



# Article Hydraulic Capacity and Efficiency of a Low-Speed Nonpressurized Coil Pump

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**Abstract:** The paper presents the results of an investigation on hydraulics of a low-speed coil pump for transport of liquids. One of the pump's advantages is its wide range of potential inclination angles for its rotating shaft, from the horizontal to an almost vertical position. A simplified hydraulic model was developed based on kinematic and geometrical considerations to determine the pump capacity. The model was verified under laboratory conditions using a low-speed coil pump composed of transparent PVC tube (15 mm outer diameter) wound around a cylindrical drum (104 mm external diameter; 550 mm long). Laboratory tests were performed for three angles of inclination of the axis of rotation (20°, 40°, and 60°) and four rotational speeds (10, 20, 30, and 40 rpm). The results of the tests showed satisfactory agreement with the hydraulic model predictions. Energetic efficiency was estimated on the base of electric power measurements and difference of water levels in the two arms of rotating transparent torus, partly filled with water. The hydraulic efficiency of the coil pump is increasing with decreasing rotational speed.

Keywords: Archimedean screw; hydraulic capacity; coil pump efficiency; pressure drop in helical coil

# 1. Introduction

The Archimedean screw is one of the oldest devices for water raising. It was described, although not necessarily invented by Archimedes [1]. The screw rotates inside a tube or trough, pushing portions of water ahead. The screw axis is inclined and its lower end picks up water from the source, whereas the upper end discharges into an irrigation ditch. The optimum angle of inclination, usually in the range of 30° to 40°, depends on the pitch and the diameter of the internal helix. The main disadvantage of the Archimedean screw is water loss through the slots between the screw and its casing [2]. Its hydraulic efficiency may reach up to 80%, including when used as a small hydro-power turbine [3–5]. In the latter case, the angle of inclination ranges between 20° and 30° to gain maximum torque.

A similar principle is applied in spiral pumps. The spiral pump, invented most probably by Wirtz (1746), is composed of a tube wound in the form of an Archimedean spiral. Its outer, broader end picks up water from the source during the rotation of the shaft. Its inner end, located at the axis of rotation, discharges liquid into a channel or pipeline. While it is easy to construct, the critical part of the spiral pump is its rotary fitting (swivel). It provides a relatively watertight seal to prevent water and pressure losses. Unfortunately, it often leaks and generates friction and losses of energy. An overall efficiency of up to 75% can be achieved for a well-designed spiral pump [6]. Apart from liquid lifting, the spiral pump may be used to remove floating pollutants from the surface of the sea [7,8]. Structures based on the idea of a spiral pump may also be applied in wastewater treatment facilities or in heat exchangers and chemical reactors thanks to efficient mixing by secondary currents [5,9–12]. A coil pump is based on the same principle as the spiral pump, but the tube is wound around either a conical drum or a cylindrical drum, as a result creating a helix. The principle is that water is picked up by the

submerged end of the tube each time it dips below the surface; however, due to the dynamic effects, the entrapped air is pressurized. The maximum pressure builds up at the outlet, while the liquid plugs rotate in a manner that suggests that they are creating a form of multiple (cascading) manometer. This type of pump, although in a nonpressurized version, was originally described by the Roman architect and engineer Vitruvius in the first century B.C. It has attracted considerable renewed interest recently, with research projects concerning this design conducted at several universities. Today the coil pump is recommended for use in stream (river current) powered irrigation pumps [13–15]. The coil pump powered by the water flow is sometimes called a sling pump. It can also be driven using other nonconventional energy sources (pedal, animal, wind or solar) or even manually. Due to the presence of air plugs, water or wastewater is partly aerated; therefore the pump can also be used for wastewater treatment [16,17], particularly for wastewater pretreated mechanically in septic tanks, e.g., in a small diameter sewerage system [18].

The majority of studies on coil pumps were performed using designs with a horizontal axis of the cylindrical drum. We have found only three studies on the role of inclination of the axis of rotation [19–21]. Inclining the axis of rotation allows elimination of the rotating seal, since the coiled tube and the inclined force main are now joined to form one component. This also means that the force main may be used as the drive shaft, enabling the pump to be driven from dry land as an alternative to the drum being driven by the stream power. The inclined coil pump is therefore more versatile in terms of the source of power.

