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Abstract: In this paper, we propose a new numerical scheme based on a variation of the standard formulation of the Runge–Kutta method using Taylor series expansion for solving initial value problems (IVPs) in ordinary differential equations. Analytically, the accuracy, consistency, and absolute stability of the new method are discussed. It is established that the new method is consistent and stable and has third-order convergence. Numerically, we present two models involving applications from physics and engineering to illustrate the efficiency and accuracy of our new method and compare it with further pertinent techniques carried out in the same order.

Keywords: numerical methods; initial value problem; autonomous equation; local truncation error; consistency; stability

MSC: 65L05; 65L07; 65L09; 65L12; 65L20



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1. Introduction

Numerical analysis is the area of mathematics that deals with computational techniques for studying and finding solutions to math problems. It is an offshoot of computer science and mathematics that develops, analyzes, and generates algorithms for numerically solving mathematical problems. Numerical methods are typically centered on the implementation of numerical algorithms. The goal of these numerical methods is to provide systematic procedures for numerically solving mathematical problems. The ordinary differential equation (ODE) is a mathematical equation that is used in natural physical processes, chemistry, engineering, and other sciences. Ordinary differential equations are notoriously difficult to approximate analytically, so the numerical treatment is crucial because it offers a powerful alternate solution method for resolving the differential equation.

We frequently use initial value problems (IVPs), such that

$$y' = f(x, y(x)), y(x_0) = y_0$$
 (1)

where *x* is the independent variable, which might also indicate time in physical problems, and y(x) is the solution. Furthermore, because y(x) can be a vector-valued function with N-dimensions, the domain and range of *f* and the solution *y* are given by

$$f : \mathbb{R} \times \mathbb{R}^{N} \to \mathbb{R}$$
$$y : \mathbb{R} \to \mathbb{R}^{N}$$

Furthermore, Equation (1), in which f is a function both of x and y, is known as a "non-autonomous" equation. However, by simply adding an extra variable, which is

always equal to *x*, it can be easily rewritten in an equivalent "autonomous" form below, in which *f* is a function of *y* only:

$$y' = f(y(x)), y(x_0) = y_0$$
 (2)

There are numerous well-known numerical methods that can be used to solve the IVP in Equation (1). One category of these is third-order methods, like Ralston's third-order Runge–Kutta method (Ralston's scheme) [1], the third-order Runge–Kutta method (RK3 scheme) [2], the third-order Runge–Kutta method based on the arithmetic mean (ARK3 scheme) [3], the third-order Heun's method (Heun's scheme) [4], and several other methods. These numerical methods have been constructed to solve Equations (1) and (2) using different techniques such as the Taylor series expansion technique, homotopy analysis technique, quadrature formulas technique, and Adomian decomposition technique; for more details, see [5–7].

In this study, we construct a new numerical method based on a variation of the standard formulation of the Runge–Kutta method using Taylor series expansion. The rest of this paper is divided into the following sections. In Section 2, we recall some definitions and auxiliary results that we will use in our work. The derivation of the new method is described in Section 3. Section 4 provides details on the local truncation error. Section 5 discusses the stability analysis of the suggested technique. The consistency of the new method and its convergence are detailed in Section 6. Several numerical models are shown in Section 7 to show the efficiency of this method and compare it with other relevant methods. Finally, Section 8 offers the discussion and the conclusions.

2. Preliminaries

Theorem 1 ([8]). (*Existence and Uniqueness Theorem*).

Let f(x, y) and $\frac{\partial f}{\partial y}$ be continuous functions of x and y at all points (x, y) in some neighborhood of the initial point (x_0, y_0) . Then, there is a unique function y(x) defined on some interval $[x_0 - \varepsilon, x_0 + \varepsilon]$ satisfying

$$y' = f(x, y(x)), \ y(x_0) = y_0, \ x \in [x_0 - \varepsilon, \ x_0 + \varepsilon], \ \varepsilon > 0$$
(3)

Definition 1 ([9]). A Taylor sum or Taylor series is a mathematical function representation in the form of a series consisting of terms calculated using the values of the function's derivation at a specific point and given by

$$\psi_T(x_n, y_n; h) = \sum_{r=0}^{\infty} \frac{h^r}{(r+1)!} \left(\frac{\partial}{\partial x} + f\frac{\partial}{\partial y}\right)^r f(x, y)$$
(4)

