

# Applying Discrete Homotopy Analysis Method for Solving Fractional Partial Differential Equations

Figen Özpınar 

Bolvadin Vocational School, Afyon Kocatepe University, 03300 Afyonkarahisar, Turkey; fozpinar@aku.edu.tr;  
Tel.: +90-272-612-6353

Received: 3 March 2018; Accepted: 27 April 2018; Published: 1 May 2018



**Abstract:** In this paper we developed a space discrete version of the homotopy analysis method (DHAM) to find the solutions of linear and nonlinear fractional partial differential equations with time derivative  $\alpha$  ( $0 < \alpha \leq 1$ ). The DHAM contains the auxiliary parameter  $\hbar$ , which provides a simple way to guarantee the convergence region of solution series. The efficiency and accuracy of the proposed method is demonstrated by test problems with initial conditions. The results obtained are compared with the exact solutions when  $\alpha = 1$ . It is shown they are in good agreement with each other.

**Keywords:** discrete homotopy analysis method; Caputo fractional derivative; fractional discrete diffusion equation; fractional discrete Schrödinger equation; fractional discrete Burgers' equation

## 1. Introduction

Fractional calculus has been of increasing interest to scientists and engineers, arising in mathematical physics, chemistry, modeling mechanical and electrical properties of real phenomena [1–6]. Fractional calculus has been recognized as a powerful instrument to discover the secret directions of various material and physical processes that deal with derivatives and integrals of arbitrary orders [7–16].

Various techniques have been investigated to solve partial differential equations of fractional order, such as the homotopy analysis method (HAM) [17–22], homotopy perturbation method (HPM) [23–26], Adomian decomposition method (ADM) [27–29], meshless method [30–33], operational matrix [34,35] and so on. In 1992, Liao introduced the homotopy analysis method, a semi-analytical method, for solving strongly nonlinear differential equations [20]. The main advantage of HAM is that it provides great freedom to choose equation type and solution expression of related linear high-order approximation equations. HAM gives rapidly convergent successive approximations of the exact solutions, if such a solution exists, otherwise approximations can be used for numerical purposes. It is an analytical approach to get the series solution of linear and nonlinear partial differential equations. Unlike the other analytical techniques, HAM is independent of small/large physical parameters. Since HAM has many advantages in comparison to other analytical methods, it is employed to solve continuous problems. Hence after the discrete ADM method [36], the discrete homotopy analysis method (DHAM) was introduced in 2010 by Zhu et al. [37]. This method can be applied to complex problems containing discontinuity in fluid characteristics and geometry of the problem. In addition, it needs little computational cost as numerical method in comparison to HAM; as an analytical approach DHAM has similar advantages to continuous HAM. By means of introducing an auxiliary parameter one can adjust and control the convergence region of the solution series. This method should be employed for solving various differential equations to highlight its high capabilities in comparison with other numerical methods.

In this study, we develop the discrete homotopy analysis method (DHAM) for the fractional discrete diffusion equation, nonlinear fractional discrete Schrödinger equation and nonlinear fractional discrete Burgers' equation with time derivative  $\alpha$  ( $0 < \alpha \leq 1$ ). The approximate analytical solutions of the test problems are obtained using initial conditions. The obtained solutions are verified by comparison with exact solution when  $\alpha = 1$ .

## 2. Preliminaries and Notations

### 2.1. Fractional Analysis

**Definition 1** ([38]). A real function  $f(x)$ ,  $x > 0$  is said to be in space  $C_\alpha$ ,  $\alpha \in \mathbb{R}$  if there exists a real number  $p > \alpha$  such that  $f(x) = x^p f_1(x)$  where  $f_1(x) \in C[0, \infty)$ .

**Definition 2** ([38]). A function  $f(x)$ ,  $x > 0$  is said to be in space  $C_\alpha^m$ ,  $m \in \mathbb{N} \cup \{0\}$  if  $f^m \in C_\alpha$ .

**Definition 3** ([5]). Let  $f \in C_\alpha$  and  $\alpha \geq -1$ , then Riemann-Liouville fractional integral of  $f(x, t)$  with respect to  $t$  of order  $\alpha$  is denoted by  $J^\alpha f(x, t)$  and is defined as

$$J^\alpha f(x, t) = \frac{1}{\Gamma(\alpha)} \int_0^t (t - \tau)^{\alpha-1} f(x, \tau) d\tau, \quad t > 0, \alpha > 0.$$

The well-known property of Riemann-Liouville operator  $J^\alpha$  is

$$J^\alpha t^\gamma = \frac{\Gamma(\gamma + 1) t^{\gamma+\alpha}}{\Gamma(\gamma + \alpha + 1)}.$$

**Definition 4** ([39]). Form to be the smallest integer that exceeds  $\alpha > 0$ , the Caputo fractional derivative of  $u(x, t)$  with respect to  $t$  of order  $\alpha > 0$  is defined as

$$D_t^\alpha u(x, t) = \frac{\partial^\alpha u(x, t)}{\partial t^\alpha} = \begin{cases} \frac{1}{\Gamma(m-\alpha)} \int_0^t (t - \tau)^{m-\alpha-1} \frac{\partial^m u}{\partial t^m} d\tau, & \text{for } m-1 < \alpha < m \\ \frac{\partial^m u(x, t)}{\partial t^m}, & \text{for } \alpha = m \in \mathbb{N}. \end{cases}$$

Note that the Caputo fractional derivative is considered due to its suitable for initial conditions of the differential equations.

The relations between Riemann-Liouville operator and Caputo fractional differential operator are given as follows

$$D^\alpha (J^\alpha f(x, t)) = f(x, t), J^\alpha (D^\alpha f(x, t)) = J^\alpha (J^{m-\alpha} f^{(m)}(x, t)) = J^m f^{(m)}(x, t) = f(x, t) - \sum_{k=0}^{m-1} f^{(k)}(x, 0) \frac{t^k}{k!}.$$

### 2.2. Discrete Homotopy Analysis Method

Consider the following general difference equation respect to  $j$

$$\mathcal{N}[u_j(t)] = 0, \quad j \in \mathbb{Z}, \quad t \in \mathbb{R}, \quad (1)$$

where  $\mathcal{N}$  is a linear or nonlinear operator,  $j$  and  $t$  denote the independent variables. Suppose that  $\Delta x = h$  and the function  $u(x, t) = u(j\Delta x, t)$  is the discrete function and denoted by  $u_j(t)$ .