Koetsier and Blauwendraat [22] provided a detailed review of historical attempts to formulate a theory of the Archimedean screw pump. Among the contributors they listed such famous names as Archimedes himself, L. da Vinci, G. Cardano, G. Galilei, D. Bernoulli, and J. Weisbach. Especially interesting and useful is a formula derived by D. Bernoulli, using which allows to calculate the height above the horizontal plane  $H(\alpha)$  of a point P on the cylindrical helix (Figure 1) as the function of  $\alpha$ , the angle of rotation during the generation of the helix [22], in the form of

$$H(\alpha) = R \left[ \alpha \tan \psi \sin \beta + \cos \beta \left( 1 + \cos \alpha \right) \right]$$
(1)

where  $\psi$  is the inclination of the cylindrical helix to the pump baseline and  $\beta$  is the angle of inclination of the axis of rotation to the horizontal.

A diagram of Equation (1) for the selected inclination angles  $\beta$  at the helix inclination angle  $\psi = 3^{\circ}$  (typical for closely wound tube) is shown in Figure 2.

In the interval  $0 < \alpha < \pi$ , corresponding to the first section of the cylindrical helix, Equation (1) has no extrema when  $\psi > 90 - \beta$ ; the helix is then an ascending curve and water does not enter the tube. The same result was obtained by Galilei, who deduced it somewhat differently. However, when  $\psi < 90 - \beta$ , Equation (1) has two extrema: there is a maximum corresponding to point *o* and a minimum corresponding to point *p*. Bernoulli remarked that in this case the volume of water in one pocket is determined by the section *opq* of the helix (Figure 1), with *q* being on the same level as *o*. He stated that the length of this pocket cannot be determined algebraically, but may be approximated in every specific case.



**Figure 1.** A diagram of a helix *opq* wound around a cylindrical drum (adapted from Koetsier and Blauwendraat (2004)).



**Figure 2.** Heights of points on the cylindrical helix on inclined ( $\beta = 40^{\circ}$ ), vertical ( $\beta = 90^{\circ}$ ), and horizontal ( $\beta = 0^{\circ}$ ) cylinder of radius R = 5.75 cm at the helix inclination angle  $\psi = 3^{\circ}$  according to Equation (1).

The maximum value of  $H(\alpha)$ , corresponding to point *o*, can be found from the first derivative of the function expressed by Equation (1), equaled to zero:

$$\frac{dH(\alpha)}{d\alpha} = R \left( \tan \psi \sin \beta - \cos \beta \sin \alpha \right) = 0$$
(2)

hence

$$\alpha_o = \sin^{-1} \left( \frac{\tan \psi}{\cot \beta} \right) \tag{3}$$

A theoretical explanation and description for the formation of gas and liquid plugs in the classical, pressurized coil pump was provided by Mortimer and Annable [15]. The theory was successfully confirmed by those researchers as well as Kassab et al. [23]. Mortimer [16] thoroughly investigated the coil pump in two variants: as a lift and as a suction (single and double) pump.

Kassab et al. [24] stated that the application of a two- or three-layer coil pump improves its performance (head and discharge), mainly due to the upper tube layer. Increasing the tube and drum diameter also led to improved coil pump performance. Rorres [25] found the optimal inner radius and pitch that maximize the volume of water lifted in one turn of the Archimedean screw. Optimal parameters' values were found to be close to those used in a screw described by Vitruvius and with values used in the design of modern Archimedean screw pumps.

This paper presents investigations of a low-speed coil pump in the form of tube wound helically around a cylindrical drum, without any force main on the outlet. Alternatively, the pump can be called a tubular (helical) hydraulic elevator [26]. Such a design eliminates leakage; hence, its volumetric efficiency is theoretically equal to 100%. The rotation speed is so low ( $\leq$ 40 rpm) that the air pressure between the liquid plugs (pockets) is close to the barometric one. The pump only allows the lifting of water over a small height (up to a few meters); therefore, in some ancient structures, the Romans applied several stages of lifting, as in the case of the Augsburg Machine [22]. The theory of the pump work is relatively simple, but, to our best knowledge, it has not been experimentally verified for low-speed coil pumps. Though so old, the coil pump design is continuously being improved and its prospects for innovation and commercialization are promising.

### 2. Materials and Methods

#### 2.1. Theoretical Considerations

In order to calculate the capacity of a nonpressurized coil pump, the following geometrical analysis was carried out.