Definition 2 ([10]). Let $I \subseteq \mathbb{R}$ be an open interval and $f: I \to \mathbb{R}$ be a *n*-times differentiable function at $\alpha \in I$. The Taylor polynomial $T_n^{\alpha}f$ of degree *n* at a point α of *f* is the polynomial

$$\Gamma_n^{\alpha} f(\alpha) = f(\alpha) + (x - \alpha) f'(\alpha) + \frac{(x - \alpha)^2}{2!} f''(\alpha) + \frac{(x - \alpha)^3}{3!} f'''(\alpha) + \dots + \frac{(x - \alpha)^n}{n!} f^{(n)}(\alpha)$$
(5)

Definition 3 ([11]). *The difference between the numerical solution*

$$y_{n+1} = y_n + h\psi(x_n, y_n; h) \tag{6}$$

and the exact solution $y(x_{n+1})$ is called the local truncation error (L.T.E.) for a one-step method x_n with step size h is given by

$$L.T.E. = y(x_{n+1}) - y_{n+1} \tag{7}$$

Definition 4 ([12]). The numerical method is said to be stable if there exists h > 0 for each differential equation such that changing the initial values by a fixed amount produces a bounded change in the numerical solution for all h

Definition 5 ([13]). A numerical method is said to be consistent when all step sizes $h \rightarrow 0$, as it will converge to the differential equation. In other words, we say the method is consistent if and only if

$$\lim_{h \to 0} \frac{L.T.E.}{h} = 0 \text{ for } h = \frac{b-a}{N}$$
(8)

3. Construction of the New Scheme

In order to construct new single-step methods to solve the IVP Equation (1), we rely on a variant of the standard form of the Runge–Kutta method given by

$$y_{n+1} = y_n + h\psi(x_n, y_n; h)$$

where

$$\psi(x_n, y_n; h) = \sum_{i=1}^{2m} w_i k_{i+1}$$

and

$$k_1 = f(x_n, y_n), \ k_i = f\left(x_n + c_i h, y_n + h \sum_{\substack{j=1 \ j \neq i}}^{i-1} a_{ij} k_j\right), \ i = 2, \ 3, \dots, \ 2m+1$$

and

$$\sum_{j=1}^{i-1} a_{ij} = c_i, \ i = 2, \ 3, \ \dots, \ 2m+1$$

By using the Taylor series expansion of any arbitrary function $\psi(x_n, y_n; h)$ for the non-autonomous Equation (1), we have

$$\psi_T(x_n, y_n; h) = \sum_{r=0}^{\infty} \frac{1}{(r+1)!} \left(\frac{\partial}{\partial x} + f \frac{\partial}{\partial y}\right)^r f(x, y)$$

and in the autonomous case Equation (2), $\psi_T(x_n, y_n; h)$ becomes

$$\psi_T(y_n;h) = \sum_{r=0}^{\infty} \frac{1}{(r+1)!} \left(f \frac{\partial}{\partial y} \right)^r f(y)$$

We now start with the numerical method by using the family of explicit Runge–Kutta methods listed below to solve the mentioned problem in (2).

$$y_{n+1} = y_n + h(w_1k_2 + w_2k_3) \tag{9}$$

with

$$k_{1} = f(x_{n}, y_{n}) k_{2} = f(x_{n} + c_{2}h, y_{n} + a_{1}hk_{1}) k_{3} = f(x_{n} + c_{3}h, y_{n} + h(a_{2}k_{1} + a_{3}k_{2}))$$
(10)

where $a_1 = a_{21}$, $a_2 = a_{31}$, and $a_3 = a_{32}$. We must first calculate the unknowns a_1 , a_2 , a_3 , w_2 and w_3 . Using Taylor series expansion around f(y), we obtain

$$k_1 = f(y_n) \tag{11}$$

$$k_{2} = \sum_{r=0}^{\infty} \frac{1}{r!} \left(ha_{1}k_{1} \frac{d}{dy} \right)^{r} f(y_{n})$$
(12)

$$k_3 = \sum_{r=0}^{\infty} \frac{1}{r!} \left(h(a_2k_1 + a_3k_2) \frac{d}{dy} \right)^r f(y_n)$$
(13)

