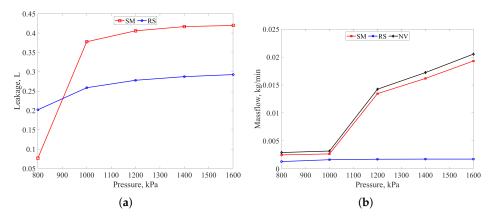
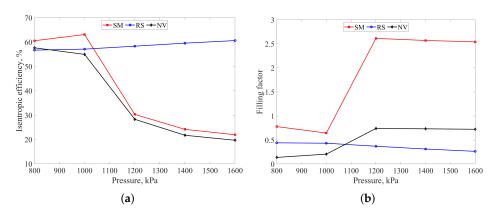


Figure 19. (a) Leakage and (b) mass flow at different inlet pressures.

It is seen that at low pressure of 800 kPa, the SM-DDRV has a minimum leakage of 0.077 L, while the RS-DDRV has a leakage of 0.20 L. However, as the pressure increases to 1000 kPa, the leakage for the SM-DDRV increases considerably to 0.38 L. The leakage continues to increase steadily and reaches a maximum value of 0.42 L at a pressure of 1600 kPa, whereas the RS-DDRV maintains a relatively constant leakage with the pressure increase and reaches a maximum leakage of 0.29 L at 1600 kPa. An increased leakage means an increased fluid flow into the expander chamber, as seen in Figure 19b. At higher pressures, the mass flowing into the chamber will be high if not properly blocked by the valve. The SM-DDRV has a higher leakage at increased pressures, so the mass flow is high in this case, but less than that of the expander with no valve. The mass flow is less for the RS-DDRV as it has lower leakage as compared to the other two cases. The gas leakage affects the performance indices of the expander, as will be seen as follows.

As seen in Equations (29) and (30), the isentropic efficiency and filling factor depend on the amount of mass flowing into the chamber. Figure 20a shows the variation of the isentropic efficiency for different inlet pressures. At low pressure from 800–1000 kPa, the isentropic efficiency for the SM-DDRV is high and reaches a maximum of 63.14%. However, at higher pressures, the efficiency reduces as the mass flow increases. On the contrary, for the RS-DDRV, the efficiency is less at lower pressure, but increases with pressure. The maximum achieved efficiency with the RS-DDRV is 60.63% at 1600 kPa. Figure 20b shows the filling factor at different pressures. As seen above, the SM-DDRV has a higher mass flow compared to the RS-DDRV, so, the filling factor will be high for this valve. However, at higher pressure, the leakage is high, which contributes to a filling factor calculation above 1 as more fluid flow is more than the intended amount. It can be deduced from Figure 20b that the maximum filling factor less than 1 is achieved by the SM-DDRV at 1000 kPa. The filling factor at low pressures with the RS-DDRV is higher than the no-valve case but less



than the SM-DDRV; however, it decreases at higher pressures.

Figure 20. (a) Isentropic efficiency and (b) filling factor at different inlet pressures.

#### 5. Results and Discussion

The case studies provided here give insights into improving the process efficiency of the expander. As seen in Figure 12, both the SM-DDRV and RS-DDRV have some delay in response due to the time taken to build up current in the windings and the resultant mechanical displacement of the valve spool. However, the SM-DDRV has a faster response initially compared to the RS-DDRV, but slows down at the end of the opening or closing cycle, which is reflected in its performance indices. At lower pressures of less than 1000 kPa, the SM-DDRV has better performance than the RS-DDRV in improving the isentropic efficiency and filling factor. At 1000 kPa, the SM-DDRV improves the isentropic efficiency by 14.88% and the filling factor by 220% compared to the ported expander. However, at higher pressures, the performance of the SM-DDRV degrades as the leakage increases. Further optimization of both valve parameters should yield better dynamic performance at elevated inlet pressures. Nevertheless, the proposed SM-DDRV valve is preferable at pressures less than 1000 kPa, which improves the overall process efficiency. But, at higher pressures, the RS-DDRV is preferable. The isentropic efficiency of 63.14% achieved with the SM-DDRV is an improvement from the previously reported study conducted by Sultan et al. [24], where they reported an efficiency of 58.67%. This study provides a good theoretical base for future applications of limaçon expanders in ORC-based plants.

# 6. Conclusions

Enhancing the efficiency of compact and renewable-based power-generation systems, like the one examined here, can significantly contribute to addressing the current energy crisis and reducing the effects of greenhouse gas emissions. The valves proposed in this work hold promise as a valuable instrument for enhancing the operational efficiency of such setups. The effectiveness of the proposed valves has been analyzed theoretically and will be validated by experiments in the future. The mathematical models and simulated outcomes will also offer essential guidelines for the future development and deployment of similar power-generation units. In light of the above comparative studies, the SM-DDRV

presents promising results in improving isentropic efficiency and the filling factor. However, the performance degrades at higher inlet pressure, whereas the RS-DDRV has moderate performance at different pressures. The proposed valves can be optimized in future work, which should improve the performance indices further.

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# Abbreviations

The following abbreviations are used in this manuscript:

ORC	organic Rankine cycle
DDRV	direct-drive rotary valve
SM-DDRV	stepper motor DDRV
RS-DDRV	rotary solenoid DDRV
WHR	waste heat recovery
CHP	combined heat and power
GWP	Global Warming Potential
ODP	Ozone Depletion Potential

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