

Article

# Influence of Carrier Gases on the Quality of Epitaxial Corundum-Structured $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> Films Grown by Mist Chemical Vapor Deposition Method

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**Abstract:** This report systematically investigates the influence of different carrier gases ( $O_2$ ,  $N_2$ , and air) on the growth of gallium oxide ( $Ga_2O_3$ ) thin films on c-plane sapphire substrates by using the mist-CVD method. Although XRD and Raman measurements show that the pure corundum-structured  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> with single (0006) plane orientation was successfully obtained for all three different carrier gases, the crystal quality could be greatly affected by the carrier gas. When  $O_2$  is used as the carrier gas, the smallest full-width at half maximum (FWHM), the very sharp absorption cutoff edge, the perfect lattice structure, the highest growth rate, and the smooth surface can be obtained for the epitaxial  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> film as demonstrated by XRD, UV-VIS, TEM, AFM (Atomic Force Microscope), and SEM measurements. It is proposed that the oxygen content in carrier gas should be responsible for all of these results. XPS (X-ray photoelectron spectroscopy) analysis also confirms that more oxygen elements can be included in epitaxial film when  $O_2$  is used as the carrier gas and thus help improve the crystal quality. The proper carrier gas is essential for the high quality  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> growth.

**Keywords:** wide-bandgap semiconductor; α-Ga<sub>2</sub>O<sub>3</sub>; mist chemical vapor deposition (mist-CVD); carrier gas; transparent semiconductor

## 1. Introduction

As an ultra-wide-bandgap semiconductor with the obvious advantages of stable physical chemistry, low dielectric constant, and high mechanical strength, gallium oxide (Ga<sub>2</sub>O<sub>3</sub>) is attracting increasing attention as a new promising competitor to III-nitrides and SiC for various applications in high-voltage and high-power electronics and ultraviolet optoelectronics [1]. Compared with other oxides such as ZnO (3.24 eV) and In<sub>2</sub>O<sub>3</sub> (3.6 eV), Ga<sub>2</sub>O<sub>3</sub> has a larger bandgap energy of approximately 5 eV, which means a shorter absorption cutoff wavelength and a much higher power application. There are a total of five different polytypes ( $\alpha$ ,  $\beta$ ,  $\varepsilon$ ,  $\delta$ , and  $\gamma$ ) for Ga<sub>2</sub>O<sub>3</sub>. Until now, the most studied polytype was  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> because it is easy to obtain bulk and film  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> materials by the conventional crystal growth or epitaxial growth techniques, such as edge-defined film-fed growth, float-zone method, Czochralski method, molecular beam epitaxy, and metal organic chemical vapor deposition.  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> has a bandgap of 4.8 eV and high Baliga's figures of merit (FOM) of 3000, which is obviously superior to GaN and SiC. However,  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> is not the best candidate in various phases for the power application considering the bandgap. Compared to  $\beta$  phase, corundum-structured  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>, another important phase for Ga<sub>2</sub>O<sub>3</sub>, has a wider bandgap of around 5.3 eV which can result in a larger Baliga's FOM,



in theory. Thus,  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> has great potential for application in power devices. A wider bandgap of about 5.3 eV means that the absorption cutoff wavelength can be shorter than 240 nm and then  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> is more suitable for ultraviolet optoelectronics. Although  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> has great application potential, the research of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> still lags far behind  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> in large part because it is more difficult to obtain the high-quality  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> material than  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>.

 $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> is a metastable phase and the bulk material still cannot be obtained. There is no commercial  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> bulk substrate to date, so its homo-epitaxial growth is still difficult. Fortunately, the heterogeneous epitaxy provides an efficient way to obtain the  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> material. Corundum-structured gallium oxide belongs to the space group of R-3c with the lattice parameters a = b = 4.98 Å, c = 13.43 Å,  $\alpha = \beta = 90^{\circ}$ , and  $\gamma = 120^{\circ}$  [2], and the lattice mismatches between  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> and  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> (sapphire) are only 4.81% and 3.54% in the a- and c-axis directions. The same crystal structure and small lattice mismatch make it easy to grow the  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> material on the  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> sapphire substrate. More importantly, the same crystal structure means that it is attractive from the viewpoint of fabricating alloys for  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> with other corundum-structured materials, such as  $\alpha$ -Al<sub>2</sub>O<sub>3</sub>, Fe<sub>2</sub>O<sub>3</sub>, and Cr<sub>2</sub>O<sub>3</sub> for bandgap and material engineering [3], which is another advantage of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>.