For simplicity, we ignore all boundary or initial conditions, which can be treated in the similar way. Similarly to continuous HAM, we first construct the so-called zeroth-order deformation equation by means of the discrete HAM (DHAM)

$$(1-p)\mathcal{L}[\varphi_j(t;p) - u_{j,0}(t)] = p\hbar H_j(t)\mathcal{N}[\varphi_j(t;p)], \quad (2)$$

where  $p \in [0, 1]$  is an embedding parameter,  $\hbar \neq 0$  is an auxiliary parameter,  $\mathcal{L}$  is an auxiliary linear operator,  $u_{j,0}(t)$  is an initial guess of  $u_j(t)$ ,  $H_j(t)$  denotes a nonzero auxiliary function,  $\varphi_j(t;p)$  is an unknown function about  $j, t, p$ . It is important that one has great freedom to choose auxiliary things in (2). Obviously, when  $p = 0$  and  $p = 1$ , it holds

$$\varphi_j(t;0) = u_{j,0}(t) = u_j(0) \text{ and } \varphi_j(t;1) = u_j(t),$$

respectively. Thus, as  $p$  increases from 0 to 1, the solution  $\varphi_j(t;p)$  varies from initial guess  $u_{j,0}(t)$  to the solution  $u_j(t)$ . Expanding  $\varphi_j(t;p)$  in Taylor series with respect to  $p$ , we have

$$\varphi_j(t;p) = u_{j,0}(t) + \sum_{m=1}^{+\infty} u_{j,m}(t)p^m, \quad (3)$$

where

$$u_{j,m}(t) = \frac{1}{m!} \left. \frac{\partial^m \varphi_j(t;p)}{\partial p^m} \right|_{p=0} \quad (4)$$

Similarly continuous HAM by Liao [21], if the auxiliary linear operator, the initial guess, the auxiliary parameter  $\hbar$ , and the auxiliary function are so properly chosen, the series (3) converges at  $p = 1$ , then we have

$$u_j(t) = u_{j,0}(t) + \sum_{m=1}^{+\infty} u_{j,m}(t). \quad (5)$$

As  $\hbar = -1$  and  $H_j(t) = 1$ , Equation (2) becomes

$$(1-p)\mathcal{L}[\varphi_j(t;p) - u_{j,0}(t)] + p\mathcal{N}[\varphi_j(t;p)] = 0;$$

which is used in the discrete homotopy perturbation method (DHPM) [16], where as the solution obtained directly, without using Taylor series.

According to Equation (4), the governing equation can be deduced from the zeroth-order deformation Equation (DHAM). Define the vector

$$\vec{u}_n = \{u_{j,0}(t), u_{j,1}(t), \dots, u_{j,n}(t)\}.$$

Differentiating the zeroth-order deformation Equation (2)  $m$  times with respect to the embedding parameter  $p$  and then setting  $p = 0$  and finally dividing them by  $m!$ , we obtain the following  $m$ th-order deformation equation:

$$\mathcal{L}[u_{j,m}(t) - \mathcal{X}_m u_{j,m-1}(t)] = \hbar H_j(t) \mathcal{R}_m [\vec{u}_{m-1}], \quad (6)$$

where

$$\mathcal{R}_m [\vec{u}_{m-1}] = \frac{1}{(m-1)!} \left. \frac{\partial^{m-1} \mathcal{N}[\varphi_j(t;p)]}{\partial p^{m-1}} \right|_{p=0} \quad (7)$$

and

$$\mathcal{X}_m = \begin{cases} 0, & m \leq 1, \\ 1, & m > 1. \end{cases}$$

It should be emphasized that it is very important to ensure that Equation (3) converges at  $p = 1$  otherwise, the Equation (5) has no meaning.

**Theorem 1** (Convergence Theorem). *As long as the series (5) is convergent, where  $u_{j,m}(t)$  is governed by the high deformation Equation (6). It must be the solution of the original Equation (1).*

**Proof.** If the series  $\sum_{m=0}^{+\infty} u_{j,m}(t)$  is convergent, we can write

$$s_j(t) = u_{j,0}(t) + \sum_{m=1}^{+\infty} u_{j,m}(t), \quad (8)$$

and it holds

$$\lim_{n \rightarrow \infty} u_{j,n}(t) = 0.$$

From the  $m$ th-order deformation Equation (6) and by using the definition of  $\mathcal{X}_m$ , it yields

$$\begin{aligned} \sum_{m=1}^{+\infty} \hbar H_j(t) \mathcal{R}_m \left[ \vec{u}_{m-1} \right] &= \sum_{m=1}^{+\infty} \mathcal{L} [u_{j,m}(t) - \mathcal{X}_m u_{j,m-1}(t)] \\ &= \lim_{n \rightarrow \infty} \sum_{m=1}^n \mathcal{L} [u_{j,m}(t) - \mathcal{X}_m u_{j,m-1}(t)] \\ &= \mathcal{L} \left[ \lim_{n \rightarrow \infty} \sum_{m=1}^n (u_{j,m}(t) - \mathcal{X}_m u_{j,m-1}(t)) \right] \\ &= \mathcal{L} \left[ \lim_{n \rightarrow \infty} (u_{j,1}(t) + (u_{j,2}(t) - u_{j,1}(t)) + (u_{j,3}(t) - u_{j,2}(t)) + \dots \right. \\ &\quad \left. + (u_{j,n}(t) - u_{j,n-1}(t))) \right] \\ &= \mathcal{L} \left[ \lim_{n \rightarrow \infty} u_{j,n}(t) \right] = 0. \end{aligned} \quad (9)$$

Since  $\hbar \neq 0$ ,  $H_j(t) \neq 0$ , then

$$\sum_{m=1}^{+\infty} \mathcal{R}_m \left[ \vec{u}_{m-1} \right] = 0. \quad (10)$$

On the other side, according to the definition (7), we have

$$\sum_{m=1}^{+\infty} \mathcal{R}_m \left[ \vec{u}_{m-1} \right] = \sum_{m=1}^{+\infty} \frac{1}{(m-1)!} \left. \frac{\partial^{m-1} \mathcal{N}[\varphi_j(t;p)]}{\partial p^{m-1}} \right|_{p=0} = \sum_{m=0}^{+\infty} \frac{1}{m!} \left. \frac{\partial^m \mathcal{N}[\varphi_j(t;p)]}{\partial p^m} \right|_{p=0} = 0. \quad (11)$$

In general,  $\varphi_j(t;p)$  doesn't satisfy the original Equation (1). Let

$$\varepsilon_j(t;p) = \mathcal{N}[\varphi_j(t;p)]$$

denote the residual error of Equation (1). Obviously,

$$\varepsilon_j(t;p) = 0$$

corresponds to the exact solution of the Equation (1).

According to the above definition, the Maclaurin series of the residual error  $\varepsilon_j(t;p)$  with respect to the embedding parameter  $p$  is

$$\varepsilon_j(t;p) = \sum_{m=0}^{+\infty} \left( \left. \frac{1}{m!} \frac{\partial^m \varepsilon_j(t;p)}{\partial p^m} \right|_{p=0} \right) p^m = \sum_{m=0}^{+\infty} \left( \left. \frac{1}{m!} \frac{\partial^m \mathcal{N}[\varphi_j(t;p)]}{\partial p^m} \right|_{p=0} \right) p^m.$$

when  $p = 1$ , the above expression gives

$$\varepsilon_j(t; 1) = \sum_{m=0}^{+\infty} \frac{1}{m!} \frac{\partial^m \mathcal{N}[\varphi_j(t; p)]}{\partial p^m} \Big|_{p=0} = 0. \quad (12)$$

That is, according to the definition of  $\varepsilon_j(t; p)$  we have the exact solution of the original Equation (1) when  $p = 1$ . Thus as long as the series

$$u_{j,0}(t) + \sum_{m=1}^{+\infty} u_{j,m}(t)$$

is convergent, it must be the solution of the original Equation (1).  $\square$

### 3. Examples

**Example 1.** Consider the time fractional discrete diffusion equation

$$D_t^\alpha u_j(t) = D_h^2 u_j(t) + jh D_h u_j(t) + u_j(t), \quad 0 < \alpha \leq 1 \quad (13)$$

with initial condition

$$u_j(0) = jh. \quad (14)$$

Discrete diffusion equation is widely used in applied sciences. For example, population growth modeled by geographical spread [40], to model ionic diffusion on a lattice [41], and so on. Moreover, the entropy production was calculated for fractional diffusion Equation [42,43].