The length of a helix curve, created by a polar radius *R* turning around at an angle  $\alpha$ , (Figure 1) can be expressed as

$$L = \int_0^\alpha \sqrt{R^2 + \frac{l^2}{4\pi^2}} \, d\alpha = \frac{\sqrt{4\pi^2 R^2 + l^2}}{2\pi} \, \alpha \tag{4}$$

where *R* is the radius of the cylindrical helix, *l* is the lead and  $\alpha$  is the angle of rotation.

Therefore the length of one coil is equal to

$$L_1 = \sqrt{4 \, \pi^2 R^2 + l^2} \tag{5}$$

A helix lead is defined as the length of one complete helix turn (coil), measured parallel to the axis of the helix. The length of a lead is equal to the product of its pitch (the lead at the closest winding) and the number of starts. The relative inner volume of one coil of a tube (for  $\alpha = 2\pi$ ) with the same circular cross-section depends on the following length ratio.

$$\frac{V_{1l}}{V_1} = \frac{L_1}{2\,\pi\,R} = \sqrt{1 + \left(\frac{l}{2\,\pi\,R}\right)^2} \tag{6}$$

where  $V_1$  is the volume of a circular coil, i.e., torus, (S = 0) of radius R,  $V_{1l}$  is the volume of an expanded cylindrical coil of length  $L_1$  (l > 0). For l/R < 1  $V_{1l}/V_1 < 1.013$ , i.e., the relative volume difference is less than 1.3%.

Assuming the closest tube winding, the length of a cylindrical helix curve is approximately equal to the length of a circular arc, namely,

$$L \approx \left(S + \frac{c^2}{4s}\right) \left(\frac{4cS}{4S^2 + c^2}\right) \tag{7}$$

where *c* is the chord length and *S* is the sagitta of one vertical or tilted liquid pocket. The latter is approximately equal to

$$S \approx \frac{h}{\cos(\beta - 0.5\,\psi)} \tag{8}$$

where  $\psi$  is the inclination of the cylindrical helix to the pump baseline; for the closest winding with a pitch at one start

$$\psi = \tan^{-1} \left[ \frac{r}{\pi (R+r)} \right]$$
(9)

where r is the outer diameter of the coil tube.

To fill a helical pocket with liquid and to carry it upwards the following conditions need to be fulfilled:

- $\beta > 90 \psi$ , where the angle unit is deg.
- the tube inlet must be submerged intermittently below the liquid level to create a plug flow.

The chord length of the circular pocket in its lower half ( $S \le R$ ) is equal to

$$c = 2 \sqrt{2 S R - S^2} \tag{10}$$

Similarly, the chord length in the upper half (S > R) may be calculated treating the sagitta as a measure of gas content at the coil top.

We assume that a single tube fill depth is equal to the submergence depth *h* (Figure 3), which is possible when the rotational speed is sufficiently low. The liquid volume in one partly filled coil with l/R < 1 and h > 2a can then be calculated as

$$V_h = 2 \pi R a^2 \cos^{-1}\left(\frac{R-h}{R}\right)$$
 (11)

where *a* is the inner radius of the tube.

Equation (11) is valid for a < h < 2 R - a.



**Figure 3.** Cross-section of one nearly vertical coil ( $\beta = 0$ ) at the lower drum end: (**a**) still and (**b**) rotating.

More generally, using Equation (7), the pump yield can be expressed by

$$Q_p = n V_h = n \pi a^2 L = n \pi a^2 \left(s + \frac{c^2}{4 s}\right) \left(\frac{4 c s}{4 s^2 + c^2}\right)$$
(12)

where *n* is rotational speed, expressed in revolutions per time unit.

A certain unusual mode of operation for the nonpressurized coil pump is to elevate small portions of liquid, which do not create plugs. The air can flow over the surface of successive liquid portions, because it is not blocked by liquid plugs. For small fillings of the coil tube the following method of liquid pocket volume estimation was applied. For a given cross-section, e.g., *B-B* in Figure 4, liquid depth in the tube is equal to

$$\delta = d - y = d - \frac{x^2}{2(R+a)}$$
(13)

in the dimensionless form

$$\delta^* = \frac{\delta}{2a} = \frac{d}{2a} - \frac{x^2}{4a(R+a)}$$
(14)

The pocket of liquid extends from  $-x_0$  to  $x_0$ , therefore assuming symmetry its volume  $V_p$  can be estimated as

$$V_p = 2\pi a^2 \int_0^{x_0} \varepsilon dx \tag{15}$$

where

$$\varepsilon = \frac{(1-2\,\delta^*) - 2\,(1-2\delta^*)\,\sqrt{\delta^*(1-\delta^*)}}{\pi}$$
(16)

Function (15) must be integrated numerically.



