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Ventilation Heat Recovery from Wood-Burning Domestic Flues. A Theoretical Analysis Based on a Triple Concentric Tube Heat Exchanger

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Abstract: This paper presents a new air-heating system concept for energy-efficient dwellings. It is a system designed to heat a low-energy building by coupling a heat-recovery ventilation system with a three-fluid heat exchanger located on the chimney of a wood-pellet stove. The proposed work focuses on the heat transfer that occurs between flue gases, the ventilation air and the combustion air within a triple concentric tube heat exchanger with no insulation at its outer surface. The main objective is to predict outlet temperature for the specific geometry of the heat exchanger studied here. Thus, the governing differential equations are derived for a counter-co-current flow arrangement of the three fluids. Then analytical solutions for the steady-state temperature distribution are obtained as well as the amount of heat transferred to the outside. An expression for the effectiveness of the heat exchanger is also proposed. Based on these results, calculations are performed on a case study to predict the fluid temperature distribution along the heat exchanger. Finally, a parametric study is carried out on this case study to assess the influence of the relevant parameters on the effectiveness of the heat exchanger. In addition, computation of heat losses to the outside justifies whether insulation is needed.

Keywords: low energy building; air heating system; wood pellet stove; heat recovery ventilation; triple concentric tube heat exchanger; analytical solution

Nomenclature:

- A heat transfer area (m^2) ;
- c_p specific heat capacity at constant pressure [J/(kg·K)];
- C heat capacity rate (J/K);
- C_r dimensionless heat capacity ratio (-);
- dA differential heat transfer area (m²);
- $d\phi$ differential heat flow rate (W);
- dx axial discretization step (m);
- *D* diameter (m);
- D_h hydraulic diameter (m);
- e thickness (m);
- *E* effectiveness (%);
- *h* convective heat transfer coefficient $[W/(m^2 \cdot K)];$
- h_r radiant heat transfer coefficient (W/(m² K));
- h_w conductive heat transfer coefficient [W/(m²·K)];

Greek Symbols

- ε emissivity (-);
- λ thermal conductivity [W/(m·K)];
- σ Stephan-Boltzmann constant [5.67 × 10⁻⁸ W/(m²·K⁴)];
- ϕ heat flow rate (W);
- η_f fin efficiency (-);

Subscripts

- 0 ambient air (surroundings);
- 1 first fluid (flue gases);
- 2 second fluid (ventilation air);
- 3 third fluid (combustion air);
- *a* internal tube;
- *b* intermediate tube;
- c external tube;
- o outside (surroundings);
- f fin(s);
- in inlet;

- L length (m);
- \dot{m} mass flow rate (kg/s);
- *N* dimensionless ratio of overall conductance *UA* and heat capacity rate *C* (-);
- *Nu* Nusselt number (-);
- NTU number of transfer units (-);
- *P* perimeter (m);
- *Pr* Prandtl number (-);
- *r* radius (m);
- *Ra* Rayleigh number (-);
- *Re* Reynolds number (-);
- *T* temperature (K);
- U overall heat transfer coefficient $[W/(m^2 \cdot K)];$
- *X* dimensionless length coordinate (-);
- *x* axial coordinate (m);
- z heat capacity rate ratio (\min/\max) (-).
- η_o overall surface efficiency (-);
- θ dimensionless temperature (-);
- ΔT temperature difference (K);
- $\Delta T_{\rm in}$ temperature difference depending on inlet temperatures and heat capacity rates (K).

out	outlet;
ext	external;
int	internal;
inner	inner;
outer	outer;
actual	actual;
т	mean;
max	maximum;
min	minimum;
tot	total

Acronyms

FG	flue gases;	TCTHE	triple concentric tube heat
VA	ventilation air;		exchanger;
CA	combustion air;	TCTHE-PI	triple concentric tube heat
OUT	outside;		exchanger with perfect insulation;
TOT	total;	TCTHE-NI	triple concentric tube heat
RSWPS	room-sealed wood pellet stove;		exchanger with no insulation;
MVHR	mechanical ventilation heat	DCTHE-NI	double concentric tube heat
	recovery;		exchanger with no insulation.

1. Introduction

Low energy buildings are characterized by a suitable orientation, a high thermal insulation and a very low air leakage [1]. One of the first results is that air-to-air heat recovery systems are increasingly integrated into the ventilation systems of buildings. In this context, Mardiana-Idayu and Riffat [2] have carried out a broad review on heat recovery technologies for building applications. First this study presents various types of heat recovery used in buildings namely fixed-plate, heat pipe, rotary wheel and run-around units. Then, the performance of heat recovery integrated in the ventilation systems of buildings including the coupling with dehumidification systems as well as the use of solar energy is presented and discussed. It should be noticed that the proposed system in this study is not mentioned in the review showing its innovative side. In addition, many studies focus on the efficiency of ventilation systems with heat recovery compared to standard ventilation systems [3,4]. Generally speaking and for the moderate climate region of middle Europe which is discussed in this study, it is shown that the use of heat recovery on air ventilation is not always competitive when operating costs are taken into account. Thus, the use of low power fans and low pressure loss systems is strongly recommended. Moreover, for very low energy buildings, traditional heating systems are no longer required nor even suited to these constructions, as highlighted by Feist [5]. In this context, heating requirements are low enough to be provided by the ventilation, with fresh air flow rates only and a maximum supply air temperature of about 50 °C. In such dwellings, a balanced mechanical ventilation system with a high efficiency heat recovery is especially recommended and the remaining heat demand could be easily covered by a renewable energy source, as a wood stove for example [5,6].

However, wood burning appliances present the disadvantage of concentrating heat only in the room where they are installed. Therefore, systems mechanically extracting hot air around the device are often used to redistribute heat to other rooms in the house. Nevertheless, the Consumer Safety Commission [7] pointed out that this type of system could create a negatively pressurised zone around the heating appliance which often leads to dust diffusion, but even can cause flue gases to be partially drawn in and redistributed in the other rooms. As a consequence, polluted air with abnormally high levels of carbon monoxide could place occupants at risk.

Instead of taking hot air above the heating appliance, as do conventional systems, the combined system presented in this study recovers heat by blowing fresh air into a specific heat exchanger, which is both integrated into the chimney of a wood pellet stove and connected to the ventilation air supply

network. Hence, while ensuring occupant safety, this combined system aims to meet the requirements of both heating and indoor air quality in energy efficient dwellings. A detailed description of this combined system comprising of a room-sealed wood pellet stove (RSWPS), a mechanical ventilation heat recovery (MVHR) and a triple concentric tube heat exchanger (TCTHE), is given in [8].

Considering that ideal configuration should combine all the advantages of the RSWPS and the MVHR while also introducing the TCTHE to better distribute heat in the house, the authors recommend coupling the three main components of the combined system as it is shown in Figure 1. As discussed in details in [8], this is a relevant configuration to standardize temperatures in the best way throughout the house and prevent any risk of ventilation air contamination by the combustion products. It should be noticed that the part of fresh air circulating in the TCTHE is called ventilation air.

The main objective of this work is to develop a theoretical model of a triple concentric tube heat exchanger in order to compute steady-state temperature distribution of the three fluids namely flue gases, ventilation air and combustion air. This will provide heat fluxes and effectiveness of the heat exchanger.



Figure 1. Scheme of the combined system configuration.