The standard central differences  $D_h u_j(t)$  and  $D_h^2 u_j(t)$  are defined by

$$D_h u_j(t) = \frac{u_{j+1}(t) - u_{j-1}(t)}{2h}, \quad D_h^2 u_j(t) = \frac{u_{j+1}(t) - 2u_j(t) + u_{j-1}(t)}{h^2}.$$

Initial value problem (13) and (14) is the discrete form of initial value problem for diffusion equation

$$D_t^\alpha u(x, t) = u_{xx}(x, t) + xu_x(x, t) + u(x, t), \quad 0 < \alpha \leq 1$$

with initial condition

$$u(x, 0) = x,$$

where  $D_t^\alpha u(x, t)$  is Caputo fractional derivative of order  $\alpha$ .

To solve Equation (13) by DHAM let us consider the following linear operator:

$$\mathcal{L}[\varphi_j(t; p)] = D_t^\alpha [\varphi_j(t; p)] = \frac{\partial^\alpha \varphi_j(t; p)}{\partial t^\alpha} \quad (15)$$

with the property that

$$\mathcal{L}[c] = 0,$$

where  $c$  is constant coefficients. We define the nonlinear operator as

$$\mathcal{N}[\varphi_j(t; p)] = D_t^\alpha \varphi_j(t; p) - D_h^2 \varphi_j(t; p) - jh D_h \varphi_j(t; p) - \varphi_j(t; p). \quad (16)$$

Using the above definition, we construct the zeroth-order deformation equation by Equation (2).

It is obvious that when the embedding parameter  $p = 0$  and  $p = 1$ , Equation (2) becomes

$$\varphi_j(t; 0) = u_{j,0}(t) = u_j(0), \quad \varphi_j(t; 1) = u_j(t),$$

respectively. Then we obtain the  $m$ th-order deformation equation for  $m \geq 1$  with

$$\mathcal{L}[u_{j,m}(t) - \mathcal{X}_m u_{j,m-1}(t)] = \hbar H_j(t) \mathcal{R}_m[\vec{u}_{m-1}] \Rightarrow u_{j,m}(t) = \mathcal{X}_m u_{j,m-1}(t) + \hbar J^\alpha [H_j(t) \mathcal{R}_m[\vec{u}_{m-1}]], \quad (17)$$

where

$$\begin{aligned} \mathcal{R}_m[\vec{u}_{m-1}] &= \frac{1}{(m-1)!} \frac{\partial^{m-1} \mathcal{N}[\varphi_j(t;p)]}{\partial p^{m-1}} \Big|_{p=0} \\ &= D_t^\alpha u_{j,m-1}(t) - D_h^2 u_{j,m-1}(t) - jh D_h u_{j,m-1}(t) - u_{j,m-1}(t). \end{aligned} \quad (18)$$

For simplicity, we select  $H_j(t) = 1$  in this problem. So, the approximations of  $u_j(t)$  are only depend on auxiliary parameter  $\hbar$ .

Solve the above equation under the initial condition

$$u_{j,0}(t) = u_j(0) = jh$$

we get

$$\begin{aligned} u_{j,1}(t) &= -jh \frac{2t^\alpha}{\Gamma(\alpha+1)} \hbar \\ u_{j,2}(t) &= -jh \frac{2t^\alpha}{\Gamma(\alpha+1)} \hbar(\hbar+1) + jh \frac{2^2 t^{2\alpha}}{\Gamma(2\alpha+1)} \hbar^2 \\ u_{j,3}(t) &= -jh \frac{2t^\alpha}{\Gamma(\alpha+1)} \hbar(\hbar+1)^2 + jh \frac{2^2 t^{2\alpha}}{\Gamma(2\alpha+1)} 2\hbar^2(\hbar+1) - jh \frac{2^3 t^{3\alpha}}{\Gamma(3\alpha+1)} \hbar^3 \\ u_{j,4}(t) &= -jh \frac{2t^\alpha}{\Gamma(\alpha+1)} \hbar(\hbar+1)^3 + jh \frac{2^2 t^{2\alpha}}{\Gamma(2\alpha+1)} 3\hbar^2(\hbar+1)^2 - jh \frac{2^3 t^{3\alpha}}{\Gamma(3\alpha+1)} 3\hbar^3(\hbar+1) + jh \frac{2^4 t^{4\alpha}}{\Gamma(4\alpha+1)} \hbar^4 \end{aligned}$$

Thus the rest of components  $u_n$ ,  $n > 4$  of the DHAM can be completely obtained. So, we approximate the analytical solution

$$\begin{aligned} u_j(t) &= u_{j,0}(t) + u_{j,1}(t) + u_{j,2}(t) + u_{j,3}(t) + u_{j,4}(t) + \dots \\ &= jh - jh \frac{2t^\alpha}{\Gamma(\alpha+1)} \hbar - jh \frac{2t^\alpha}{\Gamma(\alpha+1)} \hbar(\hbar+1) + jh \frac{2^2 t^{2\alpha}}{\Gamma(2\alpha+1)} \hbar^2 \\ &\quad - jh \frac{2t^\alpha}{\Gamma(\alpha+1)} \hbar(\hbar+1)^2 + jh \frac{2^2 t^{2\alpha}}{\Gamma(2\alpha+1)} 2\hbar^2(\hbar+1) - jh \frac{2^3 t^{3\alpha}}{\Gamma(3\alpha+1)} \hbar^3 \\ &\quad - jh \frac{2t^\alpha}{\Gamma(\alpha+1)} \hbar(\hbar+1)^3 + jh \frac{2^2 t^{2\alpha}}{\Gamma(2\alpha+1)} 3\hbar^2(\hbar+1)^2 - jh \frac{2^3 t^{3\alpha}}{\Gamma(3\alpha+1)} 3\hbar^3(\hbar+1) \\ &\quad + jh \frac{2^4 t^{4\alpha}}{\Gamma(4\alpha+1)} \hbar^4 + \dots \end{aligned}$$

Setting  $\hbar = -1$ , we get an accurate approximation solution in the following form:

