



Review

Structural-Parametric Model of Electromagnetoelastic Actuator for Nanomechanics

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Abstract: The generalized parametric structural schematic diagram, the generalized structural-parametric model, and the generalized matrix transfer function of an electromagnetoelastic actuator with output parameters displacements are determined by solving the wave equation with the Laplace transform, using the equation of the electromagnetolasticity in the general form, the boundary conditions on the loaded working surfaces of the actuator, and the strains along the coordinate axes. The parametric structural schematic diagram and the transfer functions of the electromagnetoelastic actuator are obtained for the calculation of the control systems for the nanomechanics. The structural-parametric model of the piezoactuator for the transverse, longitudinal, and shift piezoelectric effects are constructed. The dynamic and static characteristics of the piezoactuator with output parameter displacement are obtained.

Keywords: structural-parametric model; electromagnetoelastic actuator; piezoactuator; nanomechanics; deformation; parametric structural schematic diagram; displacement; transfer function

1. Introduction

Let us consider the role of an electromagnetoelastic actuator on the piezoeffect, the piezomagnetic effect, and the electrostriction or the magnetostriction effect, which are used for precise alignment in nanomechanics and adaptive optics [1–8]. A piezoactuator uses the inverse piezoeffect and plays a role in the actuation or management of mechanisms, as well as converts electrical signals into the displacement and the force. A piezoactuator is applied for research on the nanomechanics of drives in nanotechnology, microelectronics, adaptive optics, biology as well as in scanning tunnelling microscopes, scanning force microscopes, and atomic force microscopes [7–28].

In the present paper, the generalized structural-parametric model and the generalized parametric structural schematic diagram of an electromagnetoelastic actuator are constructed by solving the wave equation with the Laplace transform for the equation of the electromagnetolasticity in the general form, the boundary conditions on the loaded working surfaces of the actuator, and the strains along the coordinate axes. The transfer functions and the parametric structural schematic diagrams of the piezoactuator are obtained from the generalized structural-parametric model.

In References [6,7], the solution of the wave equation of the piezoactuator was determined. In References [6,14,15,28], the structural-parametric models and the schematic diagrams for the simplest piezoactuators were obtained, and these were transformed into the structural-parametric model of an electromagnetoelastic actuator with output displacements.

The structural-parametric model of the piezoactuator was determined in contrast to electrical equivalent circuit types Cady and Mason for the calculation of the piezoelectric transmitter and receiver, the vibration actuator, and the vibration motor [8–13]. Reference [12] presents the classic analytical two-port lumped-element model (LEM) types Cady and Mason of the piezoelectric composite circular plate with the output pressure. Reference [13] considers the development of various lumped-element models as practical tools to design and manufacture actuators with the output velocity.

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In References [8,14,15], the transfer functions of the piezoactuator were used to overcome the problem of the absolute stability condition of the strain control system for an electromagnetoelastic actuator.

The elastic compliances and the mechanical and adjusting characteristics of a piezoactuator were explored in Reference [17–19] in order to calculate its transfer functions and create the structural-parametric model. The structural-parametric model of a multilayer and compound piezoactuator was determined in References [17–22] with output displacement. In this paper, we solve the problem of building the generalized structural parametric model and the generalized parametric structural schematic diagram of an electromagnetoelastic actuator for the equation of electromagnetoelasticity in the general form. The difference of this work from other reports [23,24,26–28] is that the construction of the structure-parametric model of the electromagnetoelastic actuator is produced immediately in the general form, and not by the method of mathematical induction from individual examples of models of piezoactuators.

2. Structural-Parametric Model and Transfer Functions of Electromagnetoelastic Actuator

Let us consider the general structural-parametric model and the parametric structural schematic diagram of an electromagnetoelastic actuator with output parameter displacement. For the electromagnetoelastic actuator, six stress components are presented; T_1 , T_2 , T_3 , T_4 , T_5 , T_6 , where the components T_1 – T_3 are related to extension-compression stresses, and T_4 – T_6 are related to shear stresses. In the electromagnetoelastic actuator, its deformation corresponds to the stressed state.