As described on the diagram in Figure 2, three different fluids are involved in heat transfers that occur in the triple concentric tube heat exchanger which has been developed in this study. Indeed, the flue gases (FG) are evacuated through the inner tube, the ventilation air (VA) recovers heat by circulating in the first annulus between the inner tube and the intermediate tube, and the combustion air (CA) which ensures the correct operation of the wood pellet stove circulates from top to bottom in the second annulus between the intermediate tube and the outer tube. In that way, flue gases and ventilation air are in a counter-flow arrangement, while ventilation air and combustion air are in a parallel-flow arrangement.





2. Theoretical Analysis

After the first analysis performed by Morley [9], many studies [10–12] have been conducted on exchanges involving three fluids in parallel stream flow arrangements and steady-state operation. For example, a significant design theory has been developed by Sorlie [13] for two parallel stream arrangements of a three-fluid heat exchanger with two thermal communications. Then, Zuritz [14] developed a set of analytical equations, he solved using Laplace transformation, for the computation of temperature distributions in a triple concentric tube heat exchanger. In other various cases studied [15–19], two or three thermal communications are considered, as well as four different flow arrangements and important dimensionless parameters. Thus, an extensive review of thermal design theory and steady state performance of three-fluid heat exchangers has been established by Sekulic and Shah [20] in order to unify the approach and systematize the solutions. To the knowledge of the authors, there has been little investigation of heat transfers in a triple concentric tube heat exchanger with no insulation at the outer surface (TCTHE-NI). In that way, this work is more related to studies by Prasad [21,22] on double concentric tube heat exchangers with no insulation at the outer surface (TCTHE-NI). In that exchangers with perfect insulation at the outer surface (TCTHE-NI).

As depicted in Figure 3, the physical model of the TCTHE-NI consists of three concentric tubes forming one circular flow passage and two concentric annular flow passages. Flue gases (FG), called fluid 1, enters the inner tube at x = L and flows in negative x-direction, while ventilation air (VA) and combustion air (CA), called fluids 2 and 3, respectively, both enter the heat exchanger at x = 0 and

flow in both annulus section in positive *x*-direction. Not to mention heat transfers with ambient air on the surroundings of the heat exchanger, because there is no insulation at its outer surface (OUT).

The theoretical analysis of the TCTHE-NI in a specific combination of counter-co-current flow arrangement, as defined in [20], has been carried out with the following assumptions:

- The system operates under steady-state conditions;
- Phase change does not take place;
- Bulk mean temperatures only depend on axial (x) direction;
- The thermo-physical properties of the fluid streams are constant and uniform;
- Thermal conduction is not assumed in fluids or walls parallel to the fluid flow direction;
- Heat transfer coefficients are independent of time and axial position;
- Temperature of the outside remains constant.





2.1. Mathematical Model

Under the above assumptions, the energy balance on a control volume of length dx yields the following set of equations:

$$d\phi_1 = d\phi_{12(a)} \tag{1}$$

$$d\phi_2 = d\phi_{12(a)} + d\phi_{32(b)} \tag{2}$$

$$d\phi_3 = -d\phi_{32(b)} - d\phi_{30(c)} \tag{3}$$

where the indices 0, 1, 2 and 3 denote outside, flue gases in the inner tube, ventilation air in the first annulus and combustion air in the second annulus, respectively.

In Equations (1) and (3), the differential heat flow rates $d\phi_1$ and $d\phi_3$ represent the heat lost by hot fluids 1 and 3 between locations x and x + dx. In Equation (2), the differential heat flow rate $d\phi_2$ denotes the heat retrieved by the cold fluid 2 between locations x and x + dx. Since there is no phase change in any fluids, the differential heat flow rates $d\phi$ can be expressed in terms of mass flow rates \dot{m} specific heats c_p , and temperature differentials dT:

$$\forall n \in \{1, 2, 3\} \qquad d\phi_n = \phi_n(x + dx) - \phi_n(x) = (\dot{m}c_p)_n dT_n = C_n dT_n$$
(4)

In the meantime and on the same differential control volume, the differential heat flow rates $d\phi_{ij(k)}$ transferred from a hot fluid *i* to a cold fluid *j* through a tube wall *k* can be expressed in terms of the difference between bulk temperatures $T_i - T_j$, the value of the overall heat transfer coefficients U_k and the corresponding differential heat transfer areas dA_k :

$$d\phi_{ii(k)} = U_k \, dA_k (T_i - T_j) \tag{5}$$

with: $\forall k \in \{a, b, c\}$ $dA_k = 2\pi r_k dx = \pi D_k dx = P_k dx$ (6)

According to an electrical analogy, the thermal circuit presented in Figure 4 is used to describe the combined heat transfers in the TCTHE-NI. The conduction, convection and radiation resistances per unit area, called l/h_w , l/h and l/h_r respectively, are computed from Equations (11)–(13).

Figure 4. Electrical analogy of the combined modes of heat transfer in the TCTHE-NI.



Concerning radiation, it should be noticed that flue gases, ventilation air and combustion air are non-participating fluids. Here, while maintaining the same radiant thermal resistance, the expression of the heat transfer by radiation in the annuli is assumed to be split in two parts involving the bulk temperature of the fluid between tube walls. In this way, the thermal resistance R_k of the tube k, or the overall heat transfer coefficient U_k which is arbitrary defined here in terms of the internal area of the tube wall $A_{k \text{ int}} = 2 \pi r_{k \text{ int}} L$, may be expressed for each tube wall as:

$$R_{a} = \frac{1}{U_{a} A_{a_{\text{int}}}} = \frac{1}{h_{1} A_{a_{\text{int}}}} + \frac{1}{h_{w_{a}} A_{a_{\text{int}}}} + \frac{1}{(h_{2a} + 2 h_{r_{ab}}) A_{a_{\text{int}}}} \left(\frac{r_{a_{\text{int}}}}{r_{a_{\text{ext}}}}\right)$$
(7)

$$R_{b} = \frac{1}{U_{b} A_{b_{\text{int}}}} = \frac{1}{(h_{2b} + 2 h_{r_{ab}}) A_{b_{\text{int}}}} + \frac{1}{h_{w_{b}} A_{b_{\text{int}}}} + \frac{1}{(h_{3b} + 2 h_{r_{bc}}) A_{b_{\text{int}}}} \left(\frac{r_{b_{\text{int}}}}{r_{b_{\text{ext}}}}\right)$$
(8)

$$R_{c} = \frac{1}{U_{c} A_{c_{\text{int}}}} = \frac{1}{(h_{3c} + 2 h_{r_{bc}}) A_{c_{\text{int}}}} + \frac{1}{h_{w_{c}} A_{c_{\text{int}}}} + \frac{1}{(h_{0} + h_{r_{co}}) A_{c_{\text{int}}}} \left(\frac{r_{c_{\text{int}}}}{r_{c_{\text{ext}}}}\right)$$
(9)

Taking the example of this study where fins are added on the external wall of the inner tube *a* to enhance the heat recovered by the ventilation air, the overall surface efficiency η_o should be introduced. Thus, the convective heat transfer coefficient of the external wall of the inner tube h_{2a} should be replaced in Equation (7) by h_{2af} .

$$h_{2af} = \eta_o \frac{A_{tot}}{A_{a_{int}}} h_{2a}$$
(10)

with:
$$\eta_o = 1 - \frac{nb_f A_f}{A_{tot}} (1 - \eta_f)$$
.

where nb_f is the number of fins, A_f is the surface area of one fin, A_{tot} is the total surface area *i.e.*, both the fins and the exposed portion of the base, and η_f is the efficiency of one fin. Expressions for the fin efficiency η_f are available in the literature [26] for several common fin arrays.

