



Article Elastic Stability of Perforated Plates Strengthened with FRP under Uniaxial Compression

Xin Tao, Shuangyin Cao * and Long Zhang

Key Laboratory of Concrete and Prestressed Concrete Structures of Ministry of Education, Southeast University, Nanjing 210096, China; 230129148@seu.edu.cn (X.T.); 101000873@seu.edu.cn (L.Z.)

* Correspondence: 101000873@seu.edu.cn; Tel.: +86-158-518-67851

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Abstract: Openings are frequently introduced in plates for the purpose of inspection, maintenance, service, etc. The presence of openings reduces the buckling and ultimate capacity significantly, and pasting fiber-reinforced polymers (FRP) is an ideal technique for postponing the buckling and increasing the ultimate capacity of the plates. In this paper, the finite element (FE) method has been employed to study the buckling stress of the perforated plates strengthened with FRP under uniaxial compression, and several parameters are considered: material's geometrical and mechanical properties, boundary conditions, plate aspect ratio, hole sizes, and hole position. Then a method of calculating the buckling stress is proposed and modified based on the theory of composite plate and the numerical results. The study shows that, the stiffness modified factor α_D , which considers the orthotropic properties of FRP are a function of the reinforcement index ω and hole size d/b for Boundary conditions (BCs) of 4S and 3S1F. And it is recommended to place the big hole close to the middle area of the plate in *x*-axis. It also shows that for a small hole size, there is little effect of the hole position e_y/b on buckling coefficient K_u regardless of the BCs, and that effect becomes more pronounced as d/b increases, so it is recommended to put the holes near the middle of the plate for 4S and the simple support edge for 3S1F in *y*-axis.

Keywords: perforated plates; elastic buckling; fiber-reinforced polymer (FRP); finite element method

1. Introduction

Thin steel plate elements are the main structural components of decks and platforms on ships, webs and flanges of plate and box girder bridges, platforms of offshore structures, and architecture structure [1]. Openings are frequently introduced in plates for the purpose of inspection, maintenance, service, etc. [2]. In such cases, the presence of openings redistributes the membrane stresses in the plates and reduces their buckling and compression capacity significantly [3–5]. The reduced buckling strength of the perforated plate may be insufficient to meet the requirements of normal structural serviceability and safety. In this situation, it is necessary and important to develop a method to improve the buckling and ultimate strength of perforated steel plates.

Fiber-reinforced polymer (FRP) has large strength to weight ratios, excellent resistance to corrosion and environmental degradation. It is also very flexible to be formed into any desired shape and easy to be handled during construction. Therefore, FRP has been widely used in strengthening steel structures. In recent years, there have been a number of studies on strengthening steel structures, particularly compressed thin-walled structures. Shaat and Fam [6–9] tested hollow structural steel (HSS) specimens strengthened with carbon fiber-reinforced polymers (CFRP) to study the influence on local and global buckling behavior and column strength. The test indicated that the transverse CFRP layers can help restrain outward directed local buckling of two opposite sides for short columns. Harris [10] employed the strengthening technique of externally bonding FRP sheets to selected sides of steel T-sections in order to improve global and local buckling behavior, and the results show that FRP material provides "bracing" against the local buckling. Haedir and Zhao [11] tested ten short cold-formed steel circular hollow section (CHS) columns with externally bonded hoop and longitudinal CFRP, the results show that the reinforcing hoop fibers provide better restraining force to postpone the steel skin buckling when the local slenderness ratio is higher, it also shows that the use of a combination of hoop and longitudinal CFRP in a slender tube can promote the attainment of the yield capacity of the bare tube. Bambach [12–14] tests a CFRP wrapped square hollow sections (SHS) columns with width-to-thickness ratio varying from 12.5 to 120. The experimental results show that the axial capacity of the wrapped SHS is up to twice of that without wrapping the CFRP, due to the restraining on the development of elastic buckling deflections from the wrapping CFRP. The critical stress is up to four times of that of the steel section due to this restraining effect. Moreover, a design method is developed whereby the theoretical elastic critical stress of the composite steel–CFRP sections developed by Pister and Dong [15] is used to determine the buckling load and the ultimate load.