Recent studies have shown that the growth of crystalline  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> on an inexpensive sapphire substrate is an efficient way to obtain the  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> material. The key technology for the growth of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> is ultrasonic mist chemical vapor deposition (mist-CVD) method [4–6]. In the growth of a metal oxide, water solutions of safe and inexpensive chemicals containing the metal, for example, acetate or acetylacetonate, have been used as the source. By atomizing the source solution ultrasonically, it turns into mist particles, which are then transferred by a carrier gas to a reaction chamber. In this way, metal elements are supplied without the use of organometallic sources. They react with an oxygen source, which may be water or oxygen gas. This offers sufficient overpressure of oxygen with respect to the metal source and prevents the formation of oxygen vacancies. Therefore, mist-CVD method is suitable for epitaxial  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> on sapphire substrate and can reduce the material cost [7,8]. Based on the grown  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>, various applications of Metal Epitaxial-Semiconductor Field Effect Transistor [9], Schottky barrier diodes [10], and solar-blind photodetectors [11] have been demonstrated. For example, a high performance Schottky diode with a breakdown voltage over 1 kV and a small specific on-resistance of  $2.5 \text{ m}\Omega \cdot \text{cm}^2$  has been achieved and a normally-off MOSFET has been shown based on  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> material grown by a mist-CVD system [9,10]. However, the present device performance is greatly inferior to the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> counterparts and the main reason is still the poor  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> film quality. Thus, more attention is urgently required to improve the quality of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> now and in the future. Recently, highly crystalline  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> thin films have been successfully grown at atmospheric pressure by mist-CVD on c-sapphire substrates, whose temperatures of 400–500 °C are reasonably low and the optimal growth conditions of solution concentration, growth temperature, carrier gas velocity, and film thickness have also been investigated [12,13]. However, there is still no systematical study about how the different carrier gas affects the film quality.

In this paper, we systematically investigate the influence of different carrier gases ( $O_2$ ,  $N_2$ , and air) on the film quality for the growth of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> on c-plane sapphire substrates by using the mist-CVD method. It is demonstrated that the crystallization quality will be different when the gallium source is carried by different gases. When  $N_2$  and  $O_2$  are used as the carrier gases,  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> achieves a relative smooth surface. When  $O_2$  is the carrier gas,  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> achieves the smallest half-height width. The oxygen element in the carrier gas may be an important reason to prevent the generation of oxygen vacancies, thus influencing the quality of the thin films. The results provide constructive perspectives for the material quality improvement.

#### 2. Materials and Methods

In the present experiment, we used gallium acetylacetonate as the gallium source, which was dissolved in deionized water. A small quantity of hydrochloric acid was added to dissolve gallium acetylacetonate completely. The concentration of the solution was adjusted to 0.05 M. By atomizing

the source solution ultrasonically, it turned into mist particles (diameter of ~3  $\mu$ m at an ultrasonic frequency of 2.4 MHz), which was then carried by air, N<sub>2</sub>, and O<sub>2</sub>, respectively, to the heated reaction chamber. In the chamber, the (0001) sapphire substrate was placed on a sample holder that was kept at 400 °C, because under this condition a high quality  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> can be obtained, as shown in Figure S1 in the supplemental information. The growth time was kept at 1h and the rate of the carrier gas was set to be 6 L/min.