For polarized piezoceramics PZT the matrix state equations [11,14] with electric and elastic variables can be given by two equations, where the first equation represents the direct piezoelectric effect, and the second describes the inverse piezoelectric effect:

$$\mathbf{D} = \mathbf{dT} + \mathbf{\varepsilon}^T \mathbf{E} \tag{1}$$

$$\mathbf{S} = \mathbf{s}^E \mathbf{T} + \mathbf{d}^t \mathbf{E} \tag{2}$$

where **D** is the column matrix of electric induction; **S** is the column matrix of relative deformations; **T** is the column matrix of mechanical stresses; **E** is the column matrix of electric field strength; \mathbf{s}^E is the elastic compliance matrix for E = const; \mathbf{e}^T is the matrix of dielectric constants for T = const; and \mathbf{d}^t is the transposed matrix of the piezoelectric modules.

The piezoactuator or piezoplate has the following properties: δ is the thickness, h is the height, b is the width, and $l = \{\delta, h, b \text{ respectively denotes the length of the piezoactuator for the longitudinal, transverse, and shift piezoeffect. The direction of the polarization axis <math>P$, i.e., the direction along which polarization was performed, is usually taken as the direction of axis 3 on kinematic schemes of the piezoactuator, shown in Figure 1.

The equation of electromagnetoelasticity [11,14,24] in the general form is:

$$S_i = d_{mi} \Psi_m(t) + s_{ii}^{\Psi} T_i(x, t)$$
(3)

where $S_i = \partial \xi(x,t)/\partial x$, $\Psi_m(t) = \{E_m(t), D_m(t), H_m(t), \text{ and } S_i \text{ is the relative displacement of the cross-section of the piezoactuator along axis <math>i$, $\Psi_m(t)$ is the control parameter along axis m = 1, 2, 3 for Figure 1, $\xi(x,t)$ is the displacement of the section of the piezoactuator, d_{mi} is the piezomodule, $E_m(t) = U(t)/\delta$ is the electric field strength along axis m, U(t) is the voltage between the electrodes of the actuator, $D_m(t)$ is the electric induction along axis m, $H_m(t)$ is the magnet field strength along axis m, S_{ij}^{Ψ} is the elastic compliance for $\Psi = \text{const}$, T_j is the mechanical stress along axis j, and i, $j = 1, 2, \ldots$, 6. The main size is the length l of the piezoactuator, respectively, denotes the thickness, height, and width for the longitudinal, transverse, and shift piezoeffect in Figure 1a–c.

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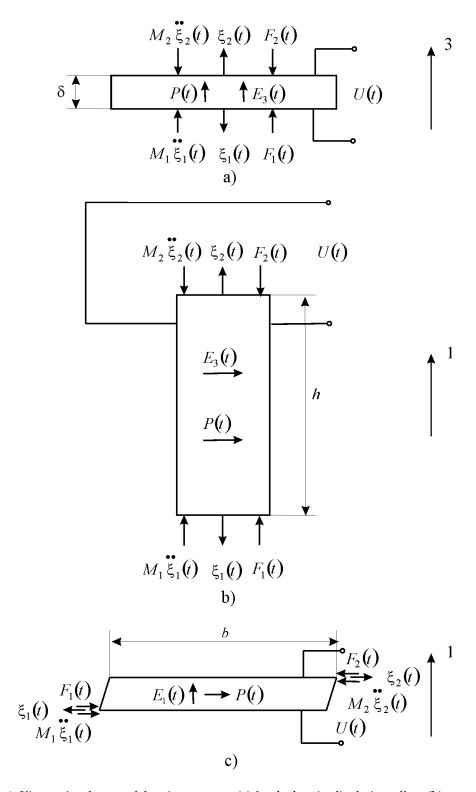


Figure 1. Kinematic schemes of the piezoactuator (a) for the longitudinal piezoeffect; (b) transverse piezoeffect; and (c) shift piezoeffect.

