## Supplementary Materials

## S1. Linear Response Functions of a Harmonic Oscillator System

We also computed linear response functions (LRFs) of a harmonic oscillator system, of which the potential is given by, $U(r)=x^{2} / 2$. However we did not present the results in the text, because these are essentially similar to those of the infinite square well potential systems. Still, as there are still differences in LRFs between these two systems, we plotted the LRFs of the harmonic oscillator system in Figure S1. As we can see from these plots, the spatial region over which LRF distributes becomes larger as Nocc increases. This is because the higher energies of electrons become, the broader the region that the electrons in exist becomes. However we should also note that, $\delta \rho(\mathrm{r}) / \delta \mathrm{v}\left(\mathrm{r}^{\prime}\right)$ decays rapidly for the distance, $\left|\mathrm{r}-\mathrm{r}^{\prime}\right|$ as Nocc increases. This conclusion is similar to that from the infinite square well potential described in the text.

Figure S1. LRFs of harmonic oscillator systems for various numbers of occupied orbitals (Nocc). (a) Nocc $=1$, (b) Nocc $=2$, (c) Nocc $=5$, (d) Nocc $=10$, (e) Nocc $=20$, and (f) $\mathrm{Nocc}=50$.


## S2. Solutions of Double-well Potential Systems

We describe how to obtain eigenvalues and analytical forms of eignfunctions.
For $E_{k}<V_{0}$, the symmetric function is given by,

$$
\Psi_{\mathrm{k}_{\mathrm{n}}^{+}}(\mathrm{r})=\left\{\begin{array}{cc}
\tan \left(2 \mathrm{k}_{\mathrm{n}}^{+}\right) \cos \left(\mathrm{k}_{\mathrm{n}}^{+} \mathrm{x}\right)+\sin \left(\mathrm{k}_{\mathrm{n}}^{+} \mathrm{x}\right) & -2 \leq \mathrm{x} \leq-1  \tag{S1}\\
\frac{\sin \left(\mathrm{k}_{\mathrm{n}}^{+} \mathrm{x}\right)}{\cosh \left(\mathrm{k}_{\mathrm{n}}^{\prime+}\right) \cos \left(2 \mathrm{k}_{\mathrm{n}}^{+}\right)} \cosh \left(\mathrm{k}_{\mathrm{n}}^{\prime+} \mathrm{x}\right) & -1 \leq \mathrm{x} \leq-1 \\
\tan \left(2 \mathrm{k}_{\mathrm{n}}^{+}\right) \cos \left(\mathrm{k}_{\mathrm{n}}^{+} \mathrm{x}\right)-\sin \left(\mathrm{k}_{\mathrm{n}}^{+} \mathrm{x}\right) & 1 \leq \mathrm{x} \leq 2
\end{array}\right.
$$

and the asymmetric functions by

$$
\psi_{\mathrm{k}_{\mathrm{n}}^{-}}(\mathrm{r})=\left\{\begin{array}{lc}
\tan \left(2 \mathrm{k}_{\mathrm{n}}^{-}\right) \cos \left(\mathrm{k}_{\mathrm{n}}^{-} \mathrm{x}\right)+\sin \left(\mathrm{k}_{\mathrm{n}}^{-} \mathrm{x}\right) & -2 \leq \mathrm{x} \leq-1  \tag{S2}\\
-\frac{\sin \left(\mathrm{k}_{\mathrm{n}}^{-} \mathrm{x}\right)}{\sinh \left(\mathrm{k}_{\mathrm{n}}^{\prime-}\right) \cos \left(2 \mathrm{k}_{\mathrm{n}}^{-}\right)} \sinh \left(\mathrm{k}_{\mathrm{n}}^{\prime-} \mathrm{x}\right) & -1 \leq \mathrm{x} \leq-1 \\
-\tan \left(2 \mathrm{k}_{\mathrm{n}}^{-}\right) \cos \left(\mathrm{k}_{\mathrm{n}}^{-} \mathrm{x}\right)+\sin \left(\mathrm{k}_{\mathrm{n}}^{-} \mathrm{x}\right) & 1 \leq \mathrm{x} \leq 2
\end{array}\right.
$$

here $k_{n}^{\prime \pm}=\sqrt{V_{0}-E_{k_{n}^{ \pm}}}$, where $E_{k_{n}^{ \pm}}$is the eigenvalue of the $n$-th symmetric $(+)$ or antisymmetric $(-)$ eigenfunction,

$$
\begin{equation*}
\mathrm{E}_{\mathrm{k}_{\mathrm{n}}^{ \pm}}=\left(\mathrm{k}_{\mathrm{n}}^{ \pm}\right)^{2} \tag{S3}
\end{equation*}
$$