Figure 4. Vertical and cross-section *B*-*B* of a tube coil partly filled with liquid.

### 2.2. Construction of the Studied Helical Coil Pump

The tested pump (Figure 5) was made of transparent PVC tube of outer/inner diameter 15/12 mm (1), wound around a cylindrical drum made of PVC (9) of external diameter 104 mm and 550 mm long. The number of coils on whole length of drum was 32. The inner tube diameter was verified by a volumetric method—as a result the mean inner tube diameter was determined to be 11.2 mm (a = 5.6 mm) instead of the factory value of 12.0 mm. The coil pump was installed in a tank with a capacity of 6.8 dm<sup>3</sup> (10) filled with water to a given level, controlled by the overflow (8). Supply water was pumped by a small centrifugal pump through the inlet (7) to maintain a constant water level. With the line of sight along the ascending drum axis, the sense of tube winding was opposite to the direction of the drum revolutions (called counter-winding) to provide a forward (upward) flow in the tube (1). The rotating drum (3) was driven by the 12VDC motor (4) and the transmission with the reduction ratio of 0.44. When lifting the liquid through the coil-pump, the water level in the following coils, can be observed. The liquid outflowing from the uppermost coil was captured by collar (10) and directed to a container (5). The container, together with the water flowing out of the coil pump, was weighed using the electronic balance (6).

To determine input power and to assess energetic efficiency of the coil pump, input voltage and current were measured using voltmeter and ammeter at the power supply terminals.



**Figure 5.** Side view and longitudinal and cross-sections of the experimental set-up: 1: tubes; 2: stiffening rods; 3: shaft; 4: electric motor; 5: container; 6: electronic balance; 7: inlet of water; 8: overflow; 9: drum; 10: collar; 11: tank.

## 2.3. Experiment

The pump capacity was tested at four rotational speeds, i.e., 10, 20, 30, and 40 rpm. The planned submergence depth of the lower part of the drum was provided using a small centrifugal pump and the overflow, through which the surplus liquid was removed. Thus, the inlet submergence was kept at a constant value h = 5.75 cm. During the successive revolutions, the tube took the liquid and air portions and lifted them up. The drum revolutions were maintained until outflow from the upper outlet filled partly the container (5). The collected water mass measurements, using the electronic balance (6), were sent to the computer and recorded with the sampling frequency 2 s.p.s.

The hydraulic model was validated by 12 experiments (see Table 1). The influence of the angle of inclination of the drum on the submergence conditions is shown in Figure 6.

			1			0						
No. of Experiment	1	2	3	4	5	6	7	8	9	10	11	12
Inclination angle, $\beta$ , deg.	20			40			60					
Rotational speed, n, rpm	10	20	30	40	10	20	30	40	10	20	30	40

Table 1. Experimental design.



**Figure 6.** Side views of partly filled coils wound on the motionless drum inclined at angle (**a**)  $\beta = 20^{\circ}$ , (**b**)  $\beta = 40^{\circ}$ , (**c**)  $\beta = 60^{\circ}$ , (**d**)  $\beta = 83^{\circ}$  to the horizontal plane; the last shows the filling of coils caused by the liquid back flow after exceeding a critical angle  $\beta = 90 - \psi$ .

In Figure 6a–c, the coil pump is shown at inclination angles 20°, 40°, and 60°. The angle of inclination of the drum up to a critical value  $\beta = 90 - \psi$  does not affect the pump capacity.

Water level difference  $\Delta h$  in two arms of rotating torus of radius R = 8.8 cm and inner tube diameter a = 6 mm was determined basically for checking the applied head-loss calculation procedure. In experiment the difference was measured from photographs.