The structural properties of Ga<sub>2</sub>O<sub>3</sub> films were investigated using a range of complementary techniques. X-ray diffraction (XRD) patterns were obtained from an X-ray diffractometer (D8 Advance, Bruker, Karlsruhe, Germany). The transmittance was measured by a dual-beam 950 UV-VIS spectrometer. XPS measurements were performed by the Escalab 250Xi (Waltham, MA, USA) with a source of monochromatic Al-Ka (1486.6 eV). The film morphologies were characterized by a field emission scanning electron microscope (SEM JSM-7800F, Tokyo, Japan), atomic force microscopy (AFM) (Agilent 5500, Palo Alto, Santa Clara, CA, USA), and high-resolution transmission electron microscopy (TEM) (Tecnai G2 F20 S-Twin, Hillsboro, OR, USA). The Raman spectra were measured using a confocal Jobin Yvon LavRam HR800 micro-Raman spectrometer (Edison, NJ, USA) with a charge-coupled device (CCD) detector.

#### 3. Results

Figure 1 shows the XRD spectra of the samples with air, N<sub>2</sub>, and O<sub>2</sub> as the carrier gases. The spectra are dominated by the diffraction peaks at 40.26° and 41.66°, which correspond to the (0006) planes of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> epilayer and sapphire substrate, respectively. No other peaks are found, which shows that all of the Ga<sub>2</sub>O<sub>3</sub> films show the obvious pure alpha phase. These XRD spectra show that the  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> films had a preferential c-axis orientation along the c-axis of the sapphire substrate and the calculated lattice constant along the c-axis is 1.34 nm. Paying attention to the sample grown with O<sub>2</sub> as the carrier gas, the full-width at half maximum (FWHM) of the  $\omega$  scan rocking curve is as small as 72 arcsec, indicating a high quality  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>. However, for the samples grown with air and N<sub>2</sub> as the carrier gases, FWHM of the  $\omega$  scan rocking curves are 88.6 arcsec and 86.4 arcsec, respectively, indicating a relatively inferior crystal quality. Considering the different oxygen content in the carrier gases, it is supposed that Ga<sub>2</sub>O<sub>3</sub> film grown with O<sub>2</sub> as the carrier gas has less defects, such as oxygen vacancies (V<sub>O</sub>), which may be the key to improving the quality of crystallization.



**Figure 1.** (a) X-ray diffraction  $2\theta/\theta$  scan spectra for  $Ga_2O_3$  films grown with different carrier gases. (b) The corresponding (0006) XRD diffraction rocking curves. The full-width at half maximum of  $Ga_2O_3$  films grown with air,  $N_2$ , and  $O_2$  as the carrier gases are 88.6 arcsec, 86.4 arcsec, and 72 arcsec, respectively.

Raman spectra of the epilayers grown with air,  $N_2$ , and  $O_2$  as the carrier gases were measured between 100 cm<sup>-1</sup> and 800 cm<sup>-1</sup> at room temperature to confirm the crystalline quality of the deposited films, and the results are presented in Figure 2. A 514 nm laser was used as the excitation source and the laser beam was focused by a microscope lens system ( $\times$ 50 ulwd) yielding a spot size of 1  $\mu$ m during the Raman measurement. The Raman peaks at 418 cm<sup>-1</sup> and 749 cm<sup>-1</sup> belong to the sapphire substrate [14]. The Raman peaks located at 431.3 cm<sup>-1</sup>, 577 cm<sup>-1</sup>, and 692 cm<sup>-1</sup> are the Raman-allowed vibrational modes of  $E_g$ ,  $A_{1g}$ , and  $E_g$  for  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>, respectively, and are consistent with the theoretical calculations [15]. The high-frequency  $A_{1g}$  mode at 577 cm<sup>-1</sup> mainly involves the vibration of oxygen atoms perpendicular to the c-axis. The linewidth of the peak  $A_{1g}$  is as narrow as 3.9 cm<sup>-1</sup> for the epilayer grown with O<sub>2</sub> as the carrier gas, strongly suggesting the high crystallinity of the epilayer. By comparing the epilayers grown with different carrier gases, the displacement and intensity of the Raman peaks do not change obviously, demonstrating that the stress in the films is mainly determined by other factors, such as lattice mismatch or growth temperature, instead of the carrier gases. It could also be observed that except the Raman peaks of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>, no other peaks are observed in the Raman spectra, which indicates that all of the epilayers are pure  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> without other phases and confirm the conclusion from the XRD measurements.