The previous description shows that openings are frequently introduced in plates and the application of FRP is an ideal technique for postponing the local buckling and increasing the ultimate capacity of the steel structure. However, the researches on pasting FRP sheets for improving the degraded buckling of perforated plates are not investigated in the published literatures. Therefore, the objective of the present paper is to investigate the material's geometrical and mechanical properties, boundary conditions, a plate aspect ratio, hole sizes, and hole position on elastic buckling of perforated plates strengthened with FRP under uniaxial compression. And a calculation method of buckling stress is proposed and modified based on the theory of composite plate and the numerical results.

2. Problem Definition

2.1. Basic Information

The problem of elastic buckling of perforated plates strengthened with FRP under uniaxial compression (Figure 1) is investigated in this study. The lengths of the plate in the longitudinal and transverse directions are *a* and *b* (in this paper, equal to 240 mm), respectively. The thicknesses of the steel plate and FRP sheet are t_s and t_f (single layer of FRP sheet's thickness $t_{f0} = 0.167$ mm and FRP Layers' number is n_f), respectively. The hole has a diameter *d*, and e_x (e_y) is the distance along *x*-axis (*y*-axis) from the center of the hole to the left (bottom) edge. The elastic modulus of steel and FRP is E_f and E_s , respectively. The Poisson ratio of FRP and steel is v_f (equal to 0.25) and v_s (equal to 0.3), respectively.



Figure 1. Geometry and Boundary conditions (BCs) of the perforated plates strengthened with fiber-reinforced polymers (FRP).

2.2. Boundary Conditions (BCs) of the Plate

The buckling performances of the plate are significantly affected by the BCs. Rectangle, T and L sections are the common cross section types used in practical, see in Figure 2. It is noted that mechanical properties of compressed members are derived from the behaviors of the individual plate

with L and T sections (Figure 2b,c) can be seen as two or three plates with three edges simple support and one edge free (3S1F), and both the BCs of 4S and 3S1F are considered in this paper, as shown in Figure 1.



Figure 2. The common cross sections and deformations under compression. (a) Rectangle section; (b) L section; (c) T section.

2.3. Load Case

In the present study, only longitudinal uniform uniaxial compression in the *x* direction is considered, as shown in Figure 1. Other load cases such as biaxial compression and edge shear loading will be discussed in the authors' subsequent research papers. The critical stress f_{cr} is defined by the average stress that is equal to the uniformly applied compressive load divided by the full cross-sectional area assuming no opening within the plate.

2.4. Parameters Considered in This Study

The object of the present study is to investigate the buckling of the perforated plates strengthened with FRP under uniaxial compression. The geometric dimensions (such as t_f , t_s), and the elastic modulus of FRP E_f , the BCs of 4S and 3S1F, the effect of plate aspect ratio β (i.e., a/b), hole size d/b, hole position e_x/b and e_y/b have been considered:

- The parameters of geometrical and mechanical properties such as t_f , t_s , E_f , E_s appear in the form $\omega = E_f t_f / E_s t_s = E_f n_f t_{f0} / (E_s t_s)$. In this equation, $E_f = 1.6 \times 10^5$ MPa, 2.4×10^5 MPa and 4.4×10^5 MPa, $E_s = 2.0 \times 10^5$ MPa, n_f varies between 0 and 3 with an increment of 1, $t_{f0} = 0.167$ mm, $t_s = 2$ mm, 3 mm and 4 mm.
- The BCs: 4S and 3S1F.
- Plate aspect ratio β : varies between 1 and 4 with increment of 1.
- Hole size *d/b*: varies between 0 and 0.7 with an increment of 0.1, in some case refines to 0.05.
- Hole position e_x/b and e_y/b . e_x/b varies between 0.05 and 0.5 β with an increment of 0.0625. e_y/b varies between 0.5 and 0.9 with an increment of 0.025 for BCs of 4S and between 0.05 and 0.95 with

an increment of 0.025 for BCs of 3S1F. In all conditions, the distance between hole's edge and plate's edge is no less than 0.05*b*.

3. The Theory of Composite Plate without Holes

The buckling problem of rectangular plate subjected to uniaxial compression, was first studied by Timoshenko [16] in 1961 and modified by Gerstle [17] in 1967. Only first order terms in the governing equations were considered, the critical stress f_{cr} at buckling can be written as

$$f_{cr} = k \frac{\pi^2}{tb^2} D \tag{1}$$

where *k* is a dimensionless buckling coefficient of the plate. *t* and *b* are the thickness and width of the plate, respectively. *D* is the flexural rigidity of the plate, and for the steel plate, it can be written as:

$$D = \frac{E_s t_s^3}{12(1 - v_s^2)}$$
(2)

3.1. The Transformed Flexural Rigidity D_t

In 1959, Pister [15] proposed a solution for the transformed flexural rigidity D_t of the perfectly bonded layered plate where each layer is assumed to be isotropic, a plate element with the associated coordinate system is shown in Figure 3, and the calculation formulas of D_t are in Equations (3)–(6).