By expanding k_1 , k_2 , and k_3 in Equations (11)–(13), we have

$$k_{1} = f$$

$$k_{2} = f + ha_{1}ff_{y} + \frac{h^{2}}{2!}(f^{2}a_{1}^{2}f_{yy}) + \frac{h^{2}}{3!}(a_{1}^{3}f^{3}f_{yyy}) + \cdots$$

$$k_{3} = f + h(a_{2} + a_{3})ff_{y} + \frac{h^{2}}{2!}((a_{2}^{2} + 2a_{2}a_{3} + a_{3}^{2})f^{2}f_{yy} + 2a_{1}a_{3}ff_{y}^{2})$$
(14)

$$+h^{3}\left(\left(\frac{1}{2}a_{2}a_{3}^{3}+a_{2}^{2}a_{3}+\frac{1}{6}a_{3}^{3}+\frac{1}{6}a_{2}^{3}\right)f^{3}f_{yyy}+a_{1}a_{2}a_{3}f^{2}f_{y}f_{yy}+\frac{1}{2}a_{1}^{2}a_{3}f^{2}f_{y}f_{yy}\right)+\cdots\right)$$

Substituting Equation (14) into Equation (9) yields

$$y_{n+1} = y_n + (w_1 + w_2)f h + ((a_2 + a_3) w_2 + a_1 w_1)ff_y h^2 + \frac{1}{2} \left(\left((a_2 + a_3)^2 w_2 + a_1^2 w_1 \right) ff_{yy} + 2 w_2 a_1 a_3 f_y^2 \right) f h^3 + \cdots$$
(15)

The Taylor series expansion of an exact solution $y(x_{n+1})$ is given by

$$y(x_{n+1}) = y(x_n) + fh + \left(\frac{1}{2}ff_y\right)h^2 + \frac{1}{6}\left(f^2f_{yy} + ff_y^2\right)h^3 + \frac{1}{24}\left(f^3f_{yy} + 4f^2f_yf_{yy} + ff_y^3\right)h^4 + \dots$$
(16)

The following system of equations is obtained by expressing k_1 , k_2 , and k_3 in the Taylor expansion, ignoring terms with powers of *h* higher than 3, and then substituting them into formula (15) and comparing them to Equation (16):

$$hf: w_1 + w_2 = 1$$
$$h^2 f f_y: (a_2 + a_3)w_2 + w_1 a_1 = \frac{1}{2}$$
$$h^3 f^2 f_{yy}: (a_2 + a_3)^2 w_2 + a_1^2 w_1 = \frac{1}{3}$$
$$h^3 f f_y^2: w_2 a_1 a_3 = \frac{1}{6}$$

This is a system with an infinite number of solutions comprising four equations and five unknowns. Assuming that $w_1 = \frac{3}{7}$, we obtain the optimal solution listed below:

$$w_1 = \frac{3}{7}$$
, $w_2 = \frac{4}{7}$, $a_1 = \frac{1}{6}$, $a_2 = -1$, $a_3 = \frac{7}{4}$

Additionally, from Equation (6), we obtain $c_2 = \frac{1}{6}$ and $c_3 = \frac{3}{4}$.

Thus, substituting the above results in Equations (9) and (10), we present the new method, and we call it the variation Runge–Kutta method of order three (VRK3), given as follows:

$$y_{n+1} = y_n + \frac{h}{7}(3k_2 + 4k_3) \tag{17}$$

with

$$k_{1} = f(x_{n}, y_{n}) k_{2} = f\left(x_{n} + \frac{h}{6}, y_{n} + \frac{h}{6}k_{1}\right) k_{3} = f\left(x_{n} + \frac{3h}{4}, y_{n} + h\left(-k_{1} + \frac{7}{4}k_{2}\right)\right)$$
(18)

4. Accuracy of the New Scheme

Here, the local truncation error of the proposed scheme is investigated as follows. The set of Equation (18), when expanded using Taylor expansion, yields

$$k_1 = f \tag{19}$$

$$k_2 = f + \frac{h}{6}ff_y + \frac{h^2}{72}f^2f_{yy} + \frac{h^3}{1296}f^3f_{yyy} + \dots$$
(20)

$$k_{3} = f + \frac{3h}{4}ff_{y} + h^{2}\left(\frac{9}{32}f^{2}f_{yy} + \frac{7}{24}ff_{y}^{2}\right) + h^{3}\left(\frac{9}{128}f^{3}f_{yyy} - \frac{77}{288}f^{2}f_{y}f_{yy}\right) + \cdots$$
(21)

Now, substituting Equations (19)-(21) into Equation (17) yields

$$y_{n+1} = yn + hf + \frac{1}{2}h^2 ff_y + \left(\frac{1}{6}f^2 f_{yy} + \frac{1}{6}ff_y^2\right)h^3 + \left(\frac{35}{864}f^3 f_{yyy} - \frac{11}{72}f^2 f_y f_{yy}\right)h^4 + \dots$$

Hence, from Definition 3, we have

$$L.T.E. = \left(\frac{1}{864}f(f^2f_{yyy} + 276ff_yf_{yy} + 36f_y^3))h^4 + O(h^5)\right)$$
(22)

As per Equation (22), our proposed method (VRK3) is of third order, with an *L*.*T*.*E*. of fourth order.