Figure 2. Raman spectra of the  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> epilayers grown with air, N<sub>2</sub>, and O<sub>2</sub> as the carrier gases.

The variation of optical transmittance spectra (200–800 nm) were performed on  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> films as shown in Figure 3. All the samples exhibited a transmittance higher than 80% in the visible to near-UV regions. The relationship between the absorption coefficient  $\alpha$  and the optical bandgap (E<sub>g</sub>) is  $\alpha$ hv = A(hv-Eg)<sup>1/2</sup>, where A is the material-dependent constant, h is the Planck's constant, and v is the frequency of the incident light [16]. The optical band gap can be evaluated from the ( $\alpha$ hv)<sup>2</sup> versus photon energy (hv) graph by linear extrapolations to zero absorption coefficient [17]. The inset shows the plot of ( $\alpha$ hv)<sup>2</sup> as function of photon-energy hv. The bandgaps of the obtained materials remained at 5.1–5.3 eV, which is obviously larger than that of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>. The larger bandgap than that of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> is obviously attributed to the crystal structure of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> being different from  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>. We note that the  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> sample grown with O<sub>2</sub> as the carrier gas is dropped more abruptly than those of air and N<sub>2</sub> in ultraviolet region. This phenomenon may be caused by less defect in  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> samples carried by O<sub>2</sub>. During the film growth, extra oxygen will ensure the ideal ratio of O to Ga atomic, which results in higher crystalline quality and leads to its bandgap being close to the ideal value of

5.3 eV. The higher crystalline quality for the sample with  $O_2$  as the carrier gas is consistent with the XRD measurement results.



**Figure 3.** Optical transmission spectra and  $(\alpha hv)^2$ -hv plots of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> films.

Figure 4a shows the electron diffraction patterns of the  $Ga_2O_3/Al_2O_3$  interface for the sample with  $O_2$  as the carrier gas. The diffraction patterns of both the  $Ga_2O_3$  film and the  $Al_2O_3$  substrate are rectangular, corresponding to the corundum-structure. All of the diffraction spots of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> are situated almost in the  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> spots. The cross-sectional HR-TEM image at the  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub> interface is shown in Figure 4b. As a result of the in-plane strain, the  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub> interface is unclearly identified. At the  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub> interface, we can observe a dark area and this is induced by the in-plane compressive strain in the  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> layer because the lattice constant of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> is larger than that of  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> substrate. The TEM images for the samples, with air and N<sub>2</sub> as the carrier gases in Figures S1 and S2 also demonstrate the similar electron diffraction patterns and obvious HR-TEM lattice structures, indicating that all of the samples achieve a pure  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> phase and confirm the conclusion from the XRD measurements again. By comparing the HR-TEM images for the three different samples, the samples with  $O_2$  and air as the carrier gases show a more complete lattice structure than the sample with  $N_2$  as the carrier gas, which shows that the adequate efficient oxygen content is essential for the high quality film growth and this result is the same with that from the XRD measurement. The  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> films deposited on sapphire substrates obtain the thicknesses of 619 nm, 318 nm, and 146 nm for the samples with O<sub>2</sub>, air, and N<sub>2</sub> as the carrier gases, respectively. These thicknesses correspond to the different growth rates of 10.3 nm/min, 5.3 nm/min, and 2.4 nm/min as shown in Figure 4c. Variable-angle Spectral Ellipsometry (SE) measurements were performed at room temperature in ambient atmosphere with an electronically controlled rotating compensator and Glan Taylor polarizers (J. A. Woollam Co., Lincoln, NE, USA). The Cauchy model was used to fit the thickness of  $Ga_2O_3$  thin film. Measurements were carried out at three different incidence angles of 55°, 65°, and 75° over the 193–1000 nm wavelength range. Considering the different oxygen contents in the carrier gases, it can be concluded that more oxygen in the carrier gas can promote the growth rate greatly, and at the same time it can also guarantee the crystal quality as demonstrated by the XRD, UV-VIS and TEM measurements.



