

Review

# Mechanical Properties and Deformation Behavior of Bulk Metallic Glasses

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**Abstract:** Metallic glasses demonstrate unique properties, including large elastic limit and high strength, which make them attractive for practical applications. Unlike crystalline alloys, metallic glasses, in general, do not exhibit a strain hardening effect, while plastic deformation at room temperature is localized in narrow shear bands. Room-temperature mechanical properties and deformation behavior of bulk metallic glassy samples and the crystal-glassy composites are reviewed in the present paper.

Keywords: metallic glasses; deformation; strength; plasticity

### 1. Introduction

Owing to their promising physical and functional properties, bulk metallic glassy alloys, which are also called bulk metallic glasses (BMGs), attract significant attention from materials scientists [1–3]. Zr-based glassy alloys have a high glass forming ability (GFA) compared to other TM-based alloys (TM-transition metals), except for some NM-based (NM-noble metals) alloys, which GFA is significantly enhanced by fluxing. The Zr-LTM (LTM-late transition metal) and Zr-LTM-NM glassy

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alloys may demonstrate a large critical diameter up to 30 mm and low critical cooling rates below 10 K/s [1–3], while Zr-LTM-Ti-Be alloys exhibit even lower critical cooling rates down to 1 K/s. The purpose of the present work is to review mechanical properties and deformation behavior of BMGs at room and cryogenic temperature.

## 2. Mechanical Properties and Deformation of Bulk Metallic Glasses and Composites at Room Temperature

Owing to the absence of crystalline lattice and dislocations, a unique deformation mechanism is realized in bulk metallic glassy alloys [4,5], which thus, exhibit high strength (~2 GPa for Cu-, Ti-, Zr-based, ~3 GPa for Ni-based, ~4 GPa for Fe-based, ~5 GPa [6] and even ~6 GPa for Co-based alloys [7]), high hardness, good wear resistance and large elastic deformation. (Fe, Co)-Cr-Mo-C-B-Tm glassy alloys prepared by Amiya and Inoue [8] in a cylindrical form with a diameter of 18 mm demonstrate an excellent GFA and high strength exceeding 4 GPa.  $Co_{43}Fe_{20}Ta_{5.5}B_{31.5}$  glassy alloy exhibited ultrahigh fracture strength exceeding 5 GPa, high Young's modulus of 270 GPa, high specific strength and high specific Young's modulus [6]. The strength and specific strength values exceed those reported for any bulk glassy alloys. The highest strength among Cu-based alloys was found for a Cu-Zr-Ti-Co bulk metallic glass [9].

Mechanical properties of bulk metallic glasses (BMGs) in terms of Yield ( $\sigma_y$ ) and Fracture Strength ( $\sigma_f$ ), Vickers Hardness as well as the Young's Modulus (E) are given in Tables 1 and 2 (together with the references) while possible correlations are shown in Figure 1. One can admit an excellent correlation between  $\sigma_y$  and HV ( $R^2 = 0.98$ ) as expected for isotropic solids like BMGs. Understandably,  $\sigma_f$  shows a less good correlation which is connected with different amount of plastic deformation in the case of each alloy. There is also a similar trend with Young's Modulus but the results are more scattered.

]	Element		С	ontent, at	%	$\sigma_{\rm y}$	$\sigma_{\rm f}$	Ε	HV	Reference
1	2	3	1	2	3					
Al	La	Cu	35	55	1		880	43	256	[10]
Al	La	Cu	30	50	20		750	32	240	[10]
Al	La	Cu	25	55	20		600	32	213	[10]
Al	La	Cu	25	50	25		535	29	208	[10]
Cu	Hf	Al	50	42.5	7.5		2370	128	673	[11]
Cu	Hf	Al	52.5	40	7.5		2345	125	661	[11]
Cu	Hf	Al	52.5	40	7.5		2344	125	661	[12]
Cu	Hf	Al	50	45	5		2262	121	627	[11]
Cu	Hf	Al	50	45	5		2260	121	627	[12]
Cu	Hf	Ti	60	25	15	2010	2160	124		[13]
Cu	Hf	Ti	60	25	15	2024	2088	124		[12]
Cu	Hf	Ti	60	25	15	1920	2130	120		[14]
Cu	Zr	Ag	50	45	5		1940	112	599	[15]
Cu	Zr	Ag	50	45	5		1940	112	599	[15]
Cu	Zr	Ag	45	47.5	7.5		1820	108	556	[15]

Table 1. Composition and mechanical properties of ternary bulk glassy alloys.