5. Stability Analysis of New Scheme

To test the absolute stability of the presented scheme (VRK3), we use the set of Equations (18) to derive the following:

$$k_1 = \lambda y_n \tag{23}$$

$$k_2 = \lambda y_n \left(1 + \frac{h\lambda}{6} \right) \tag{24}$$

$$k_3 = \lambda y_n \left(1 + \frac{11h\lambda}{4} + \frac{7h^2\lambda^2}{24} \right) \tag{25}$$

By substituting Equations (23)–(25) into (17) and allowing $z = h\lambda$, we obtain

$$y_{n+1} = y_n \left[1 + z + \frac{Z^2}{2} + \frac{z^3}{6} \right]$$
(26)

Then, from Equation (26), the stability polynomial is

$$R(z) = \frac{y_{n+1}}{y_n} = \left[1 + z + \frac{Z^2}{2} + \frac{z^3}{6}\right] + o\left(z^4\right)$$
(27)

Utilizing the MATLAB program, Figure 1 below graphically illustrates the absolute stability region of the Formula (27):



Figure 1. The absolute stability region of the VRK3 method.

6. Consistency of the New Scheme

To explain the consistency property of the newly proposed scheme, we adopt Definition 5. Therefore, by substituting Equation (22) into Equation (8), we obtain

$$\lim_{h \to 0} \frac{L.T.E.}{h} = \lim_{h \to 0} \frac{\left(\frac{1}{864} f(f^2 f_{yyy} + 276 f f_y f_{yy} + 36 f_y^{-3})\right)h^4 + O(h^5)}{h} = 0$$
(28)

According to Lambert [13], a numerical method is consistent if the order is bigger than one. Therefore, our new method is consistent since it is of order three.

Lambert also defines a numerical method as convergent if it is consistent and stable. Following from Equations (27) and (28), this method is consistent and stable. We conclude that the new method (VRK3) is convergent because it satisfies the consistency and stability properties.

7. Numerical Examples

In this section, we introduce two models of IVPs with varying step sizes *h* to compare the efficiency and the accuracy of the proposed new method (VRK3 scheme) with other third-order methods, like Ralston's scheme, RK3 scheme, ARK3 scheme, and Heun's scheme. Here, all calculations and figures are performed using MATLAB (R2022a) software. The numerical results are presented in Tables 1–8, and the error analysis is illustrated in Figures 2–7.

7.1. Problem 1 [14]

Take into consideration the first order IVP $y' = y - x^2 + 1$, y(0) = 0.5, with the exact solution $y = x^2 + 2x - 0.5e^x + 1$, $0 \le x \le 1$. Tables 1–3 show the results that were obtained. The graphs of absolute errors are shown in Figures 2–4. The comparison of CPU time between the VRK3 scheme and other relevant third-order schemes is shown in Table 7.

x_i	Ralston's Scheme	RK3 Scheme	ARK3 Scheme	Heun's Scheme	VRK3 Scheme
zero	zero	zero	zero	zero	zero
0.1	$6.2076 imes 10^{-6}$	$6.2076 imes 10^{-6}$	$8.9854 imes10^{-6}$	$3.4299 imes 10^{-6}$	$6.5207 imes 10^{-7}$
0.2	$1.2845 imes 10^{-5}$	$1.2845 imes 10^{-5}$	$1.8692 imes 10^{-5}$	$6.9968 imes 10^{-6}$	$1.1492 imes 10^{-6}$
0.3	$1.9932 imes 10^{-5}$	1.9932×10^{-5}	$2.9173 imes 10^{-5}$	$1.0692 imes 10^{-5}$	$1.4515 imes 10^{-6}$
0.4	$2.7492 imes 10^{-5}$	2.7492×10^{-5}	$4.0482 imes10^{-5}$	$1.4502 imes10^{-5}$	$1.5125 imes 10^{-6}$
0.5	$3.5546 imes 10^{-5}$	3.5546×10^{-5}	5.2680×10^{-5}	$1.8412 imes 10^{-5}$	1.2781×10^{-6}
0.6	$4.4113 imes10^{-5}$	$4.4113 imes 10^{-5}$	6.5826×10^{-5}	$2.2399 imes 10^{-5}$	$6.8563 imes 10^{-7}$
0.7	5.3212×10^{-5}	5.3212×10^{-5}	$7.9987 imes 10^{-5}$	$2.6437 imes 10^{-5}$	3.3777×10^{-7}
0.8	6.2861×10^{-5}	6.2861×10^{-5}	$9.5229 imes 10^{-5}$	$3.0492 imes 10^{-5}$	$1.8762 imes 10^{-6}$
0.9	$7.3074 imes 10^{-5}$	7.3074×10^{-5}	$1.1162 imes 10^{-4}$	$3.4524 imes10^{-5}$	$4.0265 imes 10^{-6}$
1	$8.3864 imes 10^{-5}$	$8.3864 imes 10^{-5}$	$1.2925 imes 10^{-4}$	$3.8482 imes10^{-5}$	$6.9006 imes 10^{-6}$