From the continuously differentiable conditions,

$$
\begin{equation*}
\mathrm{k}_{\mathrm{n}}^{+} \cot _{\mathrm{n}}^{+}=-\mathrm{k}_{\mathrm{n}}^{\prime+} \tanh \left(\mathrm{k}_{\mathrm{n}}^{\prime+}\right) \tag{S4}
\end{equation*}
$$

and

$$
\begin{equation*}
\mathrm{k}_{\mathrm{n}}^{-} \cot \left(\mathrm{k}_{\mathrm{n}}^{-}\right)=-\mathrm{k}_{\mathrm{n}}^{\prime-} \operatorname{coth} \mathrm{k}_{\mathrm{n}}^{\prime-} \tag{S5}
\end{equation*}
$$

we determined the $k_{n}^{ \pm}$values for symmetric $(+)$and antisymmetric ( - ) cases.
For $E_{k}<V_{0}$, the equations are slightly changed.

$$
\begin{align*}
& \Psi_{\mathrm{k}_{\mathrm{n}}^{+}}(\mathrm{r})=\left\{\begin{array}{lc}
\tan \left(2 \mathrm{k}_{\mathrm{n}}^{+}\right) \cos \left(\mathrm{k}_{\mathrm{n}}^{+} \mathrm{x}\right)+\sin \left(\mathrm{k}_{\mathrm{n}}^{+} \mathrm{x}\right) & -2 \leq \mathrm{x} \leq-1 \\
\frac{\sin \left(\mathrm{k}_{\mathrm{n}}^{+} \mathrm{x}\right)}{\cos \left(\mathrm{k}_{\mathrm{n}}^{\prime \prime+}\right) \cos \left(2 \mathrm{k}_{\mathrm{n}}^{+}\right)} \cos \left(\mathrm{k}_{\mathrm{n}}^{\prime \prime+} \mathrm{x}\right) & -1 \leq \mathrm{x} \leq-1 \\
\tan \left(2 \mathrm{k}_{\mathrm{n}}^{+}\right) \cos \left(\mathrm{k}_{\mathrm{n}}^{+} \mathrm{x}\right)-\sin \left(\mathrm{k}_{\mathrm{n}}^{+} \mathrm{x}\right) & 1 \leq \mathrm{x} \leq 2
\end{array}\right.  \tag{S6}\\
& \Psi_{\mathrm{k}_{\mathrm{n}}^{-}}(\mathrm{r})= \begin{cases}\tan \left(2 \mathrm{k}_{\mathrm{n}}^{-}\right) \cos \left(\mathrm{k}_{\mathrm{n}}^{-} \mathrm{x}\right)+\sin \left(\mathrm{k}_{\mathrm{n}}^{-} \mathrm{x}\right) & -2 \leq \mathrm{x} \leq-1 \\
-\frac{\sin \left(\mathrm{k}_{\mathrm{n}}^{-} \mathrm{x}\right)}{\sin \left(\mathrm{k}_{\mathrm{n}}^{\prime \prime-}\right) \cos \left(2 \mathrm{k}_{\mathrm{n}}^{-}\right)} \sin \left(\mathrm{k}_{\mathrm{n}}^{\prime \prime-} \mathrm{x}\right) & -1 \leq \mathrm{x} \leq-1 \\
-\tan \left(2 \mathrm{k}_{\mathrm{n}}^{-}\right) \cos \left(\mathrm{k}_{\mathrm{n}}^{-} \mathrm{x}\right)+\sin \left(\mathrm{k}_{\mathrm{n}}^{-} \mathrm{x}\right) & 1 \leq \mathrm{x} \leq 2\end{cases} \tag{S6}
\end{align*}
$$