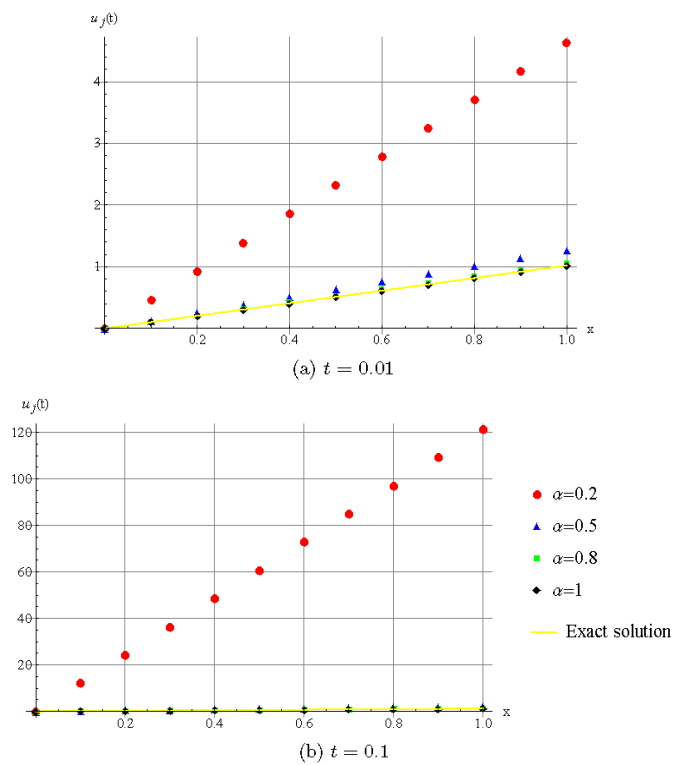
$$\begin{aligned} u_j(t) &= jh + jh \frac{2t^\alpha}{\Gamma(\alpha+1)} + jh \frac{2^2 t^{2\alpha}}{\Gamma(2\alpha+1)} + jh \frac{2^3 t^{3\alpha}}{\Gamma(3\alpha+1)} + jh \frac{2^4 t^{4\alpha}}{\Gamma(4\alpha+1)} + \dots \\ u_j(t) &= \sum_{n=0}^{\infty} (jh) \frac{2^n t^{n\alpha}}{\Gamma(n\alpha+1)} = (jh) E_\alpha(2t^\alpha), \end{aligned}$$

where  $E_\alpha$  is Mittag-Leffler function.

$$u(x, t) = x E_\alpha(2t^\alpha)$$

is the exact solution of the continuous form.

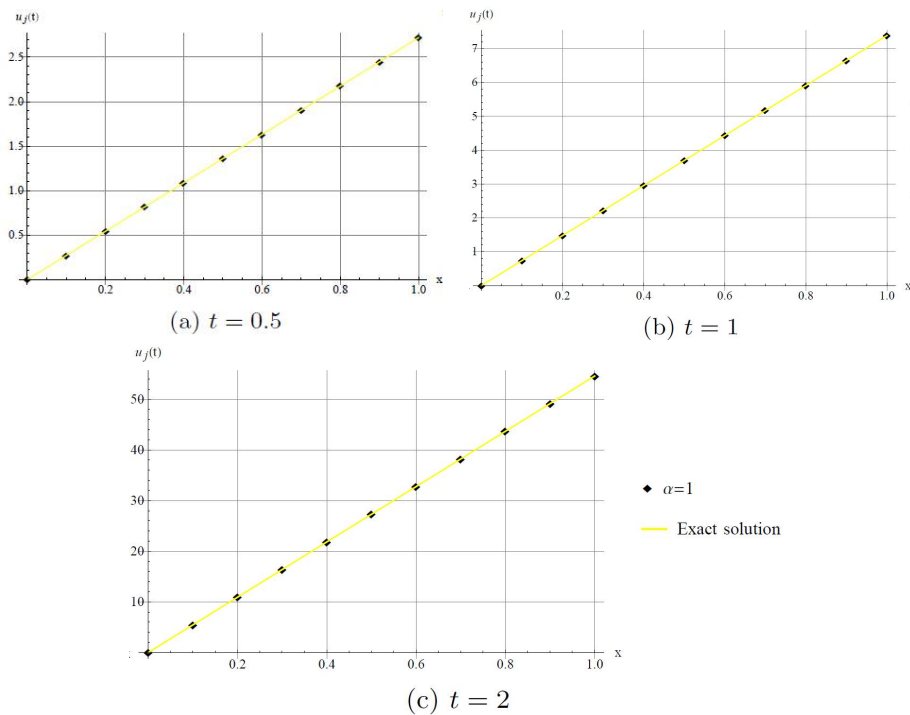
Figure 1 shows the DHAM approximate solution of  $u(x, t)$  for different values of  $\alpha$ .



**Figure 1.** Numerical illustration of approximation solution  $u(x,t)$  by discrete homotopy analysis method (DHAM). (a) For  $t = 0.01$ ; (b) For  $t = 0.1$ .

We can see that different fractional order lead to different diffusion behaviors.

In Figure 2, we show that the method has good agreement with the exact solution when  $\alpha = 1$ .



**Figure 2.** Comparison with numerical solution of  $u(x,t)$  by DHAM and the exact solution when  $\alpha = 1$ . (a) For  $t = 0.5$ ; (b) For  $t = 1$ ; (c) For  $t = 2$ .

**Example 2.** Consider the nonlinear fractional discrete Schrödinger equation

$$iD_t^\alpha u_j(t) + D_h^2 u_j(t) + q|u_j(t)|^2 u_j(t) = 0, \quad j \in \mathbb{Z}, \quad t > 0, \quad 0 < \alpha \leq 1, \quad (19)$$

with initial condition

$$u_j(0) = e^{ijkh}. \quad (20)$$

Discrete nonlinear Schrödinger equation is widely used in applied sciences. Describing the propagation of electromagnetic waves in glass fibers, one-dimensional arrays of coupled optical waveguides [18] and light-induced photonic crystal lattices [44]. Moreover, they are an established model for optical pulse propagation in various doped fibers [45,46].

Discrete nonlinear Schrödinger equations are also called lattice nonlinear Schrödinger equations [47].

The parameter  $\varepsilon = h^{-2}$  is called (discrete) dispersion and the parameter  $q$  is called anharmonicity.

Initial value problem (19) and (20) is the discrete form of initial value problem for Schrödinger equation

$$iD_t^\alpha u(x, t) + u_{xx}(x, t) + q|u(x, t)|^2 u(x, t) = 0, \quad t > 0, \quad 0 < \alpha \leq 1,$$

with initial condition

$$u(x, 0) = e^{ikx}.$$

we set  $|u_j(t)|^2 u_j(t) = u_j^2(t) \bar{u}_j(t)$ .

By means of DHAM, we choose the linear operator:

$$\mathcal{L}[\varphi_j(t; p)] = D_t^\alpha [\varphi_j(t; p)] = \frac{\partial^\alpha \varphi_j(t; p)}{\partial t^\alpha} \quad (21)$$

with property  $\mathcal{L}[c] = 0$ , where  $c$  is constant. We define the nonlinear operator as

$$\mathcal{N}[\varphi_j(t; p)] = D_t^\alpha \varphi_j(t; p) - iD_h^2 \varphi_j(t; p) - iq\varphi_j^2(t; p)\bar{\varphi}_j(t; p). \quad (22)$$

we construct the zeroth-order deformation equation by Equation (2).

For  $p = 0$  and  $p = 1$ , we can write

$$\varphi_j(t; 0) = u_{j,0}(t) = u_j(0), \quad \varphi_j(t; 1) = u_j(t),$$

respectively. Thus, we obtain the  $m$ th-order deformation equation

$$\mathcal{L}[u_{j,m}(t) - \mathcal{X}_m u_{j,m-1}(t)] = \hbar H_j(t) \mathcal{R}_m [\vec{u}_{m-1}], \quad (23)$$

where

$$\mathcal{R}_m [\vec{u}_{m-1}] = D_t^\alpha u_{j,m-1}(t) - iD_h^2 u_{j,m-1}(t) - iq \left[ \sum_{l=0}^{m-1} \sum_{n=0}^{m-l-1} u_{j,l}(t) u_{j,n}(t) \bar{u}_{j,m-n-l-1}(t) \right]. \quad (24)$$

we can select again  $H_j(t) = 1$ . Thus, the approximations of  $u_j(t)$  are only depend on auxiliary parameter  $\hbar$ .