**Figure 4.** TEM images. (a) Diffraction spots of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub> for the sample with O<sub>2</sub> as the carrier gas, (b) cross-sectional  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub> interface for the sample with O<sub>2</sub> as the carrier gas. (c) Growth rate of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> films grown by air, N<sub>2</sub> and O<sub>2</sub> carrier gas.

Figure 5 shows the AFM images for the  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> films prepared by different carrier gases of air (a), N<sub>2</sub> (b), and O<sub>2</sub>(c). The films have root-mean-square (RMS) surface roughness values of 6.6 nm, 1.16 nm and 2.18 nm, respectively, measured over an area of  $5 \times 5 \ \mu\text{m}^2$ . The surface morphology for the sample grown with air as the carrier gas is rough and composed of irregular stripes, while the surface morphologies of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> thin films grown with N<sub>2</sub> and O<sub>2</sub> as the carrier gases are much smoother. It is inferred that the complex components in air induced this rough surface, and the reason behind that requires more research in the future. On the contrary, the pure carrier gas will lead to a much smoother surface. Comparing the sample grown with O<sub>2</sub> as the carrier gas to the sample with N<sub>2</sub> as the carrier gas, although the growth rate is improved by about five times, the difference between their RMS values is relatively small. This means that the rapid growth rate with O<sub>2</sub> as the carrier gas will not degrade the surface morphology. To further investigate the microstructure, we measured SEM images of all  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> films as shown in Figure 5d–f. As demonstrated in the enlarged pictures, it can be seen that the film grown with air as the carrier gas has indeed a rough surface, and the other samples show the relatively smooth surface, which confirm the results from the AFM measurement.



Figure 5. Cont.

(d)Ga2O3 film grown with Air

carrier gas





(f) Ga<sub>2</sub>O<sub>3</sub> film grown with O<sub>2</sub> carrier gas

**Figure 5.** AFM surface images of  $Ga_2O_3$  films grown with different carrier gases: (a) air (b)  $N_2$ , and (c)  $O_2$ . SEM pictures for the  $Ga_2O_3$  films grown with different carrier gases: (d) air (e)  $N_2$ , and (f)  $O_2$ .

(e) Ga<sub>2</sub>O<sub>3</sub> film grown with N<sub>2</sub>

carrier gas

To assess the elemental composition of the  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> films, XPS measurements were conducted for the epitaxial films on sapphire. The result for the sample with  $O_2$  as the carrier gas is presented in Figure 6a and only three elements (C, Ga, and O) are observed in the films. The position of the Ga  $2p_{3/2}$  and Ga 3d binding energy peaks confirms the presence of Ga<sub>2</sub>O<sub>3</sub> [18]. It can be observed from the spectra that the  $Ga2p_{3/2}$  and  $Ga2p_{1/2}$  signal peaks are located at 1118.5 eV and 1145.5 eV respectively, The energy difference between the two signal peaks is about 27 eV, which is consistent with the Ga2*p* signal peak energy difference reported in the literature [19]. The binding energy position of Ga3d signal is 21.05 eV, which is consistent with the results reported in the literature [20]. The result for the samples with air as the carrier gases is shown in Figures S3 and S4. The same peak positions appear in both samples and no significant peak shift and peak intensity changes are observed, indicating the same pure  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> and correspond to the XRD and Raman results. Figure 6b–d show the peaks of O 1s signal around 530.6 eV for the three samples. O 1s signal peaks do not show Gaussian symmetry so Gaussian fitting was carried out for the peaks. It is found that the strong O 1s signal peaks appear at 530.6 eV and 532.1 eV. The peak at 530.6 eV originated from the oxygen element in Ga–O bond, while the peak at 532.1 eV originated from the adsorption O on the sample surface [21]. Figure 5e–f show Ga 3d spectra. The chemical compositions of the surfaces were determined from the area of the Ga 3d, O 1s, and C 1s peaks taking into account the corresponding sensitivity factors [22]. The calculated atomic ratios of O to Ga for the samples grown with air, N<sub>2</sub>, and O<sub>2</sub> as the carrier gases were 1.53, 1.51, and 1.56, respectively. It reveals that more oxygen present in the carrier gas can slightly increase the O content of the film, which confirms the guess in the XRD measurement.