A portion of water corresponding to filling half of the first coil was introduced into the initially empty coil pump. To increase the contrast, the water was stained with fluorescein. Then the motor was started at a certain speed and the flow of the liquid portion was observed over successive coils. The liquid flow phenomenon was recorded using a 1080 HD video camera (of a resolution matrix of  $1920 \times 1080$  pixels) at 30 fps. To determine the scale of the image on the surface of the coils, a millimeter scale was applied. The optical axis of the lens was positioned at the liquid level in the coils, and the tangent plane to the curved surface of the coil was perpendicular to the optical axis. This reduced the influence of the curvature of the shaft on the angle of observation and the need to take into account geometrical distortions. Video in mp4 format was processed into pictures (with a resolution of 96 dpi) on which, using the ImageJ software, the number of pixels between the dynamic, lower liquid meniscus, and the static one was measured. The calculated scale factor was 12.45 pixels per 1 mm of the observed object (around the lens axis).

Reading liquid levels was based on the analysis of images obtained from films, always for the fourth and seventh coils. The level reading position was marked on the circumference of the coils. Knowing the number of pixels and the determined image scale, the dynamic and static levels were read and the difference between the two levels was calculated. Reading the levels for specific places on the coils allowed making measurements independent of the influence of imperfections in the winding of the total 32 coils.

## 3. Results and Discussion

## 3.1. Capacity

Theoretical analysis using Equation (9) showed that, at the closest tube winding, the angle of inclination of the drum axis to the horizontal can reach its maximum values due to the smallest value of the helix angle  $\psi$  (Figure 2). In the tested coil pump, the maximum permissible inclination angle of the drum  $\beta = 90 - \psi = 86.8^{\circ}$  was provided for the lead l = DN = 15 mm (Figure 6d). At larger  $\beta$ , the liquid flowed down completely from the first (inlet) coil.

Calculations of the pump capacity were made using Equation (12) for the values given in Table 1. A comparison of calculated and measured values of the coil pump capacity shown in Figure 7 indicates their good agreement. To evaluate the hydraulic model, a coefficient of variation, based on the relative squared error, was calculated according to the following formula [27].

$$CV(RMSE) = \sqrt{\frac{\sum_{i=1}^{m} (X_c - X_{ei})^2}{m X_{ea}^2}}$$
 (17)

where  $X_c$  is the value calculated using Equation (12),  $X_{ei}$  is the *i*-th measured value,  $X_{ea}$  is the average measured value, and *m* is the number of measurements.



Figure 7. Comparison of calculated ( $Q_{p cal}$ —Equation (12)) and measured ( $Q_{p meas}$ ) coil pump yields.

Calculated values of the coefficient of variation are given in Table 2. Considering the above mentioned uncertainties, CV (RMSE) < 5% indicates a relatively good agreement between the theory and experiments. In all experiments the differences between the results of experiments and calculations fell within the limits of measurement errors. It can be stated that in the tested range the rotational speed of the drum had no significant effect on the effectiveness of coil filling. Small fluctuations of the submergence depth of the lowest coils may significantly change the volume of lifted water. The observations also confirmed the equality in the levels/volumes of water in successive coils, provided that the submergence of the lowest coil during drum rotation was constant. Indirectly, this is a confirmation for the assumption that the air plugs are not compressed by the liquid plugs. At the same time, this statement is limited only to nonpressurized pumps, i.e., the head  $H_m$  depends only on the drum length and its inclination.

Rotational Speed, n, rpm	10	20	30	40
Number of measurements, <i>m</i>	96	96	96	96
Pump yield, $Q_{p meas}$ , m <sup>3</sup> /s	$3.3  imes 10^{-6}$	$6.6  imes 10^{-6}$	$9.8  imes 10^{-6}$	$13.2 \times 10^{-6}$
Pump yield, $Q_{p cal}$ , m <sup>3</sup> /s	$3.4  imes 10^{-6}$	$6.8  imes 10^{-6}$	$10.2 \times 10^{-6}$	$13.6\times10^{-6}$
CV(RSME), % (see Equation (17))	0.20	0.16	0.59	4.15

**Table 2.** Comparison of calculated ( $Q_{p cal}$ —Equation (12) and measured ( $Q_{p meas}$ ) coil pump yields.

Relative pump yield depends to a differentiated degree on drum and helix inclinations (Figure 8). At relative small helix angles the dependence is rather weak. With increasing helix angles the relative pump yield dramatically decreases for the highest ( $\beta = 60^\circ$ ) drum inclination but it may even increase for the smallest ( $\beta = 20^\circ$ ) drum inclination to horizontal.



Figure 8. Calculated relative pump yield for different drum and helix inclinations.