Then, the following expressions are used to assess the various coefficients that describe heat transfers by conduction, convection and radiation.

Conductive heat transfer coefficient

Through a cylindrical wall k, the thermal resistance for conduction R_{wk} , or the radial conductive heat transfer coefficient h_{wk} which is defined with respect to the internal wall surface area $A_{k int}$, may be expressed in terms of thermal conductivity λ_k , radius ratio $r_{k \text{ ext}}/r_{k \text{ int}}$ and heat exchanger length L:

$$\forall k \in \{a, b, c\} \qquad \qquad R_{w_k} = \frac{1}{h_{w_k} A_k} = \frac{\ln (r_{k_{ext}} / r_{k_{int}})}{2 \pi \lambda_k L}$$
(11)

Convective heat transfer coefficient

The convective heat transfer coefficient h_i of the fluid *i* is calculated using Nusselt number Nu, fluid thermal conductivity λ_i and hydraulic diameter D_h :

$$\forall i \in \{1, 2, 3, 0\} \qquad \qquad h_i = N u \lambda_i / D_h \tag{12}$$

where $D_h = D_{outer} - D_{inner}$ for an annular space between two (inner and outer) cylindrical walls. Then the average convection coefficients are determined thanks to correlations from the literature.

For forced convection laminar flow ($Re \le 2300$), Sieder and Tate [26] correlation is used both for the inner tube and for the annulus:

$$Nu = 1.86 (Re Pr D_h / L)^{1/3}$$

For forced convection transition and turbulent flows (Re > 2300), Colburn [26] correlation is used for the inner tube and Monrad and Pelton (cited by Jakob [27] and by Garcia-Valladares [28]) correlations are used for the annulus.

Inner tube:	$Nu = 0.023 \ Re^{4/5} \ Pr^{1/3}$	
Annulus:	$Nu = 0.020 \ Re^{4/5} \ Pr^{1/3} \ (D_{outer} \ / \ D_{inner})^{0.53}$	for the inner wall
	$Nu = 0.027 \ Re^{4/5} \ Pr^{1/3}$	for the outer wall

For free convection, the following correlations [26] are used to evaluate the heat transfer coefficient at the vertical external surface of the heat exchanger:

Laminar flow:	$10^4 \le Ra \le 10^9$	$Nu = 0.59 Ra^{1/4}$
Turbulent flow:	$10^9 \le Ra \le 10^{13}$	$Nu = 0.10 Ra^{1/3}$

According to the three flow rates and the three tubes diameters studied here and presented in Section 3.1, Reynolds number is always lower than 2300 for the combustion air, whereas it is always higher than 2300 (with a maximum value of 5500) for the flue gases and the ventilation air depending on both operating modes of the wood pellet stove and the ventilation system. Hence, Sieder and Tate

correlation, Colburn correlation and Monrad and Pelton correlation are used respectively for combustion air, flue gases and ventilation air.

Radiant heat transfer coefficient •

The radiant heat transfer coefficient h_{rkl} between two diffuse gray surfaces k and l that form an enclosure can be expressed as [26]:

$$\forall (kl) \in \{(ab); (bc); (co)\} \qquad h_{r_{kl}} = \sigma F_{kl} (T_k^4 - T_l^4) \approx \sigma F_{kl} 4 T_{m\,kl}^3$$
(13)

with: $F_{ab} = \left[\frac{1}{\varepsilon_a} + \frac{1 - \varepsilon_b}{\varepsilon_b} \left(\frac{r_{a}}{r_{b}}\right)\right]^{-1} \qquad F_{bc} = \left[\frac{1}{\varepsilon_b} + \frac{1 - \varepsilon_c}{\varepsilon_c} \left(\frac{r_{b}}{r_{c}}\right)\right]^{-1} \qquad F_{co} = \varepsilon_c$

 $T_{\text{m}ab} \approx \frac{T_{\text{m}1} + T_{\text{m}2}}{2} \ T_{\text{m}bc} \approx \frac{T_{\text{m}2} + T_{\text{m}3}}{2} \ T_{\text{m}co} \approx \frac{T_{\text{m}3} + T_0}{2}$ as first approximation, and:

and where $\forall i \in \{1, 2, 3\}$ the mean bulk temperature $T_{m_i} = \frac{T_{i \text{ in }} + T_{i \text{ out }}}{2}$ is iteratively calculated.

Then, considering the differential surface areas and the overall heat transfer coefficients given by Equations (6)–(9), the set of governing Equations (1)–(3) can thus be written in the following form:

$$C_1 dT_1 = U_a dA_a (T_1 - T_2)$$
(14)

$$C_2 dT_2 = U_a dA_a (T_1 - T_2) + U_b dA_b (T_3 - T_2)$$
(15)

$$C_3 dT_3 = -U_b dA_b (T_3 - T_2) - U_c dA_c (T_3 - T_0)$$
(16)

To provide a formulation with non-dimensional parameters and for brevity, following definitions are used:

$$\forall k \in \{a, b, c\}$$
 $A_k = A_{k \text{ int}}$ and $\Delta T_{\text{in}} = \frac{C_1(T_{1 \text{ in}} - T_{2 \text{ in}}) + C_3(T_{3 \text{ in}} - T_{2 \text{ in}})}{C_1 + C_3}$

as well as the below non dimensional parameters:

$$\begin{split} X &= \frac{x}{L} \qquad \qquad C_{r1} = \frac{C_1}{C_2} \qquad \qquad C_{r3} = \frac{C_3}{C_2} \\ N_1 &= \frac{U_a A_a}{C_1} \qquad \qquad N_3 = \frac{U_b A_b}{C_3} \qquad \qquad N_0 = \frac{U_c A_c}{C_3} \\ \theta_1 &= \frac{T_1 - T_0}{\Delta T_{in}} \qquad \qquad \theta_2 = \frac{T_2 - T_0}{\Delta T_{in}} \qquad \qquad \theta_3 = \frac{T_3 - T_0}{\Delta T_{in}} \end{split}$$

Hence, the set of governing equations can be written in the form:

$$\frac{d\theta_1}{dX} = N_1\theta_1 - N_1\theta_2 \tag{17}$$

$$\frac{d\theta_2}{dX} = N_1 C_{r1} \theta_1 - (N_1 C_{r1} + N_3 C_{r3}) \theta_2 + N_3 C_{r3} \theta_3$$
(18)

$$\frac{d\theta_3}{dX} = N_3\theta_2 - (N_0 + N_3)\theta_3 \tag{19}$$

 $\theta_1(1) = \theta_{1 \text{ in}}$ $\theta_2(0) = \theta_{2 \text{ in}}$ $\theta_3(0) = \theta_{3 \text{ in}}$ with boundary conditions: (20)

2.2. Analytical Solutions

By combining Equations (17)–(19), the following set of third order ordinary differential equations is obtained:

$$\forall n \in \{1, 2, 3\} \qquad \qquad \frac{d^3 \theta_n}{dX^3} + A \frac{d^2 \theta_n}{dX^2} + B \frac{d \theta_n}{dX} + C \theta_n = 0 \tag{21}$$

where the coefficients A, B and C are given by:

 $A = N_0 + N_1(C_{r1} - 1) + N_3(C_{r3} + 1), B = N_1(N_0 + N_3)(C_{r1} - 1) + N_3C_{r3}(N_0 - N_1), C = -N_0N_1N_3C_{r3}(N_0 - N_1) + N_3C_{r3}(N_0 - N_1) + N$

General solutions of the above system of 3 linear homogeneous third order ordinary differential Equation (21) are in the form of a sum of exponential functions in X:

$$\forall n \in \{1, 2, 3\} \qquad \qquad \theta_n(X) = \alpha_n e^{\lambda_1 X} + \beta_n e^{\lambda_2 X} + \gamma_n e^{\lambda_3 X}$$
(22)

where the coefficients α_n , β_n and γ_n should be determined for n = 1, 2 and 3 with boundary conditions and where the coefficients λ_1 , λ_2 and λ_3 are solutions of the characteristic Equation (23):

$$\lambda^{3} + A\lambda^{2} + B\lambda + C = 0$$
⁽²³⁾

Cardano's method is used to solve the cubic Equation (23). Thus, the quadratic term is eliminated and the so-called *depressed cubic* form is obtained:

$$x^3 + p x^2 + q = 0 (24)$$

with $x = \lambda + \frac{A}{3}$, $p = -\frac{A}{3} + B$, $q = \frac{A}{27}(2A^2 - 9B) + C$.