Figure 3. A composite plate with the associated system.

$$D_t = C - \frac{B^2}{A} \tag{3}$$

$$A = \sum_{k} \frac{E_k}{1 - v_k^2} t_k \tag{4}$$

$$B = \sum_{k} \frac{E_k}{1 - v_k^2} \frac{t_k^2 - t_{k-1}^2}{2}$$
(5)

$$C = \sum_{k} \frac{E_{k}}{1 - v_{k}^{2}} \frac{t_{k}^{3} - t_{k-1}^{3}}{3}$$
(6)

To a pure steel plate, the flexural rigidity of D_0 can be calculated by Equations (4)–(6), and the rigidity D_0 can be expressed as Equation (7), which is the same with Equation (2), so the method proposed by Pister can also be used to calculate the flexural rigidity of the pure steel plate.

$$D_0 = C_0 - \frac{B_0^2}{A_0} = \frac{E_s t_s^3}{12(1 - v_s^2)}$$
(7)

Considering that the steel and CFRP are perfectly bonded up to the ultimate load [18] and the Young's modulus of the adhesive layers is about 1/10 of those of steel and FRP, the effect of adhesive layers is not taken into account for conservative design in this paper. Therefore, a plate strengthened with FRP can be seen as FRP-steel layered plates. Two reinforcement index α , γ are introduced as given by Equations (8) and (9).

$$\alpha = \frac{E_f}{1 - v_f^2} / \frac{E_s}{1 - v_s^2}$$
(8)

$$\gamma = \frac{t_f}{t_s} \tag{9}$$

Substitution of Equations (8) and (9) into Equations (4)~(6) and then substitution of Equations (4)~(6) into (3) leads to Equation (10), where D_t is the flexural rigidity of plate strengthened with FRP. Therefore, D_t/D_0 can be expressed as Equation (10), and shown in Figure 4. Equation (10) can be approximate simplified as Equation (11) for calculation convenient, and shown in Figure 5.

$$D_t/D_0 = 1 + 3\alpha\gamma \left(1 + \frac{4}{3}\gamma^2 + 2\gamma - \frac{\alpha\gamma(\gamma+1)^2}{1+\alpha\gamma}\right)$$
(10)

$$D_t/D_0 = \begin{cases} \frac{1+3\alpha\gamma, \gamma \leq 0.2}{\left(1+3\alpha\gamma, \frac{(\alpha+2)(\gamma+1)^3}{3}\right), between} \\ \frac{(\alpha+2)(\gamma+1)^3}{3}, \ 0.3 \leq \gamma \leq 1 \end{cases}$$
(11)

From the equations, figures and discussion above, the following can be concluded:

- The method proposed by Pister to calculate the transformed flexural rigidity of the perfectly bonded layered plate is available to calculate the flexural rigidity of the pure steel plate.
- When γ is in a range of 0 to 0.2 and α is in a range of 0.8 to 2.4, D_t/D_0 almost had a linearly related with α and γ . For $\gamma = 0.2$, D_t/D_0 will increase from 1.48 to 2.44 as α enlarges from 0.8 to 2.4.
- D_t/D_0 will have a non-linear increase as γ is enlarged from 0.2 to 0.3.
- When γ is in a range of 0.3 to 1 and α is in a range of 0.8 to 2.4: for a fixed α , D_t/D_0 will had a cubic curvilinear increase as γ is enlarged from 0.3 to 1.0, and for a fixed γ , D_t/D_0 will have a linear change as α enlarges from 0.8 to 2.4. When $\alpha = 0.8$, D_t/D_0 will increase from 2.05 to 7.47 as γ enlarges from 0.3 to 1, when $\alpha = 2.4$, D_t/D_0 will increase from 3.22 to 11.73 as γ enlarges from 0.3 to 1.



Figure 4. Variation of Stiffness enhancing ratio D_t/D_0 calculated by Equation (10) as a function of γ .



Figure 5. Variation of Stiffness enhancing ratio D_t/D_0 calculated by Equation (11) as a function of γ .