From the continuously differentiable conditions,

$$
\begin{equation*}
\mathrm{k}_{\mathrm{n}}^{+} \cot \mathrm{k}_{\mathrm{n}}^{+}=-\mathrm{k}_{\mathrm{n}}^{\prime \prime+} \tan \left(\mathrm{k}_{\mathrm{n}}^{\prime \prime+}\right) \tag{S4}
\end{equation*}
$$

and

$$
\begin{equation*}
\mathrm{k}_{\mathrm{n}}^{-} \cot \left(\mathrm{k}_{\mathrm{n}}^{-}\right)=-\mathrm{k}_{\mathrm{n}}^{\prime \prime-} \cot \mathrm{k}_{\mathrm{n}}^{\prime \prime-} \tag{S5}
\end{equation*}
$$

we determined the $\mathrm{k}_{\mathrm{n}}^{ \pm}$values for symmetric ( + ) and antisymmetric ( - ) cases.
The energies of the lowest-lying 20 states are listed in Tables S 1 and S 2 for $\mathrm{V}_{0}=25.0$ and $\mathrm{V}_{0}=100.0$ respectively.

Table S1. Energy levels of lowest-lying 20 symmetric ( $\mathrm{E}_{\mathrm{k}_{\mathrm{n}}^{+}}$) and antisymmetric ( $\mathrm{E}_{\mathrm{k}_{\mathrm{n}}^{-}}$) states for $V_{0}=25.0$.

| Principal Quantum <br> Number (n) | $\mathrm{E}_{\mathrm{k}_{\mathrm{n}}^{+}}$ | $\mathrm{E}_{\mathrm{k}_{\mathrm{n}}^{-}}$ |
| :---: | :---: | :---: |
| 0 | 6.737136611241714 | 6.738586329623146 |
| 1 | 23.85268979853237 | 24.45828371131453 |
| 2 | 29.02476750506050 | 35.92020877590885 |
| 3 | 44.69052075564606 | 53.850993048370069 |
| 4 | 63.07266603809249 | 74.021976028317598 |
| 5 | 87.57397077271618 | 102.40275031296936 |
| 6 | 117.3299647023210 | 133.20897457515955 |
| 7 | 151.3144849061653 | 171.07637476244874 |
| 8 | 191.2659933651678 | 212.21266805659995 |
| 9 | 235.0781885644307 | 259.68213865408501 |
| 10 | 284.9495480989314 | 310.94680417402724 |
| 11 | 338.6670847478374 | 368.11961813658054 |
| 12 | 398.3885518478544 | 429.40854140215914 |
| 13 | 462.0256928221838 | 496.34429193922553 |
| 14 | 531.5803957116293 | 567.60137168748467 |
| 15 | 605.1353735676119 | 644.33513174499887 |
| 16 | 684.5216795438358 | 725.52808432128687 |
| 17 | 767.9889987187931, | 812.08137221390233 |
| 18 | 857.2097266149157 | 903.19047882543600 |
| 19 | 950.5835962825196 | 999.57709334509814 |
| 20 | 1049.642608083979 | 1100.5896890165435 |

Table S2. Energy levels of lowest-lying 20 symmetric ( $\mathrm{E}_{\mathrm{k}_{\mathrm{n}}^{+}}$) and antisymmetric ( $\mathrm{E}_{\mathrm{k}_{\mathrm{n}}^{-}}$) states for $V_{0}=100.0$.

| Principal Quantum <br> Number (n) | $\mathrm{E}_{\mathrm{k}_{\mathrm{n}}^{+}}$ | $\mathrm{E}_{\mathrm{k}_{\mathbf{n}}}$ |
| :---: | :---: | :---: |
| 0 | 8.135854269472983 | 8.13585429619729 |
| 1 | 32.25340052217796 | 32.25340186525873 |
| 2 | 70.95009585920174 | 70.95063314337901 |
| 3 | 102.1164019712885 | 107.72888719571438 |
| 4 | 114.6772275841585 | 121.55218645459998 |
| 5 | 130.5279861984107 | 143.8814437261952 |
| 6 | 160.6432406009162 | 178.1895246847984 |
| 7 | 194.2447029736519 | 210.3921198528759 |
| 8 | 230.2683322281207 | 253.6459425343700 |
| 9 | 277.1063977760075 | 299.2070863045833 |
| 10 | 322.8249575631072 | 350.2985196899664 |
| 11 | 379.5046166941854, | 407.6867628400308 |
| 12 | 436.0974830259074 | 467.7111367422718 |
| 13 | 501.9033978777250 | 535.7686163834129 |
| 14 | 569.3337305404052 | 605.3886595419620 |
| 15 | 644.3223149282313 | 683.5513370682773 |
| 16 | 722.3419381761522 | 763.0618747369907 |
| 17 | 806.6915148276904 | 851.0854959366932 |
| 18 | 895.0797346089407 | 940.6004134960433 |
| 19 | 988.9455576312383 | 1038.388405913195 |
| 20 | 1087.542080170761 | 1137.941451482828 |