Therefore the solution of the  $m$ th-order deformation Equation (23) for  $m \geq 1$  become

$$u_{j,m}(t) = \mathcal{X}_m u_{j,m-1}(t) + \hbar J^\alpha \mathcal{R}_m [\vec{u}_{m-1}]. \quad (25)$$

Substituting the initial condition (20) into the system (25), we get

$$\begin{aligned} u_{j,1}(t) &= \frac{i\omega e^{ijkh} t^\alpha}{\Gamma(\alpha+1)} \hbar \\ u_{j,2}(t) &= \frac{i\omega e^{ijkh} t^\alpha}{\Gamma(\alpha+1)} \hbar(\hbar+1) - \frac{\omega^2 e^{ijkh} t^{2\alpha}}{\Gamma(2\alpha+1)} \hbar^2 \\ u_{j,3}(t) &= \frac{i\omega e^{ijkh} t^\alpha}{\Gamma(\alpha+1)} \hbar(\hbar+1)^2 - \frac{2\omega^2 e^{ijkh} t^{2\alpha}}{\Gamma(2\alpha+1)} \hbar^2(\hbar+1) \\ &\quad - \frac{i\omega^3 e^{ijkh} t^{3\alpha}}{\Gamma(3\alpha+1)} \hbar^3 \left[ 1 - \frac{2q}{\omega} + \frac{q\Gamma(2\alpha+1)}{\omega(\Gamma(\alpha+1))^2} \right] \\ &\quad \vdots \end{aligned}$$

where  $\omega = (4/h^2)\sin^2(kh/2) - q$  (discrete dispersion relation).

Thus, we can conclude that

$$\begin{aligned} u_j(t) &= u_{j,0}(t) + u_{j,1}(t) + u_{j,2}(t) + u_{j,3}(t) + \dots \\ &= e^{ijkh} + \frac{i\omega e^{ijkh} t^\alpha}{\Gamma(\alpha+1)} \hbar + \frac{i\omega e^{ijkh} t^\alpha}{\Gamma(\alpha+1)} \hbar(\hbar+1) \\ &\quad - \frac{\omega^2 e^{ijkh} t^{2\alpha}}{\Gamma(2\alpha+1)} \hbar^2 + \frac{i\omega e^{ijkh} t^\alpha}{\Gamma(\alpha+1)} \hbar(\hbar+1)^2 - \frac{2\omega^2 e^{ijkh} t^{2\alpha}}{\Gamma(2\alpha+1)} \hbar^2(\hbar+1) \\ &\quad - \frac{i\omega^3 e^{ijkh} t^{3\alpha}}{\Gamma(3\alpha+1)} \hbar^3 \left[ 1 - \frac{2q}{\omega} + \frac{q\Gamma(2\alpha+1)}{\omega(\Gamma(\alpha+1))^2} \right] + \dots \end{aligned}$$

Setting  $\hbar = -1$ , we get an accurate approximation solution in the following form:

$$\begin{aligned} u_j(t) &= e^{ijkh} \left\{ 1 - \frac{i\omega t^\alpha}{\Gamma(\alpha+1)} - \frac{\omega^2 t^{2\alpha}}{\Gamma(2\alpha+1)} \right. \\ &\quad \left. + \frac{i\omega^3 t^{3\alpha}}{\Gamma(3\alpha+1)} \left[ 1 - \frac{2q}{\omega} + \frac{q\Gamma(2\alpha+1)}{\omega(\Gamma(\alpha+1))^2} \right] + \dots \right\} \end{aligned} \quad (26)$$

For the special case  $\alpha = 1$ , the form Equation (26) is obtained discrete plane wave solution

$$\begin{aligned} u_j(t) &= e^{ijkh} \left\{ 1 - i\omega t - \frac{\omega^2 t^2}{2} + \frac{i\omega^3 t^3}{6} + \dots \right\} \\ &= e^{ijkh} \left\{ 1 + \frac{(-i)\omega t}{1!} + \frac{(-i)^2 \omega^2 t^2}{2!} + \frac{(-i)^3 \omega^3 t^3}{3!} + \dots \right\} \\ &= e^{ijkh} e^{-i\omega t} \\ &= e^{i(jkh - \omega t)}, \end{aligned}$$

which is the same solution obtained in [6].

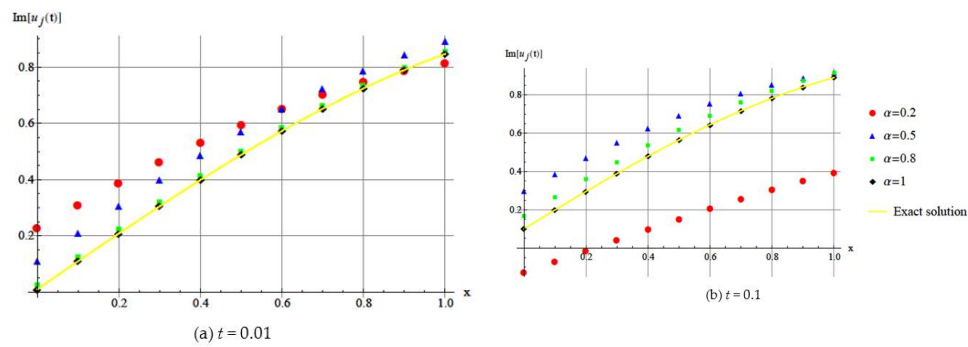
$$u(x, t) = e^{i(kx - \omega t)}, \quad x \in \mathbb{R}, \quad t > 0,$$

is the plane wave solution of the continuous form, where  $k$  is the wave number and  $\omega$  denotes the frequency.

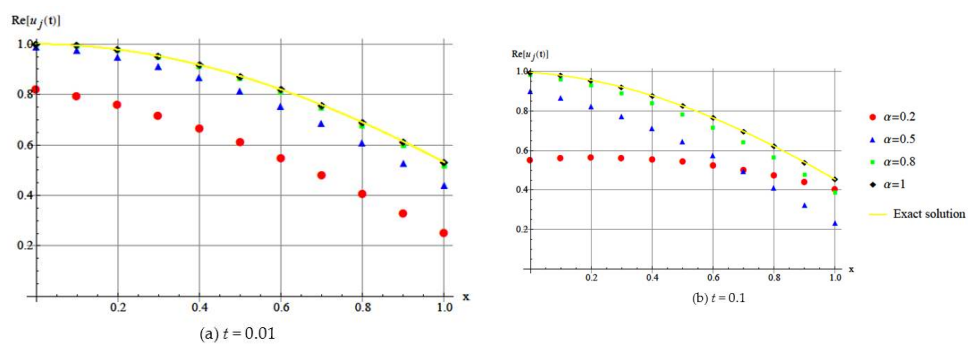
Figures 3 and 4 show the DHAM approximate solution of  $u(x, t)$  for different values of  $\alpha$ ,  $k = 1$  and  $q = 2$ .

We can observe the different behaviors of the discrete fractional Schrödinger equations, with different fractional parameters.