]	Element		Content, at %			σ	$\sigma_{\rm f}$	Ε	HV	Reference
1	2	3	1	2	3	•				
Cu	Zr	Ag	45	45	10		1810	108	542	[15]
Cu	Zr	Ag	42.5	47.5	10		1780	106	534	[15]
Cu	Zr	Ag	45	50	5		1885	111	585	[15]
Cu	Zr	Al	47.5	47.5	5	1547	2265	87		[16]
Cu	Zr	Al	55	40	5		2210	115	581	[17]
Cu	Zr	Al	52.5	42.5	5		2115	111	573	[17]
Cu	Zr	Al	50	45	5		1885	102	546	[17]
Cu	Zr	Al	46	46	8	1894	2250		580	[18]
Cu	Zr	Al	55	40	5		2210	115	581	[19,20]
Cu	Zr	Al	52.5	42.5	5		2115	111	573	[19,20]
Cu	Zr	Al	50	45	5		1885	102	546	[19,20]
Cu	Zr	Al	48	48	4	1199	1882	103		[21]
Cu	Zr	Al	47	47	6	1733	2250		580	[22]
Cu	Zr	Ga	52.5	42.5	5		1940	105	552	[23]
Cu	Zr	Ga	55	40	5		2025	109	565	[23]
Cu	Zr	Ga	52.5	40	7.5		2130	111	581	[23]
Cu	Zr	Ga	57.5	40	2.5		1910	105	547	[23]
Cu	Zr	Ti	60	30	10	1785	2150	114		[24]
Gd	Al	Ni	60	30	10		1330	67		[25]
Gd	Al	Ni	55	25	20		1300	65		[26]
Gd	Al	Ni	65	25	10		1300	63		[26]
Gd	Co	Al	60	30	10		1186	60		[25]
Gd	Co	Al	60	25	15		1250	63		[25]
Gd	Co	Al	60	20	20		1240	63		[25]
Gd	Ni	Al	60	25	15		1280	64		[25]
Gd	Ni	Al	60	20	20		1240	63		[25]
Gd	Ni	Al	50	25	25		1320	66		[26]
La	Al	Ni	45	45	10	1080	795	52	330	[27]
La	Al	Ni	45	35	20	995	720	46	305	[27]
La	Al	Ni	50	35	15	950	715	41	290	[27]
La	Al	Ni	50	30	20	930	715	41	285	[27]
La	Al	NI	55	25	20	735	515	34	225	[22]
Mg	Cu	Y	80	10	10	630	820		220	[28]
Mg	Cu	Y	75	15	10		743	50		[28]
Mg	Ni	Gd	75	15	10		929			[29]
Mg	Ni	Gd	70	20	10		880			[29]
Mg	Ni	Gd	70	15	15		965			[29]
Mg	Ni	Gd	65	25	10		884			[29]
Mg	Ni	Gd	65	20	15		909			[29]
Mg	Ni	Gd	60	25	15		869			[29]
Mg	Ni	Y	82.5	12.5	5		610	44	212	[30]
Mg	Ni	Y	80	15	5		830	46	224	[30]
Mg	Ni	Y	85	10	5		640	40	193	[30]

Table 1. Cont.

]	Element		Content, at %			σ <sub>y</sub>	$\sigma_{\rm f}$	Е	HV	Reference
1	2	3	1	2	3					
Zr	Al	Ni	70	10	20	1411	1335	61	432	[31]
Zr	Al	Ni	65	10	25	1581	1520	64.5	484	[31]
Zr	Al	Ni	65	15	20	1614	1640	70.5	494	[31]
Zr	Al	Ni	60	15	25	1640	1715	72.6	502	[31]
Zr	Al	Ni	60	20	20	1795	1720	78.2	549	[31]
Zr	Co	Al	55	30	15		1790	98	543	[32]
Zr	Co	Al	55	25	20		1750	96	530	[32]
Zr	Co	Al	55	25	20		1900	114		[33]
Zr	Cu	Al	50	40	10		1821	89		[34]
Zr	Cu	Al	50	40	10		1860	88	496	[35]
Zr	Cu	Al	52.5	37.5	10		1840	86	485	[32,36]
Zr	Cu	Al	50	37.5	12.5		1960	93	511	[32,36]
Zr	Cu	Al	50	42.5	7.5		1820	86	475	[32,36]
Zr	Cu	Al	55	35	10		1810	83	470	[36]
Zr	Cu	Al	60	30	10		1720	80	446	[36]
Zr	Cu	Al	47.5	42.5	10		1920	90	508	[36]
Zr	Ni	Al	60	25	15		1760	88	495	[31]
Zr	Ni	Al	55	25	20		1780	89	502	[31,37]
Zr	Ni	Al	55	30	15		1820	99	514	[31,37]
Zr	Ni	Al	60	20	20	1793	1720	78.2	549	[38]
Zr	Ni	Al	70	20	10	1411	1335	61	432	[31,39]
Zr	Ni	Al	65	25	10	1520	1581	64.5	484	[31,39]
Zr	Ni	Al	65	20	15	1614	1640		494	[31,39]
Zr	Ni	Al	60	25	15	1640	1715		502	[31,39]
Zr	Ni	Ti	40	37	23		1630		524	[40]

 Table 1. Cont.