Considering the range of variation of the parameters used in this study, the discriminant Δ of the reduced form is always strictly negative:

$$\Delta = q^2 + \frac{4}{27} p^3 < 0$$

Hence, Equation (24) has three real roots:

$$\forall n \in \{1, 2, 3\}$$
 $x_n = 2\sqrt{\frac{-p}{3}} \cos\left[\frac{1}{3} \arccos\left(\frac{-q}{2}\sqrt{\frac{27}{-p^3}}\right) + \frac{2n\pi}{3}\right]$ (25)

and the three solutions of Equation (23) are:

$$\forall n \in \{1, 2, 3\} \qquad \qquad \lambda_n = x_n - \frac{A}{3} \tag{26}$$

Then, using Equations (17)–(19), the boundary conditions given by Equation (20) and the form of the solutions given by Equation (22) yields a set of three equations for each fluid:

Fluid 1
$$\theta_1(0) = \alpha_1 + \beta_1 + \gamma_1$$
 (27)

$$\left. \frac{d\theta_1}{dX} \right|_{X=0} = \alpha_1 \lambda_1 + \beta_1 \lambda_2 + \gamma_1 \lambda_3 = N_1 \left[\theta_1(0) - \theta_2(0) \right]$$
(28)

$$\theta_{1}(1) = \alpha_{1}e^{\lambda_{1}} + \beta_{1}e^{\lambda_{2}} + \gamma_{1}e^{\lambda_{3}}$$
⁽²⁹⁾

$$\theta_2(0) = \alpha_2 + \beta_2 + \gamma_2 \tag{30}$$

Fluid 2

$$\frac{d\theta_2}{dX}\Big|_{X=0} = \alpha_2 \lambda_1 + \beta_2 \lambda_2 + \gamma_2 \lambda_3 = N_1 C_{r_1} \theta_1(0) - \left[N_1 C_{r_1} + N_3 C_{r_3}\right] \theta_2(0) + N_3 C_{r_3} \theta_3(0)$$
(31)

$$\frac{d\theta_3}{dX}\Big|_{X=0} = \alpha_2 \lambda_1^{"} + \beta_2 \lambda_2^{"} + \gamma_2 \lambda_3^{"} + K_0 (\alpha_1 + \beta_1 + \gamma_1) = N_3 \theta_2(0) - [N_0 + N_3] \theta_3(0)$$
(32)

where $\forall n \in \{1, 2, 3\}$ $\lambda''_n = \frac{1}{N_3 C_{r3}} \left[\lambda_n^2 + (N_1 C_{r1} + N_3 C_{r3}) \lambda_n + N_1^2 C_{r1} \right]$, and $K_0 = \frac{-N_1^2 C_{r1}}{N_3 C_{r3}}$.

Fluid 3

$$\theta_3(0) = \alpha_3 + \beta_3 + \gamma_3 \tag{33}$$

$$\left. \frac{d\theta_3}{dX} \right|_{X=0} = \alpha_3 \lambda_1 + \beta_3 \lambda_2 + \gamma_3 \lambda_3 = N_3 \theta_2(0) - \left[N_0 + N_3 \right] \theta_3(0)$$
(34)

$$\frac{d^{2}\theta_{3}}{dX^{2}} + \left[N_{0} + N_{3}\right]\frac{d\theta_{3}}{dX}\Big|_{X=0} = \alpha_{3}\lambda_{1}\dot{\lambda_{1}} + \beta_{3}\lambda_{2}\dot{\lambda_{2}} + \gamma_{3}\lambda_{3}\dot{\lambda_{3}} = N_{3}\left[N_{1}C_{r1}\theta_{1}(0) - (N_{1}C_{r1} + N_{3}C_{r3})\theta_{2}(0) + N_{3}C_{r3}\theta_{3}(0)\right]$$
(35)

where $\forall n \in \{1, 2, 3\}$ $\lambda'_n = \lambda_n + N_0 + N_3$.

After some manipulations combining equations, coefficients α_n , β_n and γ_n and solutions $\theta_n(X)$, for n = 1, 2 and 3, can be expressed as functions of the unknown $\theta_1(0)$:

$$\alpha_n = F_n^{"} \theta_1(0) + K_n^{"} \tag{36}$$

$$\forall n \in \{1, 2, 3\} \qquad \beta_n = F'_n \theta_1(0) + K'_n \qquad (37)$$

$$\gamma_n = F_n \,\theta_1(0) + K_n \tag{38}$$

$$\forall n \in \{1,2,3\} \quad \theta_n(X) = \left[F_n^{"} e^{\lambda_1 X} + F_n^{'} e^{\lambda_2 X} + F_n e^{\lambda_3 X} \right] \theta_1(0) + \left[K_n^{"} e^{\lambda_1 X} + K_n^{'} e^{\lambda_2 X} + K_n e^{\lambda_3 X} \right]$$
(39)

where F_n , F_n' , F_n'' , K_n , K_n' and K_n'' , for n = 1, 2 and 3, are coefficients which can be expressed in terms of N_0 , N_1 , N_3 , C_{r1} , C_{r3} , λ_1 , λ_2 , λ_3 , $\theta_1(1)$, $\theta_2(0)$ and $\theta_3(0)$ (cf. Appendix).

Applying $\theta_1(X)$ and $\theta_2(X)$ from Equation (39) to the right-hand side of Equation (17) and through the integration of $d\theta_1(X)/dX$ from 0 to any *X* and then from 1 to any *X*, yields:

0 to X

$$\theta_{1}(X) = \left[1 + \xi_{1}(e^{\lambda_{1}X} - 1) + \xi_{2}(e^{\lambda_{2}X} - 1) + \xi_{3}(e^{\lambda_{3}X} - 1)\right]\theta_{1}(0) + \xi_{4}(e^{\lambda_{1}X} - 1) + \xi_{5}(e^{\lambda_{2}X} - 1) + \xi_{6}(e^{\lambda_{3}X} - 1)\right]$$
(40)

1 to X

$$\theta_{1}(X) = \left[\xi_{1}(e^{\lambda_{1}X} - e^{\lambda_{1}}) + \xi_{2}(e^{\lambda_{2}X} - e^{\lambda_{2}}) + \xi_{3}(e^{\lambda_{3}X} - e^{\lambda_{3}}) \right] \theta_{1}(0) + \xi_{4}(e^{\lambda_{1}X} - e^{\lambda_{1}}) + \xi_{5}(e^{\lambda_{2}X} - e^{\lambda_{2}}) + \xi_{6}(e^{\lambda_{3}X} - e^{\lambda_{3}}) + \theta_{1}(1)$$
(41)

where $\xi_1, \xi_2, \xi_3, \xi_4, \xi_5$ and ξ_6 are coefficients which can be expressed in terms of $\lambda_1, \lambda_2, \lambda_3, N_1, F_n, F_n', F_n'', K_n, K_n'$ and K_n'' , for n = 1 and 2 (cf. Appendix).