Table 1. Comparison of the absolute errors among third-order schemes in Problem 1, for $h = 10^{-1}$.

Table 2. Comparison of the absolute errors among third-order schemes in Problem 1, for $h = 5 \times 10^{-2}$.

x_i	Ralston's Scheme	RK3 Scheme	ARK3 Scheme	Heun's Scheme	VRK3 Scheme
zero	zero	zero	zero	zero	zero
0.1	$7.9184 imes10^{-7}$	$7.9184 imes 10^{-7}$	$1.1480 imes10^{-6}$	$4.3572 imes 10^{-7}$	$7.9594 imes 10^{-8}$
0.2	$1.6379 imes 10^{-6}$	$1.6379 imes 10^{-6}$	$2.3876 imes 10^{-6}$	$8.8818 imes10^{-7}$	$1.3848 imes10^{-7}$
0.3	$2.5408 imes 10^{-6}$	$2.5408 imes 10^{-6}$	$3.7254 imes10^{-6}$	$1.3561 imes 10^{-6}$	$1.7141 imes 10^{-7}$
0.4	$3.5031 imes 10^{-6}$	3.5031×10^{-6}	$5.1684 imes10^{-6}$	$1.8377 imes 10^{-6}$	1.7229×10^{-7}
0.5	$4.5273 imes 10^{-6}$	$4.5273 imes 10^{-6}$	$6.7240 imes10^{-6}$	$2.3307 imes 10^{-6}$	$1.3400 imes 10^{-7}$
0.6	$5.6159 imes 10^{-6}$	$5.6159 imes 10^{-6}$	$8.3997 imes10^{-6}$	$2.8321 imes 10^{-6}$	$4.8292 imes10^{-8}$
0.7	$6.7710 imes 10^{-6}$	6.7710×10^{-6}	$1.0204 imes 10^{-5}$	$3.3383 imes 10^{-6}$	$9.4375 imes 10^{-8}$
0.8	$7.9947 imes10^{-6}$	$7.9947 imes 10^{-6}$	$1.2144 imes 10^{-5}$	$3.8448 imes10^{-6}$	$3.0504 imes 10^{-7}$
0.9	$9.2884 imes10^{-6}$	$9.2884 imes 10^{-6}$	$1.4231 imes10^{-5}$	$4.3460 imes 10^{-6}$	$5.9642 imes 10^{-7}$
1	$1.0653 imes 10^{-5}$	1.0653×10^{-5}	1.6472×10^{-5}	$4.8351 imes 10^{-6}$	$9.8318 imes 10^{-7}$

Table 3. Comparison of the absolute errors among third-order schemes in Problem 1, for $h = 25 \times 10^{-3}$.