## S3. Geometry of the $\alpha$-Helix Structure of Trialanine Peptide

We shall describe briefly how to determine the geometry of $\alpha$-helix structure of the trialanine peptide. We consider the trialanine capped with $\mathrm{CH}_{3} \mathrm{CO}$ - and $\mathrm{CH}_{3} \mathrm{NH}$ - for N and C terminals, respectively. Choosing the interatomic distance between the O atom of the carobonyl group in a terminal amide and the H atom in the other terminal amide plane, $R_{O-H}$, as the reaction coordinate, the umbrella sampling calculation was performed. The number of sampling windows is 13 , each for which we imposed a harmonic restraining potential, $V(r)=k\left(r-R_{\text {center }}\right)^{2}$. For 13 windows, $R_{\text {center }}$ ranges from 1.8 to $13.8 \AA$ with spacing $1.0 \AA$ and k is fixed to be $1.2 \mathrm{kcal} \cdot \mathrm{mol}^{-1} \cdot \AA^{-2}$. The ff 12 SB and the Velocity Verlet method [1] were used for the force field and the integration of molecular dynamics runs, respectively. The solvent effect is included using a modified generalized Born (GB) model [2]. The dielectric constant is 1.0 , which is more lipophilic than the usual constant (4.0). This is for enhancing the stability of the $\alpha$-helix structure of trialanine peptide, of which there is only one hydrogen bond of the $\alpha$-helix structure, resulting in fragileness of the structure as described below. Temperature was kept constant at 300 K , which was controlled with using the Berendsen thermostat [3]. In addition, the SHAKE algorithm [4] was employed to fix all bonds involving hydrogen atoms. For each window, we first minimize the geometry with the specific constraint, and then followed by 200 ps of heating, 500 ps of equilibration, and 1 ns of production MD runs. The data were collected for every 100 steps during production runs. All MD calculations and the preparations of initial geometries were carried out with the Amber 12 and Amber12 tools packages, respectively. Using the histogram constructed by the umbrella sampling runs, the wham equation is solved using WHAM [5]. We found that the potential mean force (PMF) has a local minimum at approximately $2.0 \AA$ and a global minimum at approximately $4.0 \AA$ for $R_{O-H}$. Although the details of the PMF profile we obtained are different from those in the previous study [6]:(i) The $R_{O-H}$ value of the global minimum we obtained is much smaller than that of the previous study and (ii) the energy difference between the $\alpha$-helix structure and the global minimum structure is approximately $2.0 \mathrm{kcal} / \mathrm{mol}$, which is much larger than at most $0.5 \mathrm{kcal} / \mathrm{mol}$ in [6], our result is consistent with the previous report [6] that the $\alpha$-helix structure is not the global minimum, but a local minimum. The $R_{O-H}$ value of the local minimum is also similar to the value reported in [6]. The $\alpha$-helix structure we obtained is shown in Figure 8 in the text.

## S4. Linear Response Functions of Molecules for the Perturbations that are Applied to all Atoms

Figure S2. Isosurfaces of linear response functions of tripeptide molecule for the perturbations on the atomic sites. $\delta \rho(\mathbf{r}) / \delta v_{X}=0.01$ and $=-0.01$ are described as blue and red surfaces, respectively. For carbons and hydrogens, we indicated the atomic types as follows. $\mathrm{H}\left(\mathrm{C}_{\mathrm{T}}\right)=\mathrm{H}$ of the terminal methyl group. $\mathrm{H}\left(\mathrm{C}_{\alpha}\right)=\mathrm{H}$ in the amide plane. $\mathrm{H}\left(\mathrm{C}_{\beta}\right)=\mathrm{H}$ of the methyl residue. $\mathrm{C}_{\alpha}, \mathrm{C}_{\beta}$, and $\mathrm{C}(0)$ are $\alpha$ carbon, $\beta$ carbon and carbon of the carbonyl group, respectively.