Equations (40) and (41) yield the following expression for the unknown $\theta_1(0)$:

$$\theta_{1}(0) = \frac{\xi_{4}(1 - e^{\lambda_{1}}) + \xi_{5}(1 - e^{\lambda_{2}}) + \xi_{6}(1 - e^{\lambda_{3}}) + \theta_{1}(1)}{1 + \xi_{1}(e^{\lambda_{1}} - 1) + \xi_{2}(e^{\lambda_{2}} - 1) + \xi_{3}(e^{\lambda_{3}} - 1)}$$
(42)

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Thus, for n = 1, 2 and 3, the coefficients $(\alpha_n, \beta_n, \gamma_n)$ and the solutions $\theta_n(X)$ can be completely determined for any *X* value between 0 and 1. Consequently, the bulk temperatures $T_n(X)$ can be expressed as:

$$\forall X \in \{0...1\} \text{ and } \forall n \in \{1,2,3\} \qquad T_n(X) = T_0 + \Delta T_{\text{in}} \theta_n(X)$$
(43)

Outlet temperatures of the three fluids are computed from:

$$T_{1 \text{ out}} = T_1(0) = T_0 + \Delta T_{\text{in}} \theta_1(0)$$
(44)

$$T_{2 \text{ out}} = T_2(1) = T_0 + \Delta T_{\text{in}} \theta_2(1)$$
(45)

$$T_{3 \text{ out}} = T_3(1) = T_0 + \Delta T_{\text{in}} \ \theta_3(1)$$
(46)

The heat rate $\phi_{30(c)}$ transferred across the non-adiabatic outside surface can be now determined through the integration of $\theta_3(X)$ from 0 to 1:

$$\phi_{30(c)} = U_c A_c \Delta T_{\rm in} \int_0^1 \theta_3(X) dX$$
(47)

$$\phi_{30(c)} = U_c A_c \Delta T_{in} \left[\frac{\alpha_3}{\lambda_1} (e^{\lambda_1} - 1) + \frac{\beta_3}{\lambda_2} (e^{\lambda_2} - 1) + \frac{\gamma_3}{\lambda_3} (e^{\lambda_3} - 1) \right]$$
(48)

For n = 1, 2 and 3, the heat flux ϕ_n assigned or retrieved by the fluid *n* is computed from:

$$\forall n \in \{1, 2, 3\} \qquad \qquad \phi_n = C_n (T_{\text{out}} - T_{\text{in}})_n \tag{49}$$

2.3. Effectiveness Relation

Considering an analogy with a two-fluid heat exchanger, the overall effectiveness *E* of a three-fluid heat exchanger is defined [26] as the ratio of the actual heat transfer rate ϕ_{actual} to the maximum possible heat transfer rate ϕ_{max} :

$$E = \frac{\phi_{\text{actual}}}{\phi_{\text{max}}} \tag{50}$$

As recommended by Sekulic and Kmecko [29], one should consider the particular purpose of the device to address the effectiveness issue for a three-fluid heat exchanger. As such, in this study, the actual heat transfer rate, ϕ_{actual} , is recovered by the fluid 2:

$$\phi_{\text{actual}} = \phi_2 = C_2 (T_{2 \text{ out}} - T_{2 \text{ in}})$$
(51)

Then, the maximum possible heat transfer rate, ϕ_{max} , that could be recovered by the fluid 2, from both fluids 1 and 3, can be expressed as:

$$\phi_{\max} = C_{\min(1;2)} \Delta T_{\max(1;2)} + C_{\min(3;2)} \Delta T_{\max(3;2)}$$
(52)

where $\forall (i;j) \in \{1;2;3\}^2 / i \neq j$, $C_{\min(i;j)} = \min(C_i; C_j)$, $\Delta T_{\max(i;j)} = \max(T_{i \text{ in }}; T_{j \text{ in }}) - \min(T_{i \text{ in }}; T_{j \text{ in }})$.

It is worth mentioning that owing to the fluids properties considered in this study [8], in any case $C_2 > C_1$, $C_2 > C_3$ and $C_1 > C_3$. It results that the maximum possible heat transfer rate ϕ_{max} is given by:

$$\phi_{\max} = C_1 \left(T_{1 \text{ in}} - T_{2 \text{ in}} \right) + C_3 \left(T_{3 \text{ in}} - T_{2 \text{ in}} \right)$$
(53)

Thus, heat exchanger effectiveness *E* is computed from:

$$E = \frac{C_2 (T_{2 \text{ out}} - T_{2 \text{ in}})}{C_1 (T_{1 \text{ in}} - T_{2 \text{ in}}) + C_3 (T_{3 \text{ in}} - T_{2 \text{ in}})}$$
(54)

Using previous notations, the heat exchanger effectiveness E can thus be written in the following form:

$$E = \frac{\theta_2(1) - \theta_2(0)}{C_{r_1} + C_{r_3}}$$
(55)

Since in any case $C_2 > C_1 + C_3$, the non-dimensional parameter *z* defined as the ratio of the minimum heat capacity rate to the maximum heat capacity rate can be introduced. Hence:

$$E = \frac{\theta_2(1) - \theta_2(0)}{z} \tag{56}$$

with:

$$z = C_{\min} / C_{\max} = (C_1 + C_3) / C_2 = C_{r1} + C_{r3}$$
(57)

Finally, the heat exchanger effectiveness *E* only depends on the non-dimensional parameter *z* and the temperature ratio $(T_{2out} - T_{2in})/\Delta T_{in}$:

$$E = \frac{T_{2 \text{ out}} - T_{2 \text{ in}}}{z \Delta T_{\text{in}}}$$
(58)

3. Parametric Study

3.1. Original Set of Input Data and Results

The input data used in the parametric study are given in Table 1. The diameters of the three tubes are 80, 180 and 230 mm respectively and the total length is 1.5 m. These values are close to measurements carried out on a bench test which is described in [8] and the comparison between experimental and numerical results showed that heat transfer rates are predicted with a relative difference lower than 5%.

Table 1. Input data and outlet temperatures for both TCTHE-NI and TCTHE-PI models (L = 1.5 m).

I _ 1 5	Т	- 25.90	Fluid 1	Fluid 2	Fluid 3
L = 1.5 m	1	$T_0 = 25 \text{ °C}$ (mm) (°C)	FG	VA	СА
	D _{a, b, c}	(mm)	80	180	230
Input data	$T_{in 1, 2, 3}$	(°C)	180	15	60
	$C_{1, 2, 3}$	(W/K)	10	20	5
Computed data	U _{a, b, c}	$(W/(m^2 K))$	5.0	3.2	2.0
TCTHE-NI	$T_{\rm out1,2,3}$	(°C)	153.09	31.6	38.2
TCTHE-PI	$T_{\rm out1,2,3}$	(°C)	153.13	32.1	45.2

The overall heat transfer coefficients U of the three tubes a, b, and c are calculated as explained previously in this paper. In case of perfect insulation, the parameter U_c is equal to zero. The last two rows of Table 1 give outlet temperatures obtained first by using the TCTHE model with no insulation (NI) which is presented in this paper and then by using the TCTHE with perfect insulation (PI) which has been developed by Ünal [23]. It can be noticed that inlet temperature of the combustion air (fluid 3) is higher than inlet temperature of the ventilation air (fluid 2) because the combustion air is preheated by the flue gases.