To calculate the electromagnetoelastic actuator, we used the wave equation [6,7,14] for the wave propagation in the long line with damping but without distortions. After Laplace transform is obtained, the linear ordinary second-order differential equation with the parameter p is achieved, where the

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original problem for the partial differential equation of hyperbolic type using the Laplace transform is reduced to the simpler problem for the linear ordinary differential equation.

$$\frac{d^2\Xi(x,p)}{dx^2} - \gamma^2\Xi(x,p) = 0$$
 (4)

with its solution

$$\Xi(x,p) = Ce^{-x\gamma} + Be^{x\gamma} \tag{5}$$

where $\Xi(x,p)$ is the Laplace transform of the displacement of the section of the electromagnetoelastic actuator, $\gamma = p/c^{\Psi} + \alpha$ is the propagation coefficient, c^{Ψ} is the sound speed for $\Psi = \text{const}$, α is the damping coefficient, and Ψ is the control parameter; E stands for the voltage control, D for the current control, and E for the magnet field strength control.

From Equations (3)–(5), the boundary conditions on loaded surfaces, the strains along the axes the system of equations for the generalized structural-parametric model, and the generalized parametric structural schematic diagram are determined in Figure 2 for the electromagnetoelastic actuator with output parameters and the Laplace transform for the displacements of the faces in the form:

$$\Xi_{1}(p) = \left(\frac{1}{M_{1}p^{2}}\right) \left\{ -F_{1}(p) + \left(\frac{1}{\chi_{ij}^{\Psi}}\right) \begin{bmatrix} d_{mi}\Psi_{m}(p) - \\ -\left(\frac{\gamma}{\sinh(l\gamma)}\right) \left[\cosh(l\gamma)\Xi_{1}(p) - \Xi_{2}(p)\right] \end{bmatrix} \right\}
\Xi_{2}(p) = \left(\frac{1}{M_{2}p^{2}}\right) \left\{ -F_{2}(p) + \left(\frac{1}{\chi_{ij}^{\Psi}}\right) \begin{bmatrix} d_{mi}\Psi_{m}(p) - \\ -\left(\frac{\gamma}{\sinh(l\gamma)}\right) \left[\cosh(l\gamma)\Xi_{2}(p) - \Xi_{1}(p)\right] \end{bmatrix} \right\}$$
(6)

where
$$\chi_{ij}^{\Psi} = \frac{s_{ij}^{\Psi}}{S_0} d_{mi} = \begin{cases} d_{33}, d_{31}, d_{15} \\ g_{33}, g_{31}, g_{15} \end{cases}$$
, $\Psi_m = \begin{cases} E_3, E_3, E_1 \\ D_3, D_3, D_1 \end{cases}$, $s_{ij}^{\Psi} = \begin{cases} s_{33}^{E}, s_{11}^{E}, s_{55}^{E} \\ s_{33}^{D}, s_{11}^{D}, s_{55}^{D} \\ s_{33}^{H}, s_{11}^{H}, s_{55}^{H} \end{cases}$

 $l = \{\delta, h, b, c^{\Psi} = \{c^{E}, c^{D}, c^{H}, \gamma^{\Psi} = \{\gamma^{E}, \gamma^{D}, \gamma^{H}, d_{mi} \text{ is the coefficient of the electromagnetolasticity or the piezomodule, the coefficient of magnetostriction, and <math>F_1(p), F_2(p)$ are the Laplace transform of the forces on the faces of the electromagnetoelastic actuator in Figure 1a–c.

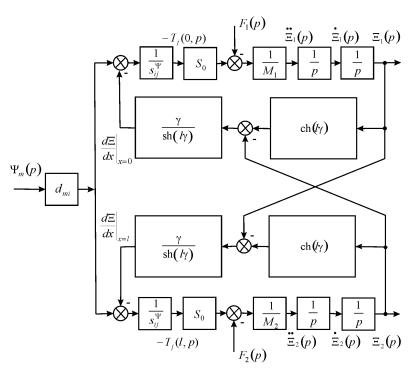


Figure 2. Generalized parametric structural schematic diagram of the electromagnetoelastic actuator.