x_i	Ralston's Scheme	RK3 Scheme	ARK3 Scheme	Heun's Scheme	VRK3 Scheme
zero	zero	zero	zero	zero	zero
0.1	$9.9973 imes 10^{-8}$	$9.9973 imes 10^{-8}$	1.4505×10^{-7}	$5.4894 imes10^{-8}$	$9.8153 imes 10^{-9}$
0.2	$2.0675 imes 10^{-7}$	$2.0675 imes 10^{-7}$	$3.0165 imes 10^{-7}$	$1.1185 imes 10^{-7}$	$1.6954 imes 10^{-8}$
0.3	$3.2066 imes 10^{-7}$	3.2066×10^{-7}	4.7062×10^{-7}	$1.7070 imes 10^{-7}$	$2.0745 imes 10^{-8}$
0.4	$4.4202 imes10^{-7}$	4.4202×10^{-7}	6.5283×10^{-7}	2.3121×10^{-7}	$2.0405 imes 10^{-8}$
0.5	$5.7114 imes 10^{-7}$	$5.7114 imes 10^{-7}$	$8.4920 imes 10^{-7}$	2.9308×10^{-7}	$1.5023 imes 10^{-8}$
0.6	$7.0830 imes 10^{-7}$	7.0830×10^{-7}	1.0607×10^{-6}	$3.5592 imes 10^{-7}$	3.5422×10^{-9}
0.7	$8.5378 imes 10^{-7}$	$8.5378 imes 10^{-7}$	$1.2883 imes 10^{-6}$	$4.1926 imes10^{-7}$	$1.5261 imes 10^{-8}$
0.8	$1.0078 imes 10^{-6}$	1.0078×10^{-6}	1.5331×10^{-6}	$4.8250 imes 10^{-7}$	$4.2799 imes 10^{-8}$
0.9	$1.1705 imes 10^{-6}$	1.1705×10^{-6}	1.7962×10^{-6}	$5.4492 imes 10^{-7}$	$8.0702 imes 10^{-8}$
1	$1.3422 imes 10^{-6}$	1.3422×10^{-6}	$2.0787 imes 10^{-6}$	$6.0565 imes 10^{-7}$	$1.3084 imes 10^{-7}$



Figure 2. The absolute errors for numerical results in Table 1.



Figure 3. The absolute errors for numerical results in Table 2.



Figure 4. The absolute errors for numerical results in Table 3.

7.2. Problem 2 (Mixture Model)

We consider here the IVP proposed in [15], which was a model of a storage tank in an oil refinery that holds 2000 gal of gasoline with 100 lb of an additive mixed within it. To prepare for winter weather, 40 gal/min of gasoline that contains 2 lb of additive per gallon is pumped into the storage tank. The well-mixed solution is pumped out at a rate of 45 gal/min. Let *y* be the amount of additive (in pounds) in the tank at time *x*. When x = 0, we know that y = 100. The mixture process is modeled by the IVP, $y' = 80 - \frac{45}{(2000-5x)}y$, y(0) = 100, and the analytic solution, $y(x) = (2000 - 5x) \{2 - \frac{3900}{(2000)^9}(2000 - 5x)^8\}$, $0 \le x \le 1$.

Tables 4–6 show the absolute errors among third-order methods and the VRK3 scheme, with different step sizes of h = 0.1, h = 0.05, and h = 0.025. Figures 5–7 depict the graphical analysis used to support the numerical results in Tables 4–6. A comparison of CPU time between the new method and other third-order schemes is shown in Table 8.

x_i	Ralston's Scheme	RK3 Scheme	ARK3 Scheme	Heun's Scheme	VRK3 Scheme
zero	zero	zero	zero	zero	zero
0.1	$3.7246 imes 10^{-9}$	$4.1133 imes 10^{-9}$	$3.8694 imes10^{-9}$	$3.3207 imes 10^{-9}$	3.1607×10^{-9}
0.2	$7.4366 imes 10^{-9}$	$8.2127 imes 10^{-9}$	$7.7257 imes 10^{-9}$	$6.6303 imes 10^{-9}$	$6.3109 imes 10^{-9}$
0.3	$1.1136 imes 10^{-8}$	1.2298×10^{-8}	$1.1568 imes 10^{-8}$	$9.9282 imes 10^{-9}$	$9.4499 imes 10^{-9}$
0.4	$1.4821 imes10^{-8}$	1.6367×10^{-8}	$1.5397 imes 10^{-8}$	$1.3214 imes10^{-8}$	1.2577×10^{-8}
0.5	$1.8494 imes10^{-8}$	2.0424×10^{-8}	$1.9213 imes10^{-8}$	$1.6488 imes10^{-8}$	$1.5694 imes 10^{-8}$
0.6	$2.2154 imes10^{-8}$	$2.4466 imes 10^{-8}$	$2.3015 imes10^{-8}$	$1.9752 imes 10^{-8}$	$1.8800 imes 10^{-8}$
0.7	$2.5801 imes 10^{-8}$	$2.8493 imes 10^{-8}$	$2.6804 imes10^{-8}$	2.3003×10^{-8}	$2.1895 imes 10^{-8}$
0.8	$2.9435 imes10^{-8}$	3.2507×10^{-8}	3.0579×10^{-8}	$2.6243 imes 10^{-8}$	$2.4979 imes 10^{-8}$
0.9	$3.3057 imes 10^{-8}$	3.6506×10^{-8}	$3.4342 imes 10^{-8}$	$2.9472 imes 10^{-8}$	$2.8053 imes 10^{-8}$
1	$3.6665 imes 10^{-8}$	$4.0491 imes 10^{-8}$	$3.8090 imes 10^{-8}$	$3.2689 imes 10^{-8}$	$3.1114 imes 10^{-8}$

Table 4. Comparison of the absolute errors among third-order schemes in Problem 2, for $h = 10^{-1}$.