3.2. The Buckling Coefficient k

The buckling coefficient *k* is a dimensionless number factor which depends on the type of loading, BCs, the plate aspect ratio β and the openings; and it has nothing to do with whether pasting FRP or not. When the plate aspect ratio is a integer, *k* is equal to 4 and 0.456 + $(b/a)^2$ for plates with BCs of 4S and 3S1F, respectively [16].

3.3. Comparison with the Experimental Results

In order to verify the theoretical method's results, a comparison with the authors' experimental results on the critical stress of perforated plates strengthened with FRP has been performed. As shown in Figure 6, a test rig capable of applying uniaxial compression on the plate was fabricated in the laboratory. And a classic perforated plate strengthened with FRP is shown in Figure 7. The unidirectional FRP has been used. The plates were constructed as following: Firstly, holes are drilled in plates if needed, then a single-layer of FRP is installed directly on the steel surface in the load direction and then another single-layer of FRP is installed in the direction of vertical load. The aim of this paper is to study the buckling of the plate, and when buckling, no delimitation of FRP has been found, so in the finite element (FE) models, the delimitation of FRP and the influence of manufacturing technology is not taken into account, and the FRP are considered an orthotropic material. The values of the basic geometry and material parameters are in Table 1, and the details of the experiments are in literature [19]. The elastic buckling loads P_{crExp} are determined from the experiments by two methods, and averaged; These two methods are the load-axial displacement curve and the load-transverse displacement squared method [14,20]. The experimental critical stress f_{crExp} can be calculated by $P_{crExp}/(bt_t)$, where t_t represents the total thickness of the plate. The comparison of experimental and theoretical results are in Table 2. From this table, it can be found that experimental and theoretical results are consistent when d/b and γ are equal to 0. It should be noted that the theoretical method cannot consider the influence of hole, which is a very important parameter to the elastic stability of perforated plates strengthened with FRP under uniaxial compression; so in calculating $f_{crEq.10}$ and $f_{crEq.11}$, k is taken as K_u , which is obtained by the numerical results in the following section. The values calculated by $f_{crEq.10}$ and $f_{crEq.11}$ are very closed, it proves the feasibility of Equation (11) in Calculating D_t . It also can be seen that when $\beta \neq 0$, the theoretical results of $f_{crEq.10}$ and $f_{crEq.11}$ are basically larger than experimental results f_{crExp} . The reason is that the theory method cannot consider the orthotropic properties of FRP, the effect of FRP's orthotropic properties is more obvious in the following numerical analysis, especially when pasting a lot of FRP.



Figure 6. Test rig setup.



Figure 7. A classic perforated plate strengthened with FRP.

 Table 1. The basic geometry and material parameters. BCs: Boundary conditions.

<i>a</i> (mm)	<i>b</i> (mm)	<i>t_s</i> (mm)	<i>t_f</i> 0 (mm)	E_s (MPa)	E _f (MPa)	v_s	v_f	f_y (MPa)	BCs
240	240	2.61	0.167	$2.0 imes 10^5$	$2.6 imes10^5$	0.3	0.25	287.9	4S

NO.	d/b	α	γ(×10 ⁻²)	D ₀ (×10 ⁵)	$D_{tEq.10}$ (×10 ⁵)	$D_{tEq.11}$ (×10 ⁵)	P _{crExp}	f _{crExp} .	f _{crEq.17}	$f_{crEq.23}$	<i>fcrFEM</i>	α _D	K _u	f _{crCal}	P_{uExp}	\mathbf{P}_{uCal}
1	0	1.225	0	3.35	3.35	3.35	55.83	89.13	88.08	88.08	87.79	1.000	4.00	88.08	94.02	87.59
2	0.1	1.225	0	3.35	3.35	3.35	53.78	85.85	83.08 *	83.08 *	83.06	1.000	3.77	83.08	90.11	85.41
3	0.3	1.225	0	3.35	3.35	3.35	47.89	76.45	72.95 *	72.95 *	72.93	1.000	3.31	72.95	79.66	80.71
4	0.5	1.225	0	3.35	3.35	3.35	45.12	72.03	71.62 *	71.62 *	71.60	1.000	3.25	71.62	74.44	80.06
5	0	1.225	6.4	3.35	4.18	4.14	64.88	97.34	103.24	102.25	95.66	0.946	4.00	96.76	101.86	97.03
6	0.1	1.225	6.4	3.35	4.18	4.14	61.90	92.87	97.39 *	96.45 *	89.84	0.934	3.77	90.13	96.64	94.13
7	0.3	1.225	6.4	3.35	4.18	4.14	58.28	87.45	85.51 *	84.68 *	76.08	0.910	3.31	77.13	92.72	87.99
8	0.5	1.225	6.4	3.35	4.18	4.14	50.54	75.83	83.95 *	83.15 *	72.44	0.887	3.25	73.77	78.35	86.30

Table 2. The comparison of experimental, theoretical, finite element (FE) models and Calculation method in this paper.