**Table 2.** Composition and mechanical properties of quaternary bulk glassy alloys.

Element			Content, at %				$\sigma_{y}$	$\sigma_{\rm f}$	Ε	HV	Reference	
1	2	3	4	1	2	3	4					
Ce	Al	Cu	Co	68	10	20	2		1180	31.34		[41]
Ce	Al	Cu	Fe	68	10	20	2		1232	32.7		[41]
Ce	Al	Cu	Nb	68	10	20	2		1165	30.95		[41]
Ce	Al	Cu	Ni	68	10	20	2		1198	31.93		[41]
Co	Fe	Та	В	38	25	5.5	31.5		5185	268		[42]
Cu	Hf	Ti	Та	56.4	23.5	14.1	6	2125	2100	104		[43]
Cu	Zr	Ag	Al	45	45	7	3		1836	110	540	[44]
Cu	Zr	Ag	Al	45	45	5	5		1890	112	556	[44]
Cu	Zr	Ag	Al	45	45	3	7		1912	112	561	[44]
Cu	Zr	Hf	Ag	45	25	20	10		2000	122	579	[41]
Cu	Zr	Ti	Be	55.5	27.75	9.25	7.5		2450	146	710	[45]
Cu	Zr	Ti	Y	58.8	29.4	9.8	2	1780	2050	115		[46]
Fe	Si	В	Nb	72	9.6	14.4	4		4200	200		[47]
La	Al	Cu	Ag	62.5	12.5	20	5		640	36	201	[48]

Element				Conter	nt, at %		σ	$\sigma_{\rm f}$	Е	HV	Reference	
1	2	3	4	1	2	3	4	•				
La	Al	Cu	Ag	55	15	20	10		758	42	208	[48]
Mg	Cu	Ni	Gd	65	5	20	10		874	54		[49]
Mg	Y	Zn	Cu	65	10	5	20		860	74		[50]
Ni	Nb	Ti	Zr	60	15	10	15		2770	156		[51]
Ni	Si	В	Nb	72	7.68	16.32	4		2510	77	870	[52]
Ni	Si	В	Та	72	7.68	16.32	4		2730	75	920	[52]
Ni	Та	Ti	Zr	60	15	15	10		3180	67		[53]
Pd	Cu	Ni	Р	40	30	10	20		1640		515	[54]
Pd	Cu	Si	Р	79	6	10	5	1475	1575	82		[55]
Pd	Pt	Cu	Р	35	15	30	20	1410			470	[56]
Pt	Cu	Ni	Р	57.5	14.7	5.3	22.5	1400	1470			[57]
Ti	Ni	Cu	Sn	50	20	25	5		2050	102	650	[58]
Ti	Ni	Cu	Sn	50	20	23	7		2200	105	670	[58]
Ti	Ni	Cu	Sn	50	22	25	3		2050	98	640	[58]
Zr	Al	Co	Cu	55	20	20	5	2000	1960	92		[59]
Zr	Al	Ni	Pd	65	7.5	10	17.5	1340	1510			[60]
Zr	Cu	Ni	Al	52	32	4	12		1780	88	501	[61]
Zr	Cu	Ni	Al	52	30	6	12		1820	93	506	[61]
Zr	Cu	Ni	Al	50	26	12	12		1878	88	498	[61]
Zr	Cu	Ni	Al	50	34	4	12		1905	91	517	[61]
Zr	Cu	Ni	Al	48	32	8	12		1894	94	513	[61]
Zr	Cu	Ni	Al	50	32	6	12		1875	92	521	[61]
Zr	Cu	Ni	Al	52	28	8	12		1798	94	512	[61]
Zr	Cu	Ni	Al	50	30	8	12		1820	92	526	[61]
Zr	Cu	Ni	Al	46	34	8	12		1777	111	562	[61]
Zr	Cu	Ni	Al	48	28	12	12		1906	102	530	[61]
Zr	Cu	Ni	Al	48	34	6	12		1899	94	529	[61]
Zr	Cu	Ni	Al	50	28	10	12		1993	92	517	[61]
Zr	Cu	Ni	Al	48	30	10	12		1378	94	520	[61]
Zr	Cu	Ni	Al	48	30	10	12		1980	92	528	[61]
Zr	Cu	Ni	Al	52	26	10	12		1960	89	509	[61]
Zr	Cu	Ni	Al	46	30	12	12		1399	106	552	[61]
Zr	Cu	Fe	Al	60	25	5	10	1643		92		[62]
Zr	Cu	Fe	Al	60	20	10	10	1708		104		[62]
Zr	Fe	Al	Cu	60	10	7.5	22.5	1718		100		[62]