Table 5. Comparison of the absolute errors among third-order schemes in Problem 2, for $h = 5 \times 10^{-2}$.

x_i	Ralston's Scheme	RK3 Scheme	ARK3 Scheme	Heun's Scheme	VRK3 Scheme
zero	zero	zero	zero	zero	zero
0.1	$4.6501 imes 10^{-10}$	$5.1354 imes 10^{-10}$	$4.8308 imes 10^{-10}$	$4.1457 imes 10^{-10}$	$3.9459 imes 10^{-10}$
0.2	$9.2889 imes 10^{-10}$	$1.0258 imes 10^{-9}$	$9.6499 imes 10^{-10}$	$8.2821 imes 10^{-10}$	$7.8832 imes 10^{-10}$
0.3	$1.3910 imes 10^{-9}$	1.5361×10^{-9}	$1.4451 imes10^{-9}$	$1.2403 imes 10^{-9}$	$1.1805 imes10^{-9}$
0.4	$1.8506 imes 10^{-9}$	$2.0437 imes 10^{-9}$	$1.9225 imes10^{-9}$	1.6499×10^{-9}	$1.5704 imes 10^{-9}$
0.5	$2.3096 imes 10^{-9}$	2.5506×10^{-9}	$2.3994 imes10^{-9}$	$2.0592 imes 10^{-9}$	$1.9600 imes 10^{-9}$
0.6	$2.7670 imes 10^{-9}$	$3.0556 imes 10^{-9}$	$2.8745 imes10^{-9}$	2.4671×10^{-9}	$2.3482 imes 10^{-9}$
0.7	3.2222×10^{-9}	$3.5584 imes 10^{-9}$	$3.3474 imes10^{-9}$	$2.8729 imes 10^{-9}$	$2.7345 imes 10^{-9}$
0.8	$3.6763 imes 10^{-9}$	$4.0598 imes 10^{-9}$	$3.8191 imes10^{-9}$	$3.2778 imes 10^{-9}$	$3.1199 imes 10^{-9}$
0.9	$4.1288 imes10^{-9}$	$4.5595 imes 10^{-9}$	$4.2892 imes10^{-9}$	$3.6813 imes10^{-9}$	$3.5040 imes10^{-9}$
1	$4.5789 imes 10^{-9}$	$5.0566 imes 10^{-9}$	$4.7568 imes 10^{-9}$	$4.0826 imes 10^{-9}$	3.8859×10^{-9}

Table 6. Comparison of the absolute errors among third-order schemes in Problem 2, for $h = 25 \times 10^{-3}$.

x _i	Ralston's Scheme	RK3 Scheme	ARK3 Scheme	Heun's Scheme	VRK3 Scheme
zero	zero	zero	zero	zero	zero
0.1	$5.7938 imes 10^{-11}$	$6.4006 imes 10^{-11}$	$6.0197 imes 10^{-11}$	$5.1628 imes 10^{-11}$	$4.9127 imes 10^{-11}$
0.2	$1.1617 imes 10^{-10}$	$1.2828 imes 10^{-10}$	$1.2069 imes 10^{-10}$	$1.0360 imes 10^{-10}$	$9.8595 imes 10^{-11}$
0.3	$1.7408 imes 10^{-10}$	$1.9220 imes 10^{-10}$	$1.8083 imes 10^{-10}$	$1.5525 imes 10^{-10}$	$1.4776 imes 10^{-10}$
0.4	$2.3087 imes 10^{-10}$	$2.5497 imes 10^{-10}$	$2.3985 imes 10^{-10}$	$2.0580 imes 10^{-10}$	$1.9583 imes 10^{-10}$
0.5	$2.8851 imes 10^{-10}$	$3.1858 imes 10^{-10}$	$2.9971 imes 10^{-10}$	$2.5724 imes 10^{-10}$	$2.4477 imes 10^{-10}$
0.6	$3.4589 imes 10^{-10}$	$3.8193 imes 10^{-10}$	$3.5934 imes 10^{-10}$	$3.0846 imes 10^{-10}$	$2.9351 imes 10^{-10}$
0.7	$4.0254 imes 10^{-10}$	$4.4454 imes 10^{-10}$	$4.1820 imes 10^{-10}$	$3.5894 imes 10^{-10}$	$3.4157 imes 10^{-10}$
0.8	$4.5941 imes 10^{-10}$	$5.0736 imes 10^{-10}$	$4.7731 imes 10^{-10}$	$4.0967 imes 10^{-10}$	$3.8989 imes 10^{-10}$
0.9	$5.1622 imes 10^{-10}$	$5.7003 imes 10^{-10}$	$5.3629 imes 10^{-10}$	$4.6035 imes 10^{-10}$	$4.3809 imes 10^{-10}$
1	$5.7196 imes 10^{-10}$	$6.3162 imes 10^{-10}$	$5.9421 imes 10^{-10}$	$5.0994 imes 10^{-10}$	$4.8530 imes 10^{-10}$