* *k* is taken as K_u , which is obtained by the numerical results in the following section.

3.4. The Proposed of Formula Correction

From the above comparison and analysis, we know that the stiffness obtained by theoretical calculation is too large due to not consider the orthotropic properties of FRP , so a modified factor of stiffness α_D is proposed, we consider α_D is a function of ω , β and d/b. It also can be seen from above analysis that the theory method cannot consider the influence of hole on buckling, so the buckling coefficient *k* needs to be modified as K_u which considers the change of d/b, β and the hole position e_x (and e_y). Therefore, the critical stress of perforated plates strengthened with FRP at buckling can be written as Equation (12). The values and change rules of α_D and K_u are the focuses of the following numerical simulation studies.

$$f_{cr} = K_u \frac{\pi^2}{t_t b^2} \left(\alpha_D D_t \right) = \alpha_D K_u \frac{\pi^2}{t_t b^2} D_t$$
(12)

4. Finite Element Model and Verification

4.1. Method of Analysis

The buckling problem of the plate is to solve the eigenvalue, the lowest eigenvalue corresponds to the critical stress while the eigenvector defines the deformed shape at buckling. A general-purpose finite element program, ABAQUS (Version: 6.11, Dassault Systemes S.A, Concord, MA, USA, 2011), has been utilized in the investigation to determine the elastic buckling loads of the perforated plates strengthened with FRP under uniaxial compression.

After performing the eigenvalue analysis, the lowest eigenvalue, λ_1 , can be expressed as Equation (13), so ($\alpha_D K_u$) can be expressed as Equation (14). It should be noted that α_D is equal to 1 as $\omega = 0$ regardless of other parameters, so we let $\omega = 0$ to study the variation of K_u as d/b, β , e_x/b , e_y/b and BCs change. We can change the form of α_D into Equation (15) to study the variation of α_D , where both the molecule and the denominator can be calculated by Equation (14) using numerical simulation results.

$$\lambda_1 = \alpha_D K_u \frac{\pi^2}{t_t b^2} D_t \tag{13}$$

$$\alpha_D(\omega, \beta, d/b, BCs)K_u(d/b, \beta, e_x/b, e_y/b, BCs) = \lambda_1 \frac{t_t b^2}{\pi^2 D_t}$$
(14)

$$\alpha_D(\omega, \beta, d/b, BCs) = \frac{\alpha_D(\omega, \beta, d/b, BCs)K_u(d/b, \beta, e_x/b, e_y/b, BCs)}{\alpha_D(0, \beta, d/b, BCs)K_u(d/b, \beta, e_x/b, e_y/b, BCs)}$$
(15)

4.2. Model Description

Four typical perforated plate models with β = 1, 2, 3 and 4 are shown in Figure 8. The plates have dimensions *a* and *b*, and circular hole with diameter *d*. The plate was discretized into a mesh consisting of S4R element, each node of the S4R element has six active degrees of freedom (DOF), including three translational DOFs and three rotational DOFs. The mesh size are dependent on the hole sizes and the extensive convergence studies in the following section.

The classical elastic-plastic material and lamina material are used to model the steel and FRP, respectively. The steel-FRP composite section can be easily created by increasing the number of layers and changing the material properties of each layer.

The BCs of the plates considered in this paper are 4S and 3S1F. For simple support edge, there will be no lateral edge displacements perpendicular to the plate's plane but rotations about the axis of each edge are allowed. And a uniform load of magnitude t_t per unit width is applied perpendicular to the width edges.



Figure 8. Examples of four typical plate models' meshes. (**a**) case of $\beta = 1$; (**b**) case of $\beta = 2$; (**c**) case of $\beta = 3$; (**d**) case of $\beta = 4$.

4.3. Model Verification

A mesh study is conducted to observe the convergence of buckling stress as the mesh size increases. Take plate with No. 8 in Table 2 for example, the Eigenvalues with respect to element numbers depended on mesh size, as shown in Figure 9. It can be seen that the curves starts to converge after an element number of 400, which with the general mesh sizes of b/10 and with mesh sizes in hole edge of $\pi d/40$. In order to reduce computational time, those meshes are adopted in the subsequent buckling analysis.