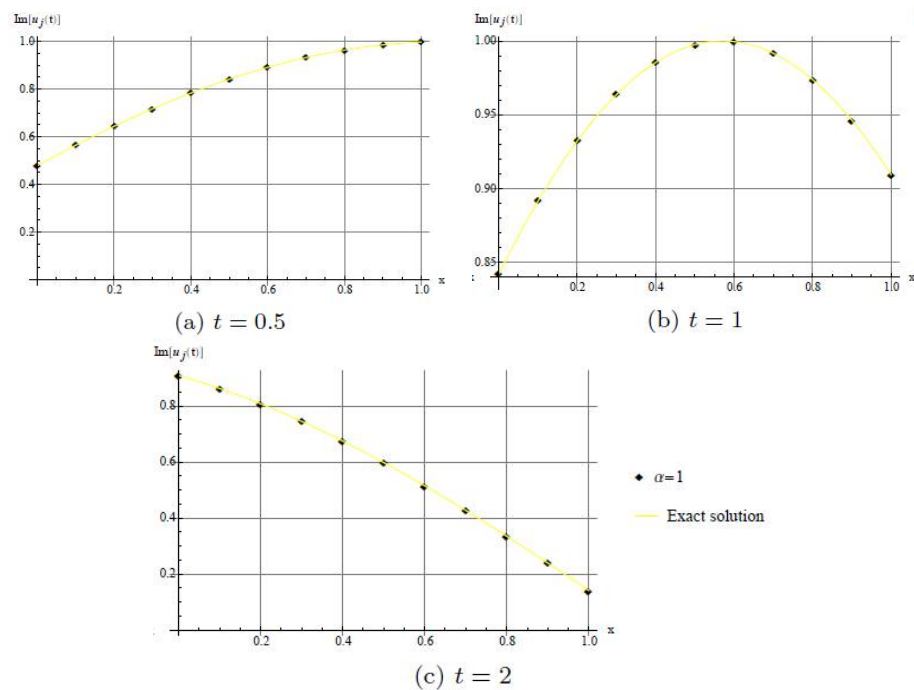
In Figure 5, we show that the method has good agreement with imaginary part of the exact solution when  $\alpha = 1$ .



**Figure 3.** Numerical illustration of imaginary part of approximation solution  $u(x,t)$  by DHAM. (a) For  $t = 0.01$ ; (b) For  $t = 0.1$ .

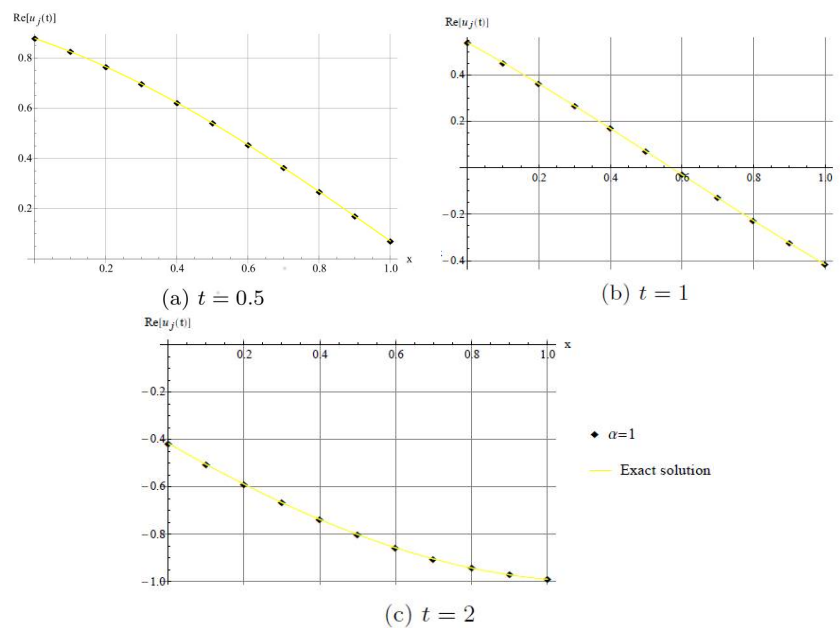


**Figure 4.** Numerical illustration of real part of approximation solution  $u(x,t)$  by DHAM. (a) For  $t = 0.01$ ; (b) For  $t = 0.1$ .



**Figure 5.** Comparison with numerical solution of  $u(x,t)$  by DHAM and the exact solution when  $\alpha = 1$ . (a) For  $t = 0.5$ ; (b) For  $t = 1$ ; (c) For  $t = 2$ .

In Figure 6, we show that the method has good agreement with real part of the exact solution when  $\alpha = 1$ .



**Figure 6.** Comparison with numerical solution of  $u(x,t)$  by DHAM and the exact solution when  $\alpha = 1$ . (a) For  $t = 0.5$ ; (b) For  $t = 1$ ; (c) For  $t = 2$ .

**Example 3.** Consider time fractional space discrete nonlinear Burgers' equation

$$D_t^\alpha u_j(t) + u_j(t)D_h u_j(t) = D_h^2 u_j(t), \quad j \in \mathbb{Z}, \quad t > 0, \quad 0 < \alpha \leq 1, \quad (27)$$

with initial condition

$$u_j(0) = \sin jh. \quad (28)$$

Initial value problem (27) and (28) is the discrete form of initial value problem for nonlinear fractional Burgers' equation

$$D_t^\alpha u(x,t) + u(x,t)u_x(x,t) = u_{xx}(x,t), \quad t > 0, \quad x \in \mathbb{R}, \quad 0 < \alpha \leq 1,$$

with initial condition

$$u(x,0) = \sin x.$$

We take into consideration the linear operator:

$$\mathcal{L}[\varphi_j(t;p)] = D_t^\alpha [\varphi_j(t;p)] = \frac{\partial^\alpha \varphi_j(t;p)}{\partial t^\alpha} \quad (29)$$

with property  $\mathcal{L}[c] = 0$ , where  $c$  is constant. We can consider the nonlinear operator as

$$\mathcal{N}[\varphi_j(t;p)] = D_t^\alpha \varphi_j(t;p) - D_h^2 \varphi_j(t;p) + \varphi_j(t;p)D_h \varphi_j(t;p). \quad (30)$$

Therefore, we construct the zeroth-order deformation equation by Equation (2).

For  $p = 0$  and  $p = 1$ , we can write

$$\varphi_j(t;0) = u_{j,0}(t) = u_j(0), \quad \varphi_j(t;1) = u_j(t),$$

respectively. Thus, we obtain the  $m$ th-order deformation equation

$$\mathcal{L}[u_{j,m}(t) - \mathcal{X}_m u_{j,m-1}(t)] = \hbar H_j(t) \mathcal{R}_m[\vec{u}_{m-1}], \quad (31)$$

where

$$\mathcal{R}_m[\vec{u}_{m-1}] = D_t^\alpha u_{j,m-1}(t) - D_h^2 u_{j,m-1}(t) + \sum_{k=0}^{m-1} u_{j,k}(t) D_h u_{j,m-1-k}(t). \quad (32)$$

For simplicity, we select again  $H_j(t) = 1$ . So, the approximations of  $u_j(t)$  are only depend on auxiliary parameter  $\hbar$ .