To determine the fill time of coils simple experiments were conducted: one coil of the same tube as in the pump (a = 6 mm), wound around a short piece of pipe of 100 mm in diameter, was rapidly submerged under water and the fill time of the coil was measured. Ten experiments showed the average fill time of 0.72 ± 0.10 s. That value was divided by two to estimate the shorter fill time of a half coil (Figure 9) and to compare it with the duration of a half-turn, during which the tube inlet is submerged. It can be seen that to provide proper filling, the maximum rotational speed should be not greater than 60–70 rpm. A slightly lower limit (50 rpm) is also valid for Archimedean screws; if the screw is rotated much faster, turbulence and sloshing prevent the buckets from being filled and the screw churns the water in the lower pool rather than lifting it [26].



Rotational speed, n, rpm

**Figure 9.** Theoretical fill time of half coil (R = 57.5 mm) of tube of inner diameter 12 mm vs. half-turn period as a function of rotational speed. Dotted lines show the confidence interval based on sample standard deviation (m = 10).

In our case, as can be deduced from Equation (4), the greater the tube lead is, the greater the volume of one coil, but the number of coils along a given drum length is smaller. In the case of the studied helices, their lead l = 0.015 m, i.e., one tube created 32 coils around a drum of 0.5 m in length. It is worth to note that liquid plugs "flow" with a velocity equal to the peripheral speed  $2\pi Rn$ . Flow velocities inside the tubes reached 14–15 cm/s at the rotational speed n = 20 rpm. Such a low linear wastewater velocity should not be harmful to a biofilm on wall surfaces of the tubes, as Hochheimer and Wheaton [28] assessed that velocities < 30 cm/s are acceptable for safe attachment of biofilm on rotating biological contactors.

## 3.2. Efficiency and Torque

Typically, overall efficiency of a fluid flow process involving a pump is related to the hydraulic, mechanical, and volumetric losses in the pump.

The hydraulic efficiency of a coil pump (excluding the bearing friction) is defined as the power ratio, as follows

$$\eta_H = \frac{P_l}{P_l + P_f} \tag{18}$$

where  $P_l$  is the power needed to lift water in the gravity field, and  $P_f$  is the power needed to surmount fluid friction within the coil pump.

The power needed to lift water by the coil pump can be estimated as

$$P_l = \rho_w \, g \, H \, Q_p \tag{19}$$

It is worth to noting that Equation (19) is equivalent to equation

$$P_1 = 2 \pi n N M \tag{20}$$

where *n* is rotational speed, rps; *N* is the number of pockets; and *M* is the moment that the weight of the water *G* exerts about the axis of the cylindrical helix, which is expressed in the following form given by Bernoulli [23].

$$M = G R \tan \psi \sin \beta \tag{21}$$

where *G* is the weight of the water in the whole pocket *opq* (Figure 1).

The power needed to surmount fluid friction within the coil pump can be estimated as

$$P_f = 2\pi R n N F_s \tag{22}$$

Assuming that the fluid velocity in the coil tube is equal to the peripheral speed, the former can be expressed as

$$v = 2 \pi R n \tag{23}$$

The shear force in the coil tube can be expressed as

$$F_s = 2\pi a L_p \tau \tag{24}$$

where  $L_p$  is length of the liquid pocket, m, and  $\tau$  is wall shear stress, N/m<sup>2</sup>.

The shear stress in coil tube can be expressed as

$$\tau = \rho_w \frac{v^2}{8} f_c \tag{25}$$

where  $f_c$  is the friction factor for liquid flow in a coil tube, -.

Numerous publications [29–33] indicate that the friction factors values for a coil tubes are higher than their counterparts in straight pipelines. Therefore, the use of classical equations to calculate the value of friction factors will give underestimated values. Nevertheless, the friction factor for a coil tube is calculated on the basis of the friction factor for a straight pipe. Several equations have been proposed to calculate the friction factor at laminar flow in a coil tube. For example, White [29] proposed the following formula, which was obtained experimentally.

$$f_c = f_s \left\{ 1 - \left[ 1 - \left( \frac{11.6}{De} \right)^{0.45} \right]^{2.2} \right\}$$
(26)

Similarly, Ito [30] proposed the following empirical formula.

$$f_c = f_s \left\{ 0.1033 D e^{0.5} \left[ \left( 1 + \frac{1.729}{D e} \right)^{0.5} - \left( \frac{1.729}{D e} \right) \right]^{-3} \right\}$$
(27)

where  $f_s$  (-) is the friction factor for a straight pipe (for laminar flow  $f_s = 64/Re$ , Re = 2av/v—Reynolds number, -; v is kinematic viscosity of the fluid, m<sup>2</sup>/s; and  $De = Re(a/R)^{0.5}$  is Dean number, -.