Unless phase separated, bulk metallic glasses are isotropic solids, and thus, their mechanical properties are much less dependent of the processing history than those of crystalline alloys which microstructure can be modified in a wide range resulting in different mechanical properties. Structural relaxation on annealing modifies their plasticity/toughness but very slightly Yield Strength or hardness values.

Although, a vast majority of bulk metallic glasses break shortly after yielding at ambient temperature, the deformation behavior of bulk metallic glassy alloys is under intense investigations at present owing to their high strength. At relatively low homologous temperatures the inhomogeneous plastic flow of glassy alloys occurs by propagation of shear bands [63] which are 10–20 nm thick and make steps on the surface up to several micrometers in height [64]. A strongly localized shear deformation at room temperature [65,66] limits practical application of such materials since a shear event may trigger a crack nucleation and rapid fracture. However, as illustrated in Figure 2, considerable apparent ductility was observed in several specific bulk metallic glassy alloys [67,68].

 $Pd_{40}Ni_{40}Si_xP_{20-x}$  [69,70], Zr-Cu-Fe-Al [71] and Zr<sub>65</sub>Al<sub>7.5</sub>Ni<sub>10</sub>Pd<sub>17.5</sub> [72] alloys exhibit significant room-temperature ductility of several percent and high fracture strength in the glassy state. Especially, ductility of the  $Pd_{40}Ni_{40}Si_4P_{16}$  BMG is drastically improved by application of  $B_2O_3$  fluxing technique [70]. Similar results were obtained for  $Pt_{57.5}Cu_{14.7}Ni_{5.3}P_{22.5}$  [73] alloy though no structural characterization was provided. Large plasticity was explained owing to high elastic moduli ratio G/B or Poisson's ratio [74]. Further investigations on this subject were taken recently and embrittlement of Mg- and La-based alloys is caused by the inclusions of oxides (oxygen is insoluble in solid Mg and La, while it dissolves very well in Zr and Ti) acting as stress concentrators [75]. Zr-Cu-Fe-Al bulk glassy alloy exhibits phase separation on heating prior to crystallization [76], which also may be responsible for its higher room-temperature ductility compared to Zr-Cu-Al alloys.

A stress-strain curve in the plastic deformation range, which is typical for ductile bulk metallic glassy alloys, [77,78] is shown in Figure 2. One can see general absence of strain hardening. The average stress drop magnitude increases with strain, reaching a sort of a critical value at break. The amount of stress drop at each serration event increases with the degree of deformation.

**Figure 2.** A part of the true stress-strain diagram of  $Zr_{60}Cu_{25}Fe_5Al_{10}$  glassy alloy tested at room temperature.



It was shown that fracture of bulk metallic glasses is very sensitive to macroscopic defects (likely pores and surface imperfections) which explains large scattering in the values of the percentage deformation changing from sample to sample [79]. One should also remember that the engineering plasticity of bulk metallic glasses depends upon their geometry and aspect ratio [80,81]. Different specimen geometries influence the shear-band formation and deformation processes, resulting in variations of mechanical properties.

*In situ* monitoring of the deformation behavior of bulk metallic glasses was performed using an infrared camera. Although temporal resolution of the existing cameras is not high enough, some weak rise of temperature was observed as a result of shear band operation [82].

Fatigue properties of BMG alloys have been also studied and fatigue-endurance limits of some Zr-based alloys are comparable with those of high-strength crystalline alloys [83].

As a localized shear deformation is a dominant plastic-deformation mode at room temperature (Figure 3), tensile ductility is not observed in bulk metallic glassy samples except in a few special cases at high strain rates [84]. Nevertheless, some BMGs demonstrate the formation of multiple shear bands (Figure 3a) while others deform through a major single shear band leading to crack formation (Figure 3b).

**Figure 3.** Two regimes of deformation observed in metallic glasses: (**a**) relatively homogeneous deformation with many shear bands in the  $Zr_{62.5}Cu_{22.5}Fe_5Al_{10}$  BMG and (**b**) stick-slip behavior in a single major shear band in the  $Zr_{67.5}Cu_{17.5}Fe_5Al_{10}$  BMG (SEM). Reprinted from [71] with permission of Elsevier.