Figure 6. Cont.



**Figure 6.** (a) X-ray photoelectron wide spectra for the  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> sample grown with O<sub>2</sub> and N<sub>2</sub> as the carrier gas. X-ray photoelectron spectra of O 1*s* peaks for the samples grown with air (b), N<sub>2</sub> (c), and O<sub>2</sub> (d) as the carrier gases. X-ray photoelectron spectra of Ga 3*d* peaks for the samples grown with air (e), N<sub>2</sub> (f), and O<sub>2</sub> (g) as the carrier gases.

We also compared the quality of the same thickness and growth rate of the  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> films grown by different carrier gases. As shown in Figure S2 in the supplemental information, the growth rate of the  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> films was set to 10 nm/min for all samples. In Figure S3 in the supplemental information, the thicknesses of all  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> samples were 500 nm. The optical transmittance spectra show that when O<sub>2</sub> is the carrier gas, the absorption edge is closer to 234 nm of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> and the downward trend of the absorption edge is more obvious. From the AFM images, we assume that with the same growth rate and film thickness, the complex components in air induces rough surface and the pure carrier gas will lead to a much smoother surface. These results fully demonstrate that the partial pressure of oxygen in the carrier gas component can improve the quality of the film, consistent with the above results.

### 4. Conclusions

In this work, we systematically investigated the influence of different carrier gases ( $O_2$ ,  $N_2$ , and air) on the grown film quality on c-plane sapphire substrates by using the mist-CVD method. Although XRD and Raman measurements show that the pure corundum-structured  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> was successfully obtained on the c-plane sapphire substrates with single (0006) plane orientation for all three different carrier gases, the crystal quality could be greatly affected by the carrier gas. When O<sub>2</sub> was used as the carrier gas, the sample showed the smallest FWHM of about 72 arcsec obtained from the XRD rocking curves. The UV-VIS measurement also showed that the sample grown with O<sub>2</sub> as the carrier gas showed a very sharp absorption cutoff edge. It is proposed that the oxygen content in carrier gas should be responsible for all of these results. More oxygen in the carrier gas can greatly increase growth rate, and at the same time guarantee the film quality. This is further verified by the smoother surface and rapid growth rate (10.3 nm/min) for O<sub>2</sub> as the carrier gas compared to air and N<sub>2</sub> as the carrier gases (5.3 nm/min and 2.4 nm/min respectively). XPS analysis confirmed that more oxygen element can be included in epitaxial film and thus help improve the crystal quality. The proper carrier gas is essential for high quality  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> growth.

**Supplementary Materials:** The following are available online at http://www.mdpi.com/1996-1944/12/22/3670/s1, Figure S1: TEM images of the sample grown with air as the carrier gas. (1) Cross-sectional  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub> interface, (2) diffraction spots of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub>, Figure S2: TEM images of the sample grown with N<sub>2</sub> as the carrier gas. (1) Cross-sectional  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub> interface, (2) diffraction spots of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub> interface, (2) diffraction spots of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub> interface, (2) diffraction spots of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub> interface, (2) diffraction spots of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub> interface, (2) diffraction spots of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub> interface, (2) diffraction spots of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub> interface, (2) diffraction spots of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub> interface, (2) diffraction spots of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>/ $\alpha$ -Al<sub>2</sub>O<sub>3</sub>.

X-ray photoelectron wide spectra for the  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> sample grown with air as the carrier gas, Figure S4: X-ray photoelectron wide spectra for the  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> sample grown with N<sub>2</sub> as the carrier gas.

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