The heat capacity rate values, noted C, correspond to flow rates of about 35, 130 and 15 m³/h for the flue gases, the ventilation air and the combustion air respectively. These values were measured during laboratory tests and ventilation air flow rate meets the regulatory requirements of the ventilation of dwellings in France. Numerical calculations have been carried out for both TCTHE-NI and PI models to predict the temperature distributions of the three fluids along the heat exchanger (Figure 5). Considering the values adopted as input, the difference between temperature distribution of the TCTHE-NI and TCTHE-PI models does not appear very clearly except for the fluid 3 which is combustion air. However, even if the results obtained with the two models differ only at the second decimal for the flue gases, outlet temperature of each fluid is actually affected when taking into account the heat transfer with the outside, as it is displayed in Table 1.



Figure 5. Temperature distribution along the TCTHE-NI and PI (L = 1.5 m).

Then, the three graphs in Figure 6 show the heat balance for the TCTHE-NI ($U_c \neq 0$), TCTHE-NI ($U_c = 0$) and TCTHE-PI models. Each graph totalizes on its left part the heat flux assigned by the flue gases and by the combustion air, and on its right part the heat flux retrieved by the ventilation air and by the outside in case of no insulation. As expected, the algebraic sum of ϕ_{OUT} , ϕ_{FG} , ϕ_{VA} and ϕ_{CA} is always equal to zero. Thus, the heat balance is verified for all models studied. Moreover, if the overall heat transfer coefficient of the external wall tube is taken very close to zero, like here with $U_c = 1.10^{-10} \text{ W/(m}^2 \text{ K})$ for failing to divide by zero and lead to errors in the code, it must be emphasized that results of the TCTHE-NI ($U_c = 0$) and TCTHE-PI models are in perfect agreements.



Figure 6. Heat balance for the three TCTHE models (L = 1.5 m).

Although heat losses to the surroundings represent about 12% of the total heat transferred in the case of the TCTHE-NI, it is worth mentioning that the heat recovered by the ventilation air is only about 3% lower than in the case of the TCTHE-PI. Indeed and as expected, it is the amount of heat assigned by the combustion air to the surroundings which is mainly affected depending on whether the heat exchanger is insulated from the outside or not. Thus, as this study is conducted with an industrial approach, adding a fourth tube containing thermal insulation is not a relevant solution to increase the heat recovered by the ventilation air. In addition, efficiency of the pellet stove is not degraded because the combustion air temperature increases highly between the outlet of the heat exchanger and the inlet of the stove. Indeed, just behind the appliance, the combustion air recovers heat directly from flue gases in the concentric tube used for connecting the chimney with the stove.

As the amount of heat recovered by the ventilation air is quite similar for both models, as well as there are identical values for $\phi_{\text{max}} = 1875$ W and z = 0.75, there is little difference between the performances of the two models. Exact values for effectiveness are 17.7% for the TCTHE-NI and 18.3% for the TCTHE-PI. Nevertheless, even if effectiveness values are quite low, the most important issues under consideration in this study are the fact that enough heat can be recovered by the ventilation air and the fact that the evacuation of flue gases is correctly maintained.

3.2. Influence of Increasing Length

It should also be emphasized that the temperature distribution appears to be quasi-linear along the exchanger on the graph in Figure 5, whereas the temperature expression is written in the form a sum of exponentials in Equation (22). However, the exponential evolution of the temperature distribution in the exchanger is stressed when increasing the tube length, as is it shown for example on the graph in Figure 7 where both TCTHE-NI and PI results are presented for the same input data as used previously, except for the length L = 10 m.



Figure 7. Temperature distribution along the TCTHE-NI and PI (L = 10 m).

On the graph in Figure 7 at X = 1, it should be noted that the outlet temperature of the combustion air (hot fluid 3) is lower than the outlet temperature of the ventilation air (cold fluid 2). This temperature cross phenomenon appears for both TCTHE-NI and TCTHE-PI models and reflects the reversal of heat exchanges from a certain point in the heat exchanger. Differences between the outlet temperatures are also more significant whether the heat exchanger is insulated or not, as evidenced by the values given in Table 2.

Of course, such a 10m length heat exchanger is not really appropriate for the combined system studied here. However, it should be stressed that in this case again, heat losses to the surroundings represent about 12% of the total heat transferred by both the flue gases and the combustion air. Here, as the maximum heat transfer rate is still the same, $\phi_{max} = 1875$ W, the effectiveness of the triple concentric tube heat exchanger reaches 53.0% if there is no insulation from the surroundings and 56.6% when perfectly insulated from the outside.

I 10	$T_0 = 25 ^{\circ}\mathrm{C}$		Fluid 1	Fluid 2	Fluid 3
L = 10 m			FG	VA	СА
Computed data	U _{a, b, c}	$(W/(m^2 K))$	5.0	2.5	1.5
TCTHE-NI	$T_{\rm out \ 1, \ 2, \ 3}$	(°C)	76.5	64.7	42.0
TCTHE-PI	$T_{\rm out1,2,3}$	(°C)	77.5	68.0	53.0

Table 2. Outlet temperatures for TCTHE-NI and TCTHE-PI models (L = 10 m).

3.3. Influence of Adding Fins

In this study, fins could be added on the external wall of the inner tube to increase the amount of heat recovered by the ventilation air on the flue gas side. Considering the original set of input data for the TCTHE-NI model, the thickness, width and length of fins are 0.4 mm, 45 mm and 1160 mm, respectively. The fins are located in the longitudinal direction of the flow and their volume represent less than 2% of the total ventilation air flow volume. So, we consider initially the flow is not affected significantly by adding fins.

Several values of C_2 are considered to reflect the various flow rates of the ventilation air that may be encountered in the heat exchanger, *i.e.*, parameter z = 0.75 to 0.375 in experiments. As expected, the graph in Figure 8 shows that adding fins increases the exchanger effectiveness, but this trend is less pronounced when more than ten fins are added. As a result of this analysis, eight fins are planned to be added to the external wall of the inner tube to enhance the exchanger effectiveness without involving excessive manufacturing costs.

Figure 8. Effectiveness as a function of the number of fins for the TCTHE-NI model with the original set of input data and with ventilation air flow variations.



3.4. Effectiveness Charts

Considering the original set of input data presented in Table 1, a short parametric study is carried out to assess how the effectiveness E of the TCTHE-NI is influenced by the four following input parameters: C_2 , D_b , $T_{2 \text{ in}}$ and T_0 . Thus, the four graphs in Figure 9 show the variation of the effectiveness with one of the four previous parameters and with the number of transfer unit (NTU), which is a non-dimensional parameter defined here with previous notations:

NTU =
$$\frac{U_a A_a + U_b A_b}{C_1 + C_3}$$
 (59)

For each graph, several values of NTU are obtained by varying the length of the heat exchanger between 1 and 30 m. However, it is worth noting that the NTU value is always lower than 0.5 when calculations are carried out with the original set of input data defined in this study.