In order to enhance the ductility of metallic glasses, the formation of heterogeneous microstructures in a composite-like manner has been found recently essential and employed in a variety of procedures combining a glassy matrix with second phase crystalline particles [85,86]; the deformation-induced nanocrystallization [87,88] is one of the possible ways illustrating this approach. One of the ways is to impede the propagation of shear bands through the sample by interaction with the phases (or microscopic pores) embedded into a glassy matrix. This enables multiplication, branching and termination of the shear bands similar to the composites where cracks are either blocked by reinforcements or blunted in ductile phases or matrix.

The blockage of the crack-tip area by nanoparticles stabilizing plastic deformation was observed in  $Cu_{50}Zr_{50}$  as well as in the  $Zr_{65}Al_{7.5}Ni_{10}Pd_{17.5}$  glassy alloys [89]. Moreover, the compressive room-temperature stress-strain curves of the  $Zr_{65}Al_{7.5}Ni_{10}Cu_{17.5}$  and  $Zr_{65}Al_{7.5}Ni_{10}Pd_{17.5}$  bulk metallic glassy alloys revealed significantly different plasticity of these alloys, though both of the alloys possess a glassy structure [90].

An alloy, which crystallizes by the nucleation and growth mechanism, is supposed to be less prone to nanocrystalsization during deformation while the alloy having pre-existing nuclei is predisposed to show such a behavior [91]. For example, the precipitation of the nanoscale  $\alpha$ -Al particles at room temperature was observed in marginal Al-TM-RE glass-formers (like Al<sub>90</sub>Fe<sub>5</sub>Gd<sub>5</sub> or Al<sub>90</sub>Fe<sub>5</sub>Ce<sub>5</sub>) within the shear bands on deformation [92,93], while a very limited fraction or complete absence of nanocrystals (depending on RE metal) was observed after significant plastic deformation in the Al<sub>85</sub>RE<sub>8</sub>Ni<sub>5</sub>Co<sub>2</sub> amorphous and glassy samples [94] (~30% reduction in thickness) which undergo a clear nucleation and growth behavior.

The nanocrystallization within the shear bands was also observed in the  $Cu_{47.5}Zr_{47.5}Al_5$  glassy alloy, which was reported to be strain-hardenable [87,95]. The effect of strain rate on the flow stress [96] and ductility was studied [97,98] and the results showed that the higher strain rate used facilitates plasticity.

In general, bulk metallic glasses exhibit so-called vein patterns on the fracture surface. On the other hand, fracture surface of a Zr-based bulk glassy alloy [99], exhibited alternating areas of vein and flat

fracture surfaces (Figure 4). A fracture surface topography showing initial cleavage-like behavior starting from the lateral surface of the rod sample has long been known [100]. However, in the present case the several cleavage-like and vein patterns alternate in the body of the sample in the direction normal to shear deformation. Thus, cleavage-like areas may be considered as the areas of significant shear deformation. The vein fracture patterns, which are specific to metallic glasses, may signify local melting of the material at the surface upon fracture [101,102], and thus, existence of considerable temperature gradients in the deforming sample. Novel fracture behavior of Pd-based bulk metallic glasses with wavy features was also reported recently by Chen [103].

Figure 4. Fracture surface of the Zr<sub>60</sub>Cu<sub>16</sub>Ni<sub>14</sub>Al<sub>10</sub> bulk glassy alloy tested at room temperature (SEM).



The mechanisms of ductilization of bulk metallic glassy alloys by crystalline and quasicrystalline inclusions have been studied before. Most commonly, they are associated with blocking and branching of the propagating localized shear bands by the structural inhomogeneities like inclusions [86]. Bulk glassy-crystal composites with considerable tensile ductility were produced in Zr-based system alloys with large Be content [104,105]. An interesting result was obtained recently for a Ni-Cu-Ti-Zr alloy [106]. The Ni<sub>40</sub>Cu<sub>10</sub>Ti<sub>33</sub>Zr<sub>17</sub> bulk composite alloy produced by copper mold casting shows a superior combination of strength and ductility not only due to a composite effect resulting from a multiphase structure, but also due to a strain-induced martensitic transformation [106]. The high ultimate strength of 2000 MPa and yield stress of 650 MPa is paired with enhanced ductility estimated at 15% compressive strain. The microstructural investigations revealed the presence of the austenitic cP2 cubic (Ni,Cu)(Ti,Zr) phase (both phases are NiTi-type) which contains other alloying elements forming solid solution is formed in the course of plastic deformation. These crystals further act as strong barriers for shear band propagation, promoting increase in stress and subsequent shear band formation and branching, thus extending the ductility and preventing a premature brittle fracture.