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The generalized parametric structural schematic diagram (Figure 2) of the electromagnetoelastic actuator corresponds to Equation (6) for the Laplace transform of the displacements of the faces. The generalized transfer functions of the electromagnetoelastic actuator are the ratio of the Laplace transform of the displacement of the face actuator and the Laplace transform of the corresponding control parameter or the force at zero initial conditions. From Equation (6), the generalized matrix equation of the deformation for the electromagnetoelastic actuator has the form:

$$\begin{pmatrix} \Xi_{1}(p) \\ \Xi_{2}(p) \end{pmatrix} = \begin{pmatrix} W_{11}(p) & W_{12}(p) & W_{13}(p) \\ W_{21}(p) & W_{22}(p) & W_{23}(p) \end{pmatrix} \begin{pmatrix} \Psi_{m}(p) \\ F_{1}(p) \\ F_{2}(p) \end{pmatrix}$$
(7)

where

$$\begin{split} W_{11}(p) &= \frac{\Xi_{1}(p)}{\Psi_{m}(p)} = \frac{V_{mi}}{A_{ij}} \left[M_{2} \chi_{ij}^{\Psi} p^{2} + \gamma \text{th} \left(\frac{l \gamma}{2} \right) \right] \\ A_{ij} &= M_{1} M_{2} \left(\chi_{ij}^{\Psi} \right)^{2} p^{4} + \frac{(M_{1} + M_{2}) \chi_{ij}^{\Psi}}{\text{cth}(l \gamma)} p^{3} + \\ &+ \left[\frac{(M_{1} + M_{2}) \chi_{ij}^{\Psi} \alpha}{\text{th}(l \gamma)} + \frac{1}{c^{2}} \right] p^{2} + \frac{2\alpha}{c} p + \alpha^{2} \\ c &= c^{\Psi} \\ W_{21}(p) &= \frac{\Xi_{2}(p)}{\Psi_{m}(p)} = \frac{V_{mi}}{A_{ij}} \left[M_{1} \chi_{ij}^{\Psi} p^{2} + \gamma \text{th} \left(\frac{l \gamma}{2} \right) \right] \\ W_{12}(p) &= \frac{\Xi_{1}(p)}{F_{1}(p)} = -\frac{\chi_{ij}^{\Psi}}{A_{ij}} \left[M_{2} \chi_{ij}^{\Psi} p^{2} + \frac{\gamma}{\text{th}(l \gamma)} \right] \\ W_{13}(p) &= \frac{\Xi_{1}(p)}{F_{2}(p)} = W_{22}(p) = \frac{\xi_{2}(p)}{F_{1}(p)} = \frac{\chi_{ij}^{\Psi} \gamma}{A_{ij} \text{sh}(l \gamma)} \\ W_{23}(p) &= \frac{\Xi_{2}(p)}{F_{2}(p)} = -\frac{\chi_{ij}^{\Psi}}{A_{ij}} \left[M_{1} \chi_{ij}^{\Psi} p^{2} + \frac{\gamma}{\text{th}(l \gamma)} \right]. \end{split}$$

Let us find the displacement of the faces the electroelastic actuator in the stationary regime for $\Psi_m(t) = \Psi_{m0} \cdot 1(t)$, $F_1(t) = F_2(t) = 0$, and inertial load. The static displacement of the faces of the electroelastic actuator $\xi_1(\infty)$ and $\xi_2(\infty)$ can be written in the following forms:

$$\xi_1(\infty) = \lim_{t \to \infty} \xi_1(t) = \frac{\nu_{mi} l \Psi_{m0}(M_2 + m/2)}{M_1 + M_2 + m}$$
(8)

$$\xi_2(\infty) = \lim_{t \to \infty} \xi_2(t) = \frac{\nu_{mi} l \Psi_{m0}(M_1 + m/2)}{M_1 + M_2 + m}$$
(9)

$$\xi_1(\infty) + \xi_2(\infty) = \lim_{t \to \infty} (\xi_1(t) + \xi_2(t)) = \nu_{mi} I \Psi_{m0}$$
 (10)

where m is the mass of the electroelastic actuator, and M_1 , M_2 are the load masses.