First, Figure 9a shows that the effectiveness increases with the growth of the heat capacity rate C_2 . Indeed, when the mass flow rate of the fluid2increases, it causes both a decrease of the non-dimensional parameter z defined in Equation (57), and an increase of convective heat transfer coefficients $h_{2 \text{ int}}$ and $h_{2 \text{ ext}}$ in the first annulus. As in the case of two-fluid heat exchangers, the effectiveness value is minimum for z = 1 and maximum for $z \rightarrow 0$. Finally, the variation of C_2 has a fairly significant influence on the effectiveness, even for low values of NTU.



Figure 9. Effectiveness charts for the TCTHE-NI model.

Then, as shown in Figure 9b, the effectiveness increases when the tube diameter D_b decreases. In simple words, reducing the flow area for the fluid 2 results in a significant increase of the convective heat transfer coefficients $h_{2 \text{ int}}$ and $h_{2 \text{ ext}}$ in the first annulus and this effect overrides the negative effect of the decrease in heat transfer surface area with fluid 3. Here again, the choice of the size of D_b has a significant influence on the effectiveness, even for low values of NTU and this observation is consistent with the findings of Ünal [24], which highlight that the relative sizes of the three tube diameters are important parameters that influence the performance of the exchanger.

Figure 9c shows that the effectiveness increases when the inlet temperature $T_{2 \text{ in}}$ decreases, *i.e.*, when the non-dimensional parameter ΔT_{in} increases. Considering the specific combination of the counter-current co-current flow arrangement studied here, the actual heat recovered by the fluid 2 increases at a greater rate than the maximum heat recoverable. Nevertheless, it is of some interest to note that the variation of $T_{2 \text{ in}}$ has a relatively limited influence on the effectiveness, especially for low values of NTU.

When the external temperature T_0 increases, Figure 9d shows that the effectiveness of the exchange also increases. In addition and as expected, the effectiveness reaches its maximum when the TCTHE is perfectly isolated (PI). However, as mentioned in the previous sections, the variation of the outside temperature mainly affects the outlet temperature of the combustion air. Thus, T_0 has little influence on the effectiveness for low NTU values, but its impact becomes more significant for higher NTU values.

The analysis of the four graphs in Figure 9 verifies the consistency of the mathematical model presented in the first part of this study and the equations derived can be used to assess the performance of a TCTHE-NI. Then the effectiveness is revealed to be more significantly influenced by parameters C_2 or D_b than by parameters $T_{2 \text{ in}}$ or T_0 .

3.5. Importance of Heat Losses

The heat transferred to the outside, ϕ_{OUT} , may be compared to the total heat assigned or retrieved, $\phi_{TOT} = \phi_{VA} + \phi_{OUT} = \phi_{FG} + \phi_{CA}$, in the TCTHE-NI. Figure10 shows the variations of the ratio ϕ_{OUT}/ϕ_{TOT} with increasing NTU values. Of course, this ratio is always equal to zero in case of perfect insulation (PI). But depending on the surroundings temperature, T_0 , heat losses to the outside may represent a significant part of heat transfer that occurs in the exchanger.

Considering the original set of input data, Figure 10a shows that heat losses remain below 20% for NTU < 1.5 and $T_0 > 10$ °C, reflecting the majority of cases that may be encountered in our application. Indeed, in the case of the combined system studied here, the TCTHE-NI is commonly located just behind the heating appliance in a closet in which the temperature remains between 15 °C and 25 °C.

Figure 10. Percentage of heat transferred to the surroundings for the TCTHE-NI model with the original set of input data (a) and with an increase of the ventilation air flow (b).



(a) Percentage of heat losses with $C_2 = 20$ W/K (b) Percentage of heat losses with $C_2 = 40$ W/K

In addition, as shown in Figure 10b, heat transfers focus on the side of the ventilation air as soon as its flow rate increases. In the TCTHE-NI, the proportion of heat losses to the surroundings compared to the total heat transferred remains therefore lower than 10% if the ambient air temperature is at least 20 °C. Hence, as heat losses to the outside remain moderate in our case study, adding a fourth tube containing insulation material cannot be justified in view of larger investment and manufacturing costs. In addition, heat losses from the heat exchanger are in the heated area. Therefore, it is not really heat losses.

4. Conclusions

Analytical solutions for the steady-state temperature of three heat exchanging fluids along the length of a triple concentric tube heat exchanger with no insulation at the outer surface are obtained for

a specific combination of counter-current and co-current arrangement. The amount of heat transferred to the outside through the external tube is also calculated. For a case study, the temperature distribution of the three fluids is graphically represented and its exponential variation is stressed by increasing the length of the heat exchanger. Heat balance calculations and comparison with the perfectly insulated model have provided validation of the mathematical model developed in this paper. With the input data values of the parametric study conducted here, the temperature of the ventilation air increases by about 15 °C in the heat exchanger while the temperature of flue gases and combustion air decreases about 30 °C and 20 °C respectively. Then the importance of adding fins to the external surface of the influence of several parameters, especially C_2 and D_b , on the heat exchanger effectiveness is also assessed through a parametric study conducted in the continuity of the case study. Heat losses to the outside represent around 10% of the total heat transferred in the TCTHE-NI for the set of input data used in the case study, so insulation of the heat exchanger is not really required if it is placed in the heat eventual used in the dwelling.

Finally, this theoretical study could be extended to the various flow arrangements encountered in a triple concentric tube heat exchanger and the equations derived here could also be used to perform sizing calculations. In addition, we are aware that this model developed in this study is quite simple but comparison between experimental and numerical results showed that heat transfer rates are predicted with a relative difference lower than 5% [8]. Furthermore, CFD approach will be used in a future work in order to study in more details the fluid flows within the heat exchanger.

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Appendix: Explanation of the Mathematical Model Coefficients

Coefficients α_n , β_n and γ_n are given by Equations (36)–(38):

 $\forall n \in \{1, 2, 3\} \qquad \alpha_n = F_n^{"} \theta_1(0) + K_n^{"} \qquad \beta_n = F_n^{'} \theta_1(0) + K_n^{'} \qquad \gamma_n = F_n^{'} \theta_1(0) + K_n^{'}$

where F_n , F_n' , F_n'' , K_n , K_n' and K_n'' , for n = 1, 2 and 3, are coefficients which can be expressed in terms of N_0 , N_1 , N_3 , C_{r1} , C_{r3} , λ_1 , λ_2 , λ_3 , $\theta_1(1)$, $\theta_2(0)$ and $\theta_3(0)$. $\boxed{n=1}$

 $\begin{aligned} \alpha_{1} &= F_{1}^{"} \theta_{1}(0) + K_{1}^{"} & K_{1}^{"} &= G_{1}^{"} \theta_{2}(0) + H_{1}^{"} \theta_{1}(1) & F_{1}^{"} &= 1 - (F_{1}^{'} + F_{1}) \\ \beta_{1} &= F_{1}^{'} \theta_{1}(0) + K_{1}^{'} & K_{1}^{'} &= G_{1}^{'} \theta_{2}(0) + H_{1}^{'} \theta_{1}(1) & G_{1}^{"} &= -(G_{1}^{'} + G_{1}) \\ \gamma_{1} &= F_{1}^{'} \theta_{1}(0) + K_{1} & K_{1} &= G_{1}^{'} \theta_{2}(0) + H_{1}^{'} \theta_{1}(1) & H_{1}^{"} &= -(H_{1}^{'} + H_{1}) \end{aligned}$