Unlike common observations of serrated plastic flow associated with the large scale shear bands in monolithic metallic glasses, propagation of the shear bands in the studied alloy does not lead to the stress drops, although the shear bands were readily seen on the lateral surface of the samples tested to fracture in addition to multiple slip bands linked to each cP2 crystalline grain. By analogy with deformation-induced martensitic transformations, one can suppose that the formation of the mP4 martensitic phase initiates within the cP2 crystalline phase [107,108]. The most of plastic deformation attained at stresses below 1400 MPa is transformation-aided. At the second stage of plastic deformation, there is an increasing strain hardening and intensive formation of shear bands and slip markings near the fracture surface.

#### 3. Deformation of Bulk Metallic Glasses at Cryogenic Temperature

Low-temperature measurements have also been performed [109–111] and demonstrated the increase in the maximum compressive stress with decreasing testing temperature. The stress-strain diagrams for the typical glassy  $Zr_{64.13}Ni_{10.12}Cu_{15.75}Al_{10}$  alloy tested at room temperature and close to boiling liquid nitrogen (LN<sub>2</sub>) temperature are shown in Figure 5. The sample tested in liquid nitrogen exhibits notably higher compression strength compared to the sample tested at room temperature.

**Figure 5.** Stress-strain diagrams of  $Zr_{64,13}Ni_{10,12}Cu_{15,75}Al_{10}$  glassy alloy tested at room temperature and close to boiling  $LN_2$  temperature (the sample which showed the maximum strength). The data is taken from Reference 111.



The samples tested in  $LN_2$  to different degree of deformation still exhibited formation of the localized shear deformation bands (Figure 6). Although, the stress-strain curves of the BMGs studied at room temperature (Figures 2 and 5) reveal regular stress drops the serrations do not show up (at least up to the sensitivity of the experimental setup) during testing at cryogenic temperatures even in a close-up view of the stress-strain curve of Figure 5. The samples tested in  $LN_2$  show a moderate shear plasticity while the mirror looking cleavage-like areas on the fracture surface are mixed up with typical vein patterns in nearly equal proportions like it is shown in Figure 4 for the sample tested at room temperature.

**Figure 6.** SEM image in secondary electrons of the lateral surface of the  $Zr_{64.13}Ni_{10.12}Cu_{15.75}Al_{10}$  glassy alloy sample after mechanical test at boiling  $LN_2$  temperature.



The mechanical behavior and the kinetics of shear deformation in bulk metallic glassy alloys were also investigated at room and liquid nitrogen temperature using the acoustic emission technique. It was demonstrated that the intensive acoustic emission reflecting the activity of strongly localized shear bands at room temperature vanishes at the transition from serrated to non-serrated plastic flow at LN<sub>2</sub> temperature [112]. The disappearance of acoustic emission signals clearly suggests that the shear band propagation velocity significantly decreases at low temperature likely owing to the increased viscosity of the deforming regions.

#### 4. Strain-Rate Sensitivity

While the strain-rate sensitivity (SRS) of polycrystalline alloys is generally positive (except for special cases) [113], for bulk metallic glasses (BMGs) different results were reported. A negative SRS was reported for  $Zr_{57}Ti_5Cu_{20}Ni_8Al_{10}$  alloy at strain rates ranging from  $10^{-4}$  to  $3 \times 10^3$  s<sup>-1</sup> [114]. In addition, a negative value was obtained [115] for a similar alloy composition, e.g.,  $Zr_{52.5}Ti_5Cu_{17.9}Ni_{14.6}Al_{10}$  (Vit105) and for Vit105-based composite, at strain rates of  $3.3 \times 10^{-4}$  and  $3.7 \times 10^{-3}$  s<sup>-1</sup>. In Zr-based metallic-glass coatings with micrometer-scale thickness for the shallow indentation, in which the substrate effect can be neglected, the increase of the penetration rate also leads to the decrease of the hardness [116]. At the same time, nanoindentation studies indicate positive SRS parameter values [117].

In order to illustrate strain-rate sensitivity of BMG alloys, the  $Zr_{65}Cu_{20}Fe_5Al_{10}$  rod (as one of relatively ductile BMG alloys) glassy samples were tested at room temperature and the strain rate was suddenly changed upon deformation.