$1 \mathrm{H}\left(\mathrm{C}_{\mathrm{T}}\right)$

$5 \mathrm{C}(\mathrm{O})$

$9 \mathrm{C}_{\alpha}$

$13 \mathrm{H}\left(\mathrm{C}_{\beta}\right)$

$2 \mathrm{C}_{\mathrm{T}}$


60

$10 \mathrm{H}\left(\mathrm{C}_{\alpha}\right)$

$14 \mathrm{H}\left(\mathrm{C}_{\beta}\right)$

$3 \mathrm{H}\left(\mathrm{C}_{\mathrm{T}}\right)$


7 N

$11 \mathrm{C}_{\beta}$

$12 \mathrm{H}\left(\mathrm{C}_{\beta}\right)$

$15 \mathrm{C}(\mathrm{O})$


16 O

Figure S2. Cont.


17 N

$21 \mathrm{C}_{\beta}$


25 C (O)

$29 \mathrm{C}_{\alpha}$

$33 \mathrm{H}\left(\mathrm{C}_{\beta}\right)$

$18 \mathrm{H}(\mathrm{N})$

$22 \mathrm{H}\left(\mathrm{C}_{\beta}\right)$


26 O

$30 \mathrm{H}\left(\mathrm{C}_{\alpha}\right)$

$34 \mathrm{H}\left(\mathrm{C}_{\beta}\right)$


27 N

$31 C_{\beta}$


35 C (O)


$24 \mathrm{H}\left(\mathrm{C}_{\beta}\right)$

$28 \mathrm{H}(\mathrm{N})$

$32 \mathrm{H}\left(\mathrm{C}_{\mathrm{B}}\right)$


36 O

Figure S2. Cont.


Figure S3. Isosurfaces of linear response functions of butane molecule for the perturbations on the atomic sites. $\delta \rho(\mathbf{r}) / \delta \mathrm{v}_{\mathrm{x}}=0.035$ and $=-0.035$ are described as blue and red surfaces, respectively.

1 C

4 C

2 C

5 H

7 H


8 H

3 C

6 H


Figure S3. Cont.


10 H


13 H


11 H


14 H

Figure S4. Isosurfaces of linear response functions of butadiene molecule for the perturbations on the atomic sites. $\delta \rho(\mathbf{r}) / \delta \mathrm{v}_{\mathrm{x}}=0.035$ and $=-0.035$ are described as blue and red surfaces, respectively.


1 C


4 C


7 H


2 C


5 H


8 H


3 C


6 H


9 H


## References

1. Swope, W.C.; Andersen, H.C.; Berens, P.H.; Wilson, K.R. A computer simulation method for the calculation of equilibrium constants for the formation of physical clusters of molecules: Appllication to small water cluster. J. Chem. Phys. 1982, 76, 637.
2. Onufriev, A.; Bashford, D.; Case, D.A. Exploring protein native states and large-scale conformational changes with a modified generalized born model. Proteins 2004, 55, 383-394.
3. Berendsen, H.J.C.; Postma, J.P.M.; van Gunsteren, W.F.; DiNola, A.; Haak, J.R. Molecular dynamics with coupling to an external bath. J. Chem. Phys. 1984, 81, 3684-3690.
4. Ryckaert, J.-P.;Ciccotti, G.; Berendsen, H.J.C. Numerical integration of the cartesian equations of motion of a system with constraints: molecular dynamics of n-alkanes. J. Comp. Phys. 1977, 23, 327-341.
5. Grossfield, A. WHAM: The Weighted Histogram Analysis Method, Version 2.0.6. Available online: http://membrane.urmc.rochester.edu/content/wham (accessed on 28 August 2014).
6. Boczko, E.M.; Brooks, C.L., III. Constant-Temperature Free Energy Surfaces for Physical and Chemical Processes. J. Phys. Chem. 1993, 97, 4509-4513.