On the other hand, Hart et al. [31] developed the following formula.

$$f_c = f_s \left( 1 + \frac{0.09 \ De^{1.5}}{70 + De} \right) \tag{28}$$

Other formulae for the friction factor at laminar flow in coil tubes can be found in the review provided by Ghobadi and Muzychka [32] and for transitional and turbulent flows by Fsadni and Whitty [33].

Comparison of values of head losses  $\Delta h$ , calculated by Darcy-Weisbach formula, with relevant experimental data (Figure 10) show acceptable agreement.

Results of calculations using Equations (18)–(26), made for a single start coil pump (R = 57.5 mm, a = 6 mm) for chosen values of rotational speed n, at the static head H = 1000 mm, are shown in Figure 10. The highest hydraulic efficiency values were achieved for the inclination angle  $\beta \approx 70^{\circ}$  (Figure 11). With the increase of rotational speed, a decrease in efficiency can be observed, which results from the linear increase in shear stress along with the increase of fluid velocity. In contrast, in the case of centrifugal pumps the hydraulic efficiency decreases with rotational speed [34–36]. For the inclination angle  $\beta \approx 20^{\circ}$ , the efficiencies are comparable to those of the Archimedean screw [3,4]. At smaller inclination angles they are even lower due to the increasing drum length and relatively large number of coils. It should be noted that energetic efficiency in the case of wastewater transport may decrease in time as a result of biofilm formation on the wall of coil tubes.



**Figure 10.** Comparison of calculated frictional head losses and those measured in a rotating torus of radius R = 8.8 mm (black dots).



**Figure 11.** Hydraulic (theoretical) efficiency of a single start coil pump (R = 57.5 mm; a = 6 mm; r = 7.5 mm) for chosen values of rotational speed *n*, static head H = 1000 mm, and  $\psi = 3 \text{ deg}$ .

Components of power balance for the experimental coil pump inclined at  $\beta = 40^{\circ}$  are presented in Figure 12. It can be seen that the hydraulic efficiency of the coil pump is increasing with decreasing rotational speed, however the total efficiency is relatively low due to high electric and mechanical losses. The values of measured mechanical and electrical losses are specific to the tested structure of the coil pump. The use of other bearings on the shaft axis, transmission and engine will change the mutual proportions of the three power components.

Comparing our pump with famous ancient constructions of tubular Archimedean pumps, e.g., with the Augsburg Machine [23], one can observe that in the latter case the drum inclination angle was relatively small ( $\beta \approx 24^\circ$ ), the lead ratio was relatively high ( $l/D \approx 6$ ,  $\psi = 45^\circ$ ), and the drum length to the tube outer diameter ratio was relatively low ( $L_D/D \approx 18$ ). However, probably due to ancient heavy or fragile materials and primitive power transmissions, the ancient architects decided to build the machine in seven stages.

Providing several starts would increase its lead to pitch ratio, helix inclination angle  $\psi$  and power. This problem is similar to determination of the optimal number of blades in Archimedean screw-pumps, which typically have 2–8 blades [26].



**Figure 12.** Components of power balance for the experimental coil pump inclined at  $\beta = 40^{\circ}$ .

Further studies in this direction as well as to optimize the device for a specific application are needed.

Low-speed coil pumps, similarly to Archimedean screw pumps [37], have a great potential to be applied in remote places and be powered by renewable energy sources (wind and solar energy). They may also be driven by running water (sling pumps).

# 3.3. Scale Effects

The presented results of experimental research were carried out in a laboratory scale reactor. Due to the small diameter of the PVC tube (15 mm) from which the coils were made, it was checked whether there is a capillary effect. For this purpose, a tube section (15 cm) from identical material was moistened and immersed vertically in a beaker with water at an ambient temperature of 25 °C. The inner meniscus of height < 1 mm has occurred at the same level as the outer one. With coils with smaller inner diameters (several mm), the meniscus may deform at the liquid/gas interface, a phenomenon known as pinning [38]. Dimensionless numbers in relationships for  $f_c$  (*Re* and *De*) can be treated as dynamic similarity criteria which regard themselves main scale effects.