The static displacements of the faces of the piezoactuator for the longitudinal piezoeffect (Figure 1a) and $m \ll M_1$ and $m \ll M_2$ are obtained in the following forms:

$$\xi_{1}(\infty) = \lim_{\substack{p \to 0 \\ \alpha \to 0}} \frac{pW_{11}(p)U_{0}}{\delta p} = \frac{d_{33}U_{0}M_{2}}{M_{1} + M_{2}}$$
(11)

$$\xi_{2}(\infty) = \lim_{\substack{p \to 0 \\ \alpha \to 0}} \frac{pW_{21}(p)U_{0}}{\delta p} = \frac{d_{33}U_{0}M_{1}}{M_{1} + M_{2}}$$
(12)

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Let us consider a numerical example for a piezoactuator from PZT $d_{33}=4\cdot 10^{-10}$ m/V, U=400 V, $M_1=1$ kg, and $M_2=4$ kg. We obtain the static displacements of the faces of the piezoactuator for the longitudinal piezoeffect: $\xi_1(\infty)=128$ nm, $\xi_2(\infty)=32$ nm, $\xi_1(\infty)+\xi_2(\infty)=160$ nm.

The static displacements of the faces of the piezoactuator for the transverse piezoeffect (Figure 1b) and $m \ll M_1$ and $m \ll M_2$ are determined in the following forms:

$$\xi_{1}(\infty) = \lim_{\substack{p \to 0 \\ \alpha \to 0}} \frac{pW_{11}(p)U_{0}}{\delta p} = \frac{d_{31}hU_{0}M_{2}}{\delta(M_{1} + M_{2})}$$
(13)

$$\xi_{2}(\infty) = \lim_{\begin{subarray}{c} p \to 0 \\ \alpha \to 0 \end{subarray}} \frac{pW_{21}(p)U_{0}}{\delta p} = \frac{d_{31}hU_{0}M_{1}}{\delta(M_{1} + M_{2})}$$
(14)

For the piezoactuator from PZT under the transverse piezoeffect at $m << M_1$ and $m << M_2$, $d_{31} = 2.5 \cdot 10^{-10}$ m/V, $h/\delta = 20$, U = 360 V, $M_1 = 1$ kg, and $M_2 = 4$ kg. The static displacements of the faces are obtained: $\xi_1(\infty) = 1440$ nm, $\xi_2(\infty) = 360$ nm, $\xi_1(\infty) + \xi_2(\infty) = 1800$ nm.

From Equation (7), the transfer function of the piezoactuator for the transverse piezoelectric effect for one rigidly fixed face of the piezoactuator at $M_1 \to \infty$ is determined in the form:

$$W_{21}(p) = \Xi_2(p)/E_3(p) = d_{31}h/\left[M_2h\chi_{11}^E p^2 + h\gamma \operatorname{cth}(h\gamma)\right]$$
(15)

From Equation (15), the resonance condition at $M_1 \rightarrow \infty$ and $M_2 = 0$ is obtained in the form:

$$ctgkh = 0 (16)$$

where $k = \omega/c^E$ is the frequency coefficient; ω is the circular frequency.

Because

$$k_i h = \pi (2i - 1)/2 \tag{17}$$

where index i = 1, 2, 3, ...

Therefore, the piezoactuator is the quarter-wave vibrator with the resonance frequency:

$$f_1 = c^E/(4h).$$
 (18)

For the piezoactuator from PZT under the transverse piezoeffect at $c^E = 3 \cdot 10^3$ m/s, $h = 5 \cdot 10^{-2}$ m, the resonance frequency $f_1 = 15$ kHz is obtained. The experimental and calculated values for the piezoactuator are in agreement up to an accuracy of 5%.

For the approximation of the hyperbolic cotangent by two terms of the power series in transfer function (Equation (15)), the following expressions of the transfer function of the piezoactuator are obtained for the elastic-inertial load at $M_1 \to \infty$, $m << M_2$ under the transverse piezoeffect, and the control voltage for the resistance R=0 of the voltage source (Figure 3) is given in the form:

$$W(p) = \frac{\Xi_{2}(p)}{U(p)} = \frac{d_{31}h/\delta}{\left(1 + C_{e}/C_{11}^{E}\right)\left(T_{t}^{2}p^{2} + 2T_{t}\xi_{t}p + 1\right)}$$

$$T_{t} = \sqrt{M_{2}/\left(C_{e} + C_{11}^{E}\right)}, \, \xi_{t} = \alpha h^{2}C_{11}^{E}/\left(3c^{E}\sqrt{M\left(C_{e} + C_{11}^{E}\right)}\right)$$
(19)

where U(p) is the Laplace transform of the voltage, T_t is the time constant, and ξ_t is the damping coefficient of the piezoactuator.