Figure 5. The absolute errors for numerical results in Table 4.



Figure 6. The absolute errors for numerical results in Table 5.



Figure 7. The absolute errors for numerical results in Table 6.

Table 7. Comparisons of CPU time in Problem 1, for different step size

Step Size			CPU Time		
	Ralston's Scheme	RK3 Scheme	ARK3 Scheme	Heun's Scheme	VRK3 Scheme
h = 0.1	0.003325	0.003508	0.004631	0.005653	0.001558
h = 0.05	0.003407	0.004827	0.003017	0.005066	0.001027
<i>h</i> = 0.025	0.003655	0.003144	0.005021	0.004364	0.001273

Table 8. Comparisons of CPU time in Problem 2, for different step sizes *h*.

Step Size			CPU Time		
	Ralston's Scheme	RK3 Scheme	ARK3 Scheme	HEUN'S SCHEME	VRK3 Scheme
h = 0.1	0.004160	0.003116	0.004698	0.004627	0.001358
h = 0.05	0.003328	0.004782	0.004879	0.003469	0.001003
<i>h</i> = 0.025	0.003506	0.004392	0.003230	0.004234	0.001296

8. Discussion and Conclusions

In this study, we introduced an innovative third-order method designed for solving initial value problems (IVPs). Our approach is rooted in a novel adaptation of the standard formulation employed in Runge–Kutta methods, incorporating Taylor series expansion. To validate the effectiveness of this new method, we employed two distinct numerical models, effectively showcasing its fundamental capabilities. It is important to underscore that all our numerical findings, including the accompanying tables and figures, were calculated using MATLAB (R2022a) software on a dedicated computer system operating with Windows 11 Pro. The system uses an 11th Generation Intel(R) Core (TM) i7-11800H processor running at 2.30 GHz, backed by 16.0 GB of RAM (15.7 GB usable).

A comprehensive numerical assessment was conducted using Tables 1–6, which present an intricate comparison of absolute errors across various step sizes, specifically $h = 10^{-1}$, $h = 5 \times 10^{-2}$, and $h = 25 \times 10^{-3}$. Through the graphical representations found in Figures 2–7, we were able to discern that our novel method, referred to as VRK3, consistently outperformed several benchmark techniques including Ralston's scheme, RK3 scheme, ARK3 scheme, and Heun's scheme. This superiority primarily stems from the reduced local truncation error of VRK3. Additionally, our investigation revealed a significant insight regarding the impact of step size on accuracy. As we decreased the step size, the error progressively approached zero, strongly indicating that precision increased with smaller step sizes. This observation reinforces the importance of carefully selecting step sizes to achieve higher levels of accuracy in numerical solutions. Turning our attention to computational efficiency, Tables 7 and 8 provided valuable insights. The VRK3 scheme consistently demonstrated reduced CPU time compared to its counterparts, further validating its utility in practical applications. Furthermore, Figure 1 depicts the stability region of our third-order VRK3 scheme, establishing its equivalence to similar methodologies. Importantly, we substantiated the convergence of our VRK3 scheme, as it satisfies both the consistency and stability criteria.

In conclusion, our newly proposed third-order method exhibits a commendable blend of efficiency and reliability. The method's stability and high accuracy render it particularly robust for a wide range of applications. This research contributes to the field of numerical methods for IVPs by presenting an innovative approach that holds promise for improving computational accuracy and efficiency. Future research directions might explore the extension of this method to more complex problems or its integration into broader computational frameworks.

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