$$F_{1}' = \frac{1}{\lambda_{2} - \lambda_{1}} \Big[(N_{1} - \lambda_{1}) - F_{1} (\lambda_{3} - \lambda_{1}) \Big] \qquad F_{1} = -F \Big[(\lambda_{2} - \lambda_{1}) e^{\lambda_{1}} + (e^{\lambda_{2}} - e^{\lambda_{1}}) (N_{1} - \lambda_{1}) \Big] \\ G_{1}' = \frac{-1}{\lambda_{2} - \lambda_{1}} \Big[N_{1} + G_{1} (\lambda_{3} - \lambda_{1}) \Big] \qquad G_{1} = F (e^{\lambda_{2}} - e^{\lambda_{1}}) N_{1} \qquad H_{1} = F (\lambda_{2} - \lambda_{1}) \\ H_{1}' = \frac{-1}{\lambda_{2} - \lambda_{1}} \Big[H_{1} (\lambda_{3} - \lambda_{1}) \Big] \qquad F = \Big[(\lambda_{2} - \lambda_{1}) (e^{\lambda_{3}} - e^{\lambda_{1}}) - (\lambda_{3} - \lambda_{1}) (e^{\lambda_{2}} - e^{\lambda_{1}}) \Big]^{-1}$$

$$n = 2$$

$$\begin{aligned} \alpha_{2} &= F_{2}^{"} \theta_{1}(0) + K_{2}^{"} \\ \beta_{2} &= F_{2}^{'} \theta_{1}(0) + K_{2}^{'} \\ \beta_{2} &= F_{2}^{'} \theta_{1}(0) + K_{2}^{'} \\ \gamma_{2} &= F_{2}^{'} \theta_{1}(0) + K_{2}^{'} \\ \gamma_{2} &= F_{2}^{'} \theta_{1}(0) + K_{2}^{'} \\ K_{2}^{'} &= -G \left[\frac{(N_{1} C_{r1} + N_{3} C_{r3} + \lambda_{1})}{\lambda_{2} - \lambda_{1}} + \frac{(N_{3} - \lambda_{1}^{'})}{\lambda_{2}^{'} - \lambda_{1}^{''}} \right] \theta_{2}(0) + G \left[\frac{N_{3} C_{r3}}{\lambda_{2} - \lambda_{1}} + \frac{N_{0} + N_{3}}{\lambda_{2}^{'} - \lambda_{1}^{''}} \right] \theta_{3}(0) \\ F_{2}^{"} &= -(F_{2}^{'} + F_{2}^{'}) \\ F_{2}^{'} &= -(F_{2}^{'} + F_{2}^{'}) \\ F_{2}^{'} &= G \left[\frac{N_{1} C_{r1}}{\lambda_{2} - \lambda_{1}} \left[N_{1} C_{r1} - F_{2} (\lambda_{3} - \lambda_{1}) \right] \\ F_{2} &= G \left[\frac{N_{1} C_{r1}}{\lambda_{2} - \lambda_{1}} + \frac{K_{0}}{\lambda_{2}^{''} - \lambda_{1}^{''}} \right] \\ \text{with } F_{1}^{"} + F_{1}^{'} + F_{1}^{'} = 1 \\ \end{aligned}$$

$$n = 3$$

$$\alpha_{3} = F_{3}^{"} \theta_{1}(0) + K_{3}^{"} \qquad K_{3}^{"} = G_{3}^{"} \theta_{2}(0) + H_{3}^{"} \theta_{1}(1) \qquad F_{3}^{"} = 1 - (F_{3}^{'} + F_{3}^{'})$$

$$\beta_{3} = F_{3}^{'} \theta_{1}(0) + K_{3}^{'} \qquad K_{3}^{'} = G_{3}^{'} \theta_{2}(0) + H_{3}^{'} \theta_{1}(1) \qquad G_{3}^{"} = -(G_{3}^{'} + G_{3}^{'})$$

$$\gamma_{3} = F_{3} \theta_{1}(0) + K_{3} \qquad K_{3} = G_{3} \theta_{2}(0) + H_{3} \theta_{1}(1) \qquad H_{3}^{"} = -(H_{3}^{'} + H_{3}^{'})$$

$$\begin{split} F_{3}^{'} &= \frac{-1}{\lambda_{2} - \lambda_{1}} \Big[(N_{0} + N_{3} + \lambda_{1}) + F_{3} (\lambda_{3} - \lambda_{1}) \Big] & F_{3} = H \Big[(\lambda_{2} - \lambda_{1}) (N_{3}^{2} C_{r3} - \lambda_{1} \dot{\lambda}_{1}) + (\lambda_{2} \dot{\lambda}_{2} - \lambda_{1} \dot{\lambda}_{1}) (N_{0} + N_{3} + \lambda_{1}) \Big] \\ G_{3}^{'} &= \frac{1}{\lambda_{2} - \lambda_{1}} \Big[N_{3} - G_{3} (\lambda_{3} - \lambda_{1}) \Big] & G_{3} = -H N_{3} \Big[(\lambda_{2} \dot{\lambda}_{2} - \lambda_{1} \dot{\lambda}_{1}) + (N_{1} C_{r1} + N_{3} C_{r3}) (\lambda_{2} - \lambda_{1}) \Big] \\ H_{3}^{'} &= \frac{-1}{\lambda_{2} - \lambda_{1}} \Big[H_{3} (\lambda_{3} - \lambda_{1}) \Big] & H = \Big[(\lambda_{3} \dot{\lambda}_{3} - \lambda_{1} \dot{\lambda}_{1}) (\lambda_{2} - \lambda_{1}) - (\lambda_{3} - \lambda_{1}) (\lambda_{2} \dot{\lambda}_{2} - \lambda_{1} \dot{\lambda}_{1}) \Big]^{-1} \end{split}$$

As presented in Equations (22) and (39), general solutions of the system of 3 linear homogeneous third order ordinary differential equations are in the form of a sum of exponential functions in X.

$$\forall n \in \{1, 2, 3\} \qquad \qquad \theta_n(X) = \alpha_n e^{\lambda_1 X} + \beta_n e^{\lambda_2 X} + \gamma_n e^{\lambda_3 X}$$

$$\forall n \in \{1, 2, 3\} \qquad \qquad \theta_n(X) = \left[F_n^{"} e^{\lambda_1 X} + F_n^{'} e^{\lambda_2 X} + F_n^{'} e^{\lambda_3 X}\right] \theta_1(0) + \left[K_n^{"} e^{\lambda_1 X} + K_n^{'} e^{\lambda_2 X} + K_n^{'} e^{\lambda_3 X}\right]$$

$$= \left[F_n^{"} e^{\lambda_1 X} + F_n^{'} e^{\lambda_2 X} + F_n^{'} e^{\lambda_3 X}\right] \theta_1(0) + \left[K_n^{"} e^{\lambda_1 X} + K_n^{'} e^{\lambda_2 X} + K_n^{'} e^{\lambda_3 X}\right]$$

with

$$\theta_{1}(0) = \frac{\xi_{4}(1 - e^{\lambda_{1}}) + \xi_{5}(1 - e^{\lambda_{2}}) + \xi_{6}(1 - e^{\lambda_{3}}) + \theta_{1}(1)}{1 + \xi_{1}(e^{\lambda_{1}} - 1) + \xi_{2}(e^{\lambda_{2}} - 1) + \xi_{3}(e^{\lambda_{3}} - 1)}$$

where $\xi_1, \xi_2, \xi_3, \xi_4, \xi_5$ and ξ_6 are coefficients which can be expressed in terms of $\lambda_1, \lambda_2, \lambda_3, N_1, F_n, F_n', F_n'', K_n, K_n'$ and K_n'' , for n = 1 and 2.

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