Figure 7 shows the true stress-plastic strain curves calculated from engineering stress values. After an initial region of apparent strain hardening, the magnitude of the stress drops increases slightly. The value of strain rate sensitivity (SRS) parameter (m) has been calculated from the true stress-strain curves according to the relationship:

$$m = \left(\frac{\partial \ln \sigma}{\partial \ln \varepsilon}\right)_{\varepsilon, \mathrm{T}}$$
(1)

where  $\sigma$  is the flow stress and  $\dot{\epsilon}$  is the strain rate. It is important to note that the strain rate changes have been performed in the regions where the stress-strain curves are generally linear which may correspond to the formation of multiple shear bands and macroscopically homogeneous deformation of the sample.

**Figure 7.** A part of the true stress-plastic strain curve (elastic strain is subtracted) obtained at the strain rates of  $5 \times 10^{-3}$  and  $5 \times 10^{-2}$  s<sup>-1</sup>. The data is taken from Reference [118] with permission of Elsevier.



The mean and maximum flow stresses and the confidence interval (probability p = 0.95) have been calculated in 0.25% plastic deformation interval using true stress values at both sides of each point of the strain rate change (Figure 7). An average maximum stress before each stress satisfied normality tests and produced similar results. Not only the flow stress but also the average maximum stress before each stress drop corresponding to a serration for each test is practically the same at each loading rate in the vicinity of the strain rate change point and the differences are practically within the confidence interval which indicates that the  $Zr_{65}Cu_{20}Fe_5Al_{10}$  BMG is strain rate insensitive (m~0) in the range of strain rates from  $5 \times 10^{-6}$  to  $5 \times 10^{-4}$  s<sup>-1</sup>. Between  $5 \times 10^{-4}$  and  $5 \times 10^{-3}$  s<sup>-1</sup> the average flow stress is similar,  $1684.4 \pm 0.5$  MPa and  $1685.2 \pm 1.6$  MPa, respectively, and the difference is within the confidence interval. It indicates that the alloy is not strain rate sensitive at these conditions. However, when the strain rate increases from  $5 \times 10^{-3} \pm 2.0$  MPa, and according to Equation 1 the SRS is negative (m = -0.0026).

From the statistical analysis of the mean and maximum flow stress at each loading rate it has been calculated that the alloy does not exhibit SRS within the confidence interval from  $5 \times 10^{-6}$  to  $5 \times 10^{-3}$  s<sup>-1</sup>. The SRS values derived by using the confidence interval are smaller (m~0) than the value of m reported for Zr<sub>52.5</sub>Ti<sub>5</sub>Cu<sub>17.9</sub>Ni<sub>14.6</sub>Al<sub>10</sub> at room temperature (e.g., about -0.002) by [119] and Cu<sub>50</sub>Zr<sub>50</sub> BMG by [120]. However, the SRS becomes meaningfully negative when the strain rate increases from  $5 \times 10^{-3}$  s<sup>-1</sup> to  $5 \times 10^{-2}$  s<sup>-1</sup> because the strain rate is so fast that the relaxation time is not enough to build up the stress which explains the decrease in the flow stress.

Despite on the minimum strain rate at which this alloy exhibits a negative SRS value is  $5 \times 10^{-3}/5 \times 10^{-2} \text{ s}^{-1}$  it should be taken into account that it depends on the loading conditions. For example, while  $Zr_{52.5}Ti_5Cu_{17.9}Ni_{14.6}Al_{10}$  one exhibits negative SRS in compression by Reference [121] it is strain rate insensitive under indentation by Reference [122]. The different behavior could be

mainly attributed to the stress state, which is closer to three-axial compression when tested by indentation. Moreover, there is a strong size effect on deformation behavior upon indentation in which the deformed volume is significantly smaller while local plastic deformation is larger than in compression.

#### 5. In situ Room-Temperature Tensile Deformation in TEM

At the same time, metallic glasses are proven to deform homogeneously on nanoscale. Plastic deformation in metallic glasses has been performed using the in-situ TEM technique [123–125]. After initial observation, the beam current was reduced and the sample deformed in a straining holder in tension. This is an important precaution to prevent electron-beam heating of the sample upon deformation. This procedure (straining and subsequent observation) has been repeated several times until some cracks had become apparent. High-resolution in-situ TEM investigation showed that the Zr<sub>65</sub>Ni<sub>10</sub>Cu<sub>5</sub>Al<sub>7.5</sub>Pd<sub>12.5</sub> ribbon sample retains its glassy structure after in-situ deformation and crack propagation (Figure 8). No nanocrystallization was observed even near the crack.