The solution of the  $m$ th-order deformation equation for  $m \geq 1$  give rise to

$$u_{j,m}(t) = \mathcal{X}_m u_{j,m-1}(t) + \hbar J^\alpha \mathcal{R}_m[\vec{u}_{m-1}]. \quad (33)$$

when we use the initial condition (28) along with (33), we attain the first three of terms of (33) as following:

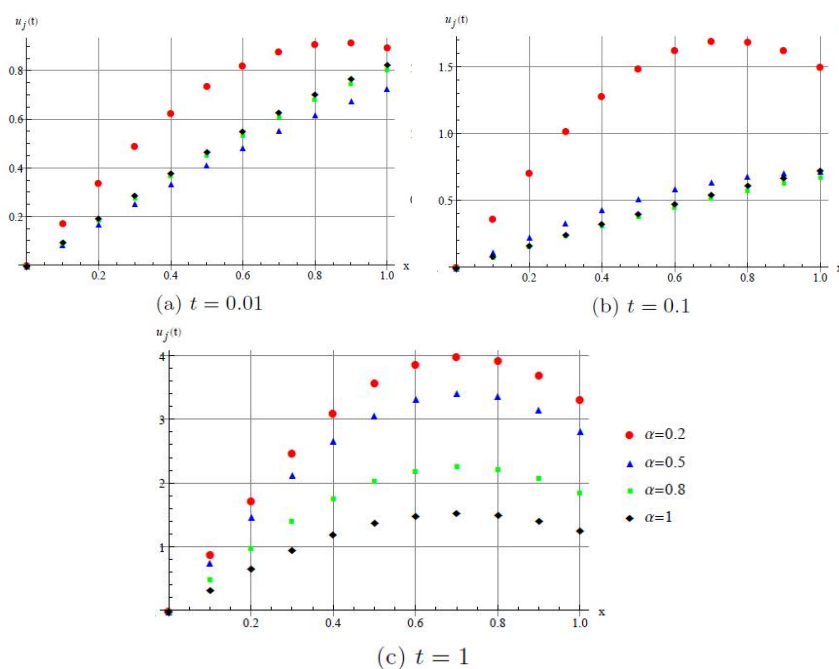
$$\begin{aligned} u_{j,0}(t) &= \sin jh \\ u_{j,1}(t) &= \left[ \frac{\sin h}{2h} \sin 2jh - \frac{2(\cos h - 1)}{h^2} \sin jh \right] \frac{t^\alpha}{\Gamma(\alpha+1)} \hbar \\ u_{j,2}(t) &= \left[ \frac{\sin h}{2h} \sin 2jh - \frac{2(\cos h - 1)}{h^2} \sin jh \right] \frac{t^{2\alpha}}{\Gamma(\alpha+1)} \hbar (\hbar + 1) \\ &\quad + \left[ \frac{\sin^2 h}{2h^2} \sin 2jh \cos jh + \frac{\sin h \sin 2h}{2h^2} \sin jh \cos 2jh \right. \\ &\quad \left. - \frac{2 \sin h (\cos h - 1) (\cos h + 2)}{h^3} \sin 2jh + \frac{4(\cos h - 1)^2}{h^4} \sin jh \right] \frac{t^{2\alpha}}{\Gamma(2\alpha+1)} \hbar^2 \\ &\quad \vdots \end{aligned}$$

and so on.

Thus, we can conclude that

$$u_j(t) = u_{j,0}(t) + u_{j,1}(t) + u_{j,2}(t) + \dots$$

Figure 7 shows the DHAM approximate solution of  $u(x, t)$  for different values of  $\alpha$ .



**Figure 7.** Numerical illustration of approximation solution  $u(x, t)$  by DHAM. (a) For  $t = 0.01$ ; (b) For  $t = 0.1$ ; (c) For  $t = 1$ .

We can see that the different behaviors of the discrete fractional Burgers' equations for different fractional parameters.

#### 4. Discussion and Conclusions

In this paper, the discrete HAM is successfully applied to find the solutions of linear and nonlinear fractional partial differential equations with time derivative  $\alpha$  ( $0 < \alpha \leq 1$ ). In contrast to all other analytic methods, it provides us with a simple way to adjust and convergence region of solution series by introducing an auxiliary parameter  $\hbar$ . This is an obvious advantage of the DHAM. We can simply choose the fractional operator  $D_t^\alpha$  as the auxiliary linear operator. In this way, we obtained solutions in power series. Also we obtained the exact solutions in special case  $\alpha = 1$ ,  $\hbar = -1$  for some equations. However, it is well-known that a power series often has a small convergence radius. The results of test problems show that the DHAM is effective and reliable. It may also be a promising method to solve other nonlinear partial differential equations.

**Funding:** This research received no external funding.

**Acknowledgments:** The author would like to thank the reviewers for their valuable comments and suggestions to improve this paper.

**Conflicts of Interest:** The author declares no conflict of interest.

#### References

1. Ablowitz, M.J.; Clarkson, P.A. *Solutions Nonlinear Evolution Equations and Inverse Scattering*; Cambridge University Press: New York, NY, USA, 1991.
2. Herrmann, R. *Fractional Calculus—An Introduction for Physicists*; World Scientific: Singapore, 2011.
3. Luchko, Y. Fractional Schrödinger equation for a particle moving in a potential well. *J. Math. Phys.* **2013**, *54*, 012111. [\[CrossRef\]](#)
4. Mickens, R.E. *Nonstandard Finite Difference Models of Differential Equations*; World Scientific: Singapore, 1994.
5. Onorato, M.; Osborne, A.R.; Serio, M.; Bertone, S. Freak waves in random oceanic sea states. *Phys. Rev. Lett.* **2012**, *86*, 5831–5834. [\[CrossRef\]](#) [\[PubMed\]](#)
6. Samko, S.G.; Kilbas, A.A.; Marichev, O.I. *Fractional Integrals and Derivatives, Theory and Applications*; Gordon and Breach: Amsterdam, The Netherlands, 1993.
7. Babiarez, A.; Niezabitowski, M. Controllability Problem of Fractional Neutral Systems: A Survey. *Math. Probl. Eng.* **2017**, *2017*, 4715861. [\[CrossRef\]](#)
8. Babiarez, A.; Łęgowski, A.; Niezabitowski, M. Robot Path Control with Al-Alaoui Rule for Fractional Calculus Discretization. In *Theory and Applications of Non-Integer Order Systems*; Lecture Notes in Electrical Engineering; Babiarez, A., Czornik, A., Klamka, J., Niezabitowski, M., Eds.; Springer: Cham, Switzerland, 2017; Volume 407.
9. Bayin, S.Ş. On the consistency of the solutions of the space fractional Schrodinger equation. *J. Math. Phys.* **2012**, *53*, 042105. [\[CrossRef\]](#)
10. El-Nabulsi, R.A.; Torres, D.F.M. Fractional action-like variational approach. *J. Math. Phys.* **2008**, *49*, 053521. [\[CrossRef\]](#)
11. El-Nabulsi, R.A. Fractional calculus of variations from extended Erdelyi-Kober operator. *Int. J. Mod. Phys B.* **2009**, *23*, 3349–3361. [\[CrossRef\]](#)
12. El-Nabulsi, R.A. Fractional variational problems from extended exponentially fractional integral. *Appl. Math. Comp.* **2011**, *217*, 9492–9496.
13. Herrmann, R. The fractional Schrodinger equation and the infinite potential well—Numerical results using the Riesz derivative. *Gam. Ori. Chron. Phys.* **2013**, *1*, 1–12.
14. Luchko, Y.; Gorenflo, R. An operational method for solving fractional differential equations with the Caputo derivative. *Acta Math. Vietnam.* **1999**, *24*, 207–233.
15. Malinowska, A.B.; Torres, D.F.M. *Introduction to the Fractional Calculus of Variations*; Imperial College Press: London, UK, 2012.
16. Ubriaco, M.R. Entropies based on fractional calculus. *Phys. Lett. A* **2009**, *373*, 2516–2519. [\[CrossRef\]](#)