# 4. Conclusions

Based on the idea of the Archimedean screw and our own experiments, a hydraulic model of a low-speed coil pump was developed. A new formula to calculate the pump capacity was successfully validated using measurement results obtained from a laboratory physical model at rotational speeds of 10 and 40 rpm.

The performed experiments confirmed the assumption that with a constant level of submergence of the drum, successive coils with a constant radius are filled to the same level regardless of the coil length. In the tested range of rotational speed of the drum this did not affect the effectiveness of coil filling. Additional experiments allowed us to determine the speed of free filling of coils. At speeds greater than 60–70 rpm filling may be hindered, as the liquid will not be able to fill the coils of inner radius a = 6 mm in time.

With an open tube outlet and a constant pump drum length the pump static head increases proportionally to the tangent of the drum inclination angle. However, the latter cannot be greater than the angle of a helix; therefore the maximum lead depends on the drum inclination angle. In practice, an inclination angle of the drum axis (to the horizontal) close to 87° (at the closest tube winding) makes

it possible to install the coil pump almost vertically in standard manholes of sewerage systems. Such installations in small diameter gravity sewerage systems can replace traditional pumping stations and may play the role of in-line treatment units. The low-speed nonpressurized coil pump is a leak-proof variant of an Archimedean screw. Its main disadvantage, namely the lack of direct access to the inner space of the tubes wound around the drum, can be overcome by applying the Archimedean screw with a helically wound lateral strip.

Relative pump yield depends to a differentiated degree on drum and helix inclinations. At relative small helix angles the dependence is rather weak. With increasing helix angles the relative pump yield dramatically decreases for the highest ( $\beta = 60^{\circ}$ ) drum inclination but it may even increase for the smallest ( $\beta = 20^{\circ}$ ) drum inclination to horizontal.

Low-speed coil pumps, similarly to Archimedean screw pumps [29], have great potential to be applied in remote places and be powered by renewable energy sources (wind and solar energy). They may also be driven by running water (sling pumps).

Further studies are needed to improve the efficiency as well as optimize the device for specific applications.

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## List of symbols

а	inner radius of the coil tube, m
С	chord length, m
D	diameter of cylindrical helix, m
D	maximum fluid depth in partially filled coil tube, m
De	Dean number,
$f_c$	friction factor for coil tube, -
$f_s$	friction factor for a straight pipe, -
$F_s$	shear force in the coil tube, N
8	gravitational constant, m s <sup>-2</sup>
h	submergence depth of the left of tube inlet, m
Η	static head or height of the helix, m
$H_m$	mean static head, m
1	lead, m
L	length of helical curve, m
$L_1$	length of cylindrical coil (torus), m
$L_p$	length of liquid pocket, m
Μ	moment of force G about the axis of cylindrical helix, Nm
т	number of measurements, -
Ν	number of pockets, -
п	rotational speed, rps or rpm
$P_f$	power needed to surmount fluid friction within the coil pump, W
$P_1$	power needed to lift water in the gravity field, W
$Q_p$	pump yield, m <sup>3</sup> s <sup>-1</sup>
r	outer radius of the tube, m
R	radius of cylindrical helix or torus, m

S	sagitta of a circular pocket, m
S	length of circular arc, m
υ	fluid velocity in coil tube, m $s^{-1}$
$V_1$	volume of a circular coil (torus), m <sup>3</sup>
$V_{1l}$	nner volume of one helical tube coil, m <sup>3</sup>
V <sub>h</sub>	liquid volume in one partly filled coil, m <sup>3</sup>
$V_p$	volume of liquid in a shallow pocket ( $d < 2a$ ), m <sup>3</sup>
X <sub>c</sub>	value calculated using the model,
$X_{ea}$	average measured value,
X <sub>ei</sub>	<i>i</i> -th measured value,
α	angle of rotation, rad
β	angle of inclination of the axis of rotation to the horizontal, deg.
δ	fluid depth in partially filled coil tube, m
$\Delta h$	water level difference in a coil during rotation, m
ε	cross sectional area of flow factor, -
$\eta_H$	hydraulic efficiency, -
ν	kinematic viscosity of the fluid, m <sup>2</sup> /s
$ ho_w$	density of water, kg m <sup><math>-3</math></sup>
τ	wall shear stress, N/m <sup>2</sup>
$\psi$	angle of inclination of the cylindrical helix to pump baseline, deg.

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