**Figure 8.** High-resolution TEM image obtained upon *in-situ* deformation in a microscope near the crack in the  $Zr_{65}Ni_{10}Cu_5Al_{7.5}Pd_{12.5}$  glassy alloy.



The homogeneous deformation of the thinned  $Zr_{65}Ni_{10}Cu_5Al_{7.5}Pd_{12.5}$  alloy is clearly seen in front of the crack tip upon loading. The deformation rate estimated by crack opening speed was quite low—about tens of nanometers per second or even less. A narrow featureless deformation zone forms ahead of the crack tip, with the brightness different from that in the rest of the sample. However, this zone is much wider than the typical (10–20 nm) width of a shear band. Similar results were obtained for submicron (about 200 nm) thick  $Zr_{52.5}Cu_{17.9}Al_{10}Ni_{14.6}Ti_5$  alloy [126].

The *in situ* experiments demonstrate a fundamental possibility of a homogeneous-like like plastic flow to occur in a localized area of glassy metal thin foil without aid of temperature. Indeed, any deformation induced heating effects can be neglected. According to recent calculations the thermal energy release, and thus a rise in the sample temperature, depends on the deformation rate and shear offset, in addition to the thermal conductivity and the strength of the glassy phase [127–129]. However, the deformation areas observed in the ribbon sample tested *in situ* are much wider compared to the highly localized deformation bands observed in bulk glassy sample. Thus, no significant heating could have taken place as the heat generated upon the deformation is dissipated over a larger volume.

#### 6. Effect of Cyclic and Long-Term Creep Deformation in the Elastic Region

The beginning of nanocrystallization was observed in the  $Zr_{62.5}Fe_5Cu_{22.5}Al_{10}$  bulk metallic glassy cubic samples (edge size equal to 4 mm) under low stress of 586 ± 242 MPa in the elastic region while no such an effect was seen in the rod shaped samples of 2 mm in diameter [130]. This alloy also showed phase separation on heating prior to crystallization [76,131]. Later the  $Zr_{62.5}Fe_5Cu_{22.5}Al_{10}$  BMG samples of 3 mm in diameter and 3 mm in height were cyclically loaded at higher stresses of 800 ± 200 MPa and 1000 ± 200 MPa, which are still significantly lower than the Yield Strength. It was found that kinetically frozen anelastic deformation accumulates at room temperature and causes crystallization of metallic glassy phase. There was an increase in the crystallization enthalpy after cyclic loading for 1000 cycles at 1000 ± 200 MPa from 44 J/g to 48 J/g is likely connected with the formation of the nuclei of non-equilibrium phase which together with the residual glassy phase may represent a state with higher free energy. An anelastic contribution can arise from localized viscous flow of structural defects [132,133]. The volume changes were locally accommodated by cumulative anelastic deformation and led to subsequent partial nanocrystallization of the glass.

Uniaxial compression of a BMG close to Yield Strength induced creep of the samples, accompanied by an increase in the stored heat of relaxation [134]. The sample exhibited an increase in compressive plasticity on the subsequent test [135]. This creep is homogeneous and leads to decreases in density and in elastic moduli. The increase in heat of relaxation and in plasticity in BMGs resulting from static loading are reported to depend on the atomic packing density [136].

#### 7. Conclusive Remarks

In general, bulk metallic glasses (glassy alloys) have high strength about two times higher than that of their crystalline counterparts but poor plasticity. However, the selected bulk metallic glasses and crystal-glassy composites on their base exhibit both high strength and good room-temperature plasticity. Moreover, in the case of the Pd-based alloys their plasticity can be drastically improved by fluxing with B<sub>2</sub>O<sub>3</sub>. On deformation, bulk metallic glasses exhibit periodic drops of stress indicating so-called serrated flow, while no such a behavior is seen on testing at cryogenic temperature and in case of various composite samples. Disappearance of the stress drops may be connected with a more homogeneous deformation. On the other hand, submicron-size Zr<sub>65</sub>Ni<sub>10</sub>Cu<sub>5</sub>Al<sub>7.5</sub>Pd<sub>12.5</sub> glassy samples demonstrate homogeneous deformation at room temperature. Some metallic glasses which have pre-existing nuclei demonstrate deformation-induced nanocrystallization within heavily deformed areas while alloys which exhibit nucleation and growth behavior on heating are less predisposed to demonstrate such a behavior. The nanocrystallization was also observed upon cyclic loading. Low-temperature measurements demonstrated the increase in the maximum compressive stress with decreasing testing temperature and absence of stress drops. BMGs showed negative SRS in compression. In addition, metallic glasses were found to deform homogeneously on nanoscale which makes them attractive for applications in micro-electro-mechanical systems.

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