17. Das, S.; Vishal, K.; Gupta, P.K.; Saha Ray, S. Homotopy analysis method for solving fractional diffusion equation. *Int. J. Appl. Math. Mech.* **2011**, *7*, 28–37.
18. Dehghan, M.; Manafian, J.; Saadatmandi, A. The solution of the linear fractional partial differential equations using the homotopy analysis method. *Z. Naturforsch.* **2010**, *65*, 935–949.
19. Dehghan, M.; Manafian, J.; Saadatmandi, A. Solving nonlinear fractional partial differential equations using the homotopy analysis method. *Numer. Methods Part. Differ. Equ.* **2010**, *26*, 448–479. [[CrossRef](#)]
20. Kivshar, Y.S.; Agarwal, G.P. *Optical Solitons: From Fibers to Photonic Crystals*; Academic Press: San Diego, CA, USA, 2003.
21. Liao, S.J. The Proposed Homotopy Analysis Technique for the Solution of Nonlinear Problems. Ph.D. Thesis, Shanghai Jiao Tong University, Shanghai, China, 1992.
22. Scott, A. *Nonlinear Science: Emergence & Dynamics of Coherent Structures*; Oxford University Press: Oxford, UK, 1999.
23. Gepreel, K.A. The homotopy perturbation method applied to the nonlinear fractional Kolmogorov-Petrovskii-Piskunov equations. *Appl. Math. Lett.* **2011**, *24*, 1428–1434. [[CrossRef](#)]
24. He, J.H. Homotopy perturbation method: A new nonlinear analytic technique. *Appl. Math. Comput.* **2003**, *135*, 73–79. [[CrossRef](#)]
25. He, J.H. An elementary introduction to the homotopy perturbation method. *Comput. Math. Appl.* **2009**, *57*, 410–412. [[CrossRef](#)]
26. Sezer, A.S.; Yildirim, A.; Mohyud-Din, S.T. He's homotopy perturbation method for solving the fractional KdV-Burgers-Kuramoto equation. *Int. J. Numer. Method Heat Fluid Flow* **2011**, *21*, 448–458. [[CrossRef](#)]
27. Ahmad, J.; Mohyud-Din, S.T. Solving fractional vibrational problem using restarted fractional Adomian's decomposition method. *Life Sci. J.* **2013**, *10*, 210–216.
28. Daftardar-Gejji, V.; Jafari, H. Solving a multi-order fractional differential equation using Adomian decomposition method. *Appl. Math. Comput.* **2007**, *189*, 541–548.
29. Daftardar-Gejji, V.; Bhalekar, S. Solving a multi-term linear and non-linear diffusion wave equations of fractional order by Adomian decomposition method. *Appl. Math. Comput.* **2008**, *202*, 113–120. [[CrossRef](#)]
30. Dehghan, M.; Abbaszadeh, M.; Mohebbi, A. An implicit RBF meshless approach for solving time fractional nonlinear Sine-Gordon and Klein-Gordon equation. *Eng. Anal. Bound. Elem.* **2015**, *50*, 412–434. [[CrossRef](#)]
31. Miller, K.S.; Ross, B. *An Introduction to the Fractional Calculus and Fractional Differential Equations*; Wiley: New York, NY, USA, 1993.
32. Mohebbi, A.; Abbaszadeh, M.; Denghan, M. The use of a meshless technique based on collocation and radial basis functions for solving the time fractional Schrödinger equation arising in quantum mechanics. *Eng. Anal. Bound. Elem.* **2013**, *37*, 475–485. [[CrossRef](#)]
33. Shirzadi, A.; Ling, L.; Abbasbandy, S. Meshless simulations of the two dimensional fractional-time convection-diffusion-reaction equations. *Eng. Anal. Bound. Elem.* **2012**, *38*, 1522–1527. [[CrossRef](#)]
34. Dehghan, M.; Safarpour, M. The dual reciprocity boundary elements method for the linear and nonlinear two-dimensional time-fractional partial differential equations. *Math. Methods Appl. Sci.* **2016**, *39*, 3979–3995. [[CrossRef](#)]
35. Podlubny, I. *Fractional Differential Equations*; Academic Press: San Diego, CA, USA, 1999.
36. Bratsos, A.; Ehrhardt, M.; Famelis, I.T. A Discrete Adomian decomposition method for discrete nonlinear Schrödinger equations. *Appl. Math. Comput.* **2008**, *197*, 190–205. [[CrossRef](#)]
37. Zhu, H.; Shu, H.; Ding, M. Numerical solutions of partial differential equations by discrete homotopy analysis method. *Appl. Math. Comput.* **2010**, *216*, 3592–3605. [[CrossRef](#)]
38. Lu, Z.; Takeuchi, Y. Global asymptotic behavior in single-species discrete diffusion systems. *J. Math. Biol.* **1993**, *32*, 67–77. [[CrossRef](#)]
39. Caputo, M. Linear models of dissipation whose Q is almost independent II. *Geophys. J. R. Astron.* **1967**, *13*, 529–539. [[CrossRef](#)]
40. Liao, S.J. *Beyond Perturbation: Introduction to Homotopy Analysis Method*; Chapman and Hall/CRC Press: Boca Raton, FL, USA, 2003.
41. Fath, G. Propagation failure of traveling waves in a discrete bistable medium. *Physica D* **1998**, *116*, 176–190. [[CrossRef](#)]
42. Prehl, J.; Essex, C.; Hoffmann, K.H. Tsallis relative entropy and anomalous diffusion. *Entropy* **2012**, *14*, 701–716. [[CrossRef](#)]

43. Prehl, J.; Boldt, F.; Hoffmann, K.H.; Essex, C. Symmetric fractional diffusion and entropy production. *Entropy* **2016**, *18*, 275. [[CrossRef](#)]
44. Eisenberg, H.S.; Silberberg, Y.; Morandotti, R.; Boyd, A.R.; Aitchison, J.S. Discrete Spatial Optical Solitons in Waveguide Arrays. *Phys. Rev. Lett.* **1998**, *81*, 3383–3386. [[CrossRef](#)]
45. Efremidis, N.K.; Sears, S.; Christodoulides, D.N.; Fleischer, J.W.; Segev, M. Discrete solitons in photorefractive optically induced photonic lattices. *Phys. Rev. E* **2002**, *66*, 046602. [[CrossRef](#)] [[PubMed](#)]
46. Gatz, S.; Herrmann, J. Soliton propagation in materials with saturable nonlinearity. *J. Opt. Soc. Am. B* **1991**, *8*, 2296–2302. [[CrossRef](#)]
47. Gatz, S.; Herrmann, J. Soliton propagation and soliton collision in doubledoped fibers with a non-Kerr-like nonlinear refractive-index change. *Opt. Lett.* **1992**, *17*, 484–486. [[CrossRef](#)] [[PubMed](#)]



© 2018 by the author. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (<http://creativecommons.org/licenses/by/4.0/>).