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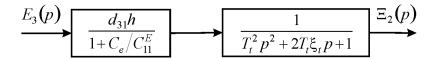


Figure 3. Parametric structural schematic diagram of the voltage-controlled piezoactuator under the transverse piezoeffect for elastic-inertial load.

The expression for the transient response of the voltage-controlled piezoactuator for the elastic-inertial load under the transverse piezoeffect and the resistance R = 0 of the voltage source is determined in the following form:

$$\xi(t) = \xi_m \left[1 - \frac{e^{-\frac{\xi_t t}{T_t}}}{\sqrt{1 - \xi_t^2}} \sin(\omega_t t + \varphi_t) \right]$$

$$\xi_m = \frac{d_{31}(h/\delta)U_m}{1 + C_e/C_{11}^E}, \quad \omega_t = \frac{\sqrt{1 - \xi_t^2}}{T_t}, \quad \varphi_t = \arctan\left(\frac{\sqrt{1 - \xi_t^2}}{\xi_t}\right)$$
(20)

where ξ_m is the steady-state value of displacement of the piezoactuator, and U_m is the amplitude of the voltage.

For the voltage-controlled piezoactuator from the piezoceramics PZT under the transverse piezoelectric effect for the elastic-inertial load $M_1 \rightarrow \infty$, $m << M_2$ and input voltage with amplitude $U_m = 200 \text{ V}$ at $d_{31} = 2.5 \cdot 10^{-10} \text{ m/V}$, $h/\delta = 20$, $M_2 = 9 \text{ kg}$, $C_{11}^E = 2 \cdot 10^7 \text{ N/m}$, $C_e = 0.5 \cdot 10^7 \text{ H/m}$, values $\xi_m = 800 \text{ nm}$, $T_t = 0.6 \cdot 10^{-3} \text{ s}$ are obtained.

3. Results and Discussions

The generalized structural-parametric model, the generalized parametric structural schematic diagram, and the matrix equation of an electromagnetoelastic actuator with output parameters displacements are obtained from the solutions of the wave equation with the Laplace transform and from its deformations along the coordinate axes.

From the generalized matrix equation for the transfer functions of the electromagnetoelastic actuator, the matrix equations of the piezoactuator for the longitudinal, transverse, and shift piezoelectric effects are constructed after algebraic transformations.

The dynamic and static characteristics of the piezoactuator are constructed in order to elucidate the nanomechanics, with regard to its physical parameters and the external load.

The structural-parametric model and the parametric structural schematic diagrams of the voltage-controlled piezoactuator for the longitudinal, transverse, and shift piezoelectric effects are determined from the generalized structural-parametric model of the electromagnetoelastic actuator for the nanomechanics by replacing the generalized parameters with the parameters of the piezoactuator.

4. Conclusions

The generalized structural-parametric model, the generalized parametric structural schematic diagram, and the matrix equation of an electromagnetoelastic actuator with output parameters displacements for the nanomechanics were obtained. The structural-parametric model, the matrix equation for the transfer functions, and the parametric structural schematic diagram of the piezoactuator for the transverse, longitudinal, and shift piezoelectric effects were determined from the generalized structural-parametric model of the electromagnetoelastic actuator.

From the solution of the wave equation with the Laplace transform, the equation of the electromagnetolasticity in the general form, the deformations along the coordinate axes of the generalized structural-parametric model, and the generalized parametric structural schematic diagram of the electromagnetoelastic actuator with output parameters displacements were constructed for the nanomechanics. The deformations of the actuator were described by the matrix equation for the transfer functions of the actuator.

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To calculate control systems in the nanomechanics of nanotechnology, microelectronics, nanobiology, astronomy, and adaptive optics, the transfer functions of the electromagnetoelastic actuator were obtained.

The deformations of the piezoactuator for the nanomechanics were described by the matrix equation and the transfer functions of the piezoactuator. The structural-parametric model, the parametric structural schematic diagrams, and the transfer functions of the piezoactuator make it possible to describe the dynamic and static characteristics of the piezoactuator for the nanomechanics with regard to its physical parameters and the external load.

Conflicts of Interest: The authors declare no conflict of interest.

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