In order to verify FE models used in this study, a comparison with the authors' experimental results [19] on the critical stress of perforated plates strengthened with FRP has been performed, and the results are in Table 2. It can be seen that the FE models have a very limited relative absolute error in f_{cr} , in the consider cases, the errors are 13.0% on maximum, 4.05% on average, and 0.108 on the coefficient of variation. Therefore, the FE models are very accurate and enough for the studies in the next section.



Figure 9. Finite element (FE) models converge using mesh refinement.

5. Discussion of the Results

5.1. Buckling Mode Shapes

Figures 10–13 show the buckling mode shapes of perforated plates strengthened with FRP under uniaxial compression with different e_x/b , β , d/b and BCs. The following are found:

When the hole center is in the middle of the plate, the buckling mode shapes transform from two half-waves to three half-waves as d/b is enlarged from 0 to 0.7 (seen from Figure 10), a similar phenomenon occurs in all plates with $e_x/b = 1$ and 2.

- From Figures 10a and 11a, it can be seen that when *d/b* = 0.1, the buckling mode shapes are very close to plates without hole; and as *d/b* is enlarged, the max deformations of the plate are firstly centered near the hole then gradually transferred between hole and loading edge for 4S (seen from Figures 10b–d and 11b–d).
- From Figure 12, it can be seen that as *β* increases, the deformations of the existing part have been basically unchanged, and the deformations of the adding part are very small.
- From Figure 13, it can be seen that the mode shapes are essentially unchanged as *d*/*b* is enlarged for plates with BCs of 3S1F.



Figure 10. The model shapes of plates with $e_x/b = 1$, $\beta = 2$ and BCs of 4S. (**a**) d/b = 0.1; (**b**) d/b = 0.3; (**c**) d/b = 0.5; (**d**) d/b = 0.7.



Figure 11. The model shapes of plates with $e_x/b = 0.5$, $\beta = 2$ and BCs of 4S. (**a**) d/b = 0.1; (**b**) d/b = 0.3; (**c**) d/b = 0.5; (**d**) d/b = 0.7.



Figure 12. The model shapes of plates with $e_x/b = 0.5$, BCs of 4S and different β from 1 to 4. (a) $\beta = 1$; (b) $\beta = 2$; (c) $\beta = 3$; (d) $\beta = 4$.



Figure 13. The model shapes of plates with $e_x/b = 1$, $\beta = 3$ and BCs of 3S. (**a**) d/b = 0.1; (**b**) d/b = 0.3; (**c**) d/b = 0.5; (**d**) d/b = 0.7.

5.2. The Stiffness Modified Factor α_D

Due to the orthotropic properties of FRP, a modified factor of stiffness α_D has been proposed. And a reinforcement index of FRP has been defined as:

$$\omega = \frac{E_f t_f}{E_s t_s} = \frac{n_f E_f t_{f0}}{E_s t_s} \tag{16}$$

In order to study the variation of α_D , a large number of parametric analysis have been performed in this section, the parameters include reinforcement index ω , plate aspect ratio β , hole size d/b, and BCs. The parametric analysis' results are handled by the above approach in Equations (14) and (15), and plotted in Figures 14 and 15 for BCs of 4S and 3S1F, respectively.

Figure 14 shows the variation of α_D as the reinforcement index ω increases from 0 to 0.8 for different d/b from 0 to 0.7 and different plate aspect ratios β from 1 to 4 with the BCs of 4S. It can be seen that the curves in Figure 14 with the same d/b and different β are almost the same and the values in Figure 14a are slightly smaller. Therefore, α_D of a perforated rectangular plate strengthened by FRP with an integer aspect ratio could be conservatively taken as that of a perforated square plate, so only plates with $\beta = 1$ are analyzed in the following. From Figure 14a, it can be seen that α_D decreases as ω is enlarged and the curves can be coupled by Equation (17), where m, n are the undetermined second-order and first-order coefficients. For different d/b, m and n are calculated in Table 3, and R^2 is the correlation coefficient.

$$\alpha_D = m\omega^2 + n\omega + 1 \tag{17}$$

d/b	0	0.1	0.2	0.3	0.4	0.5	0.6	0.7
т	0.69	0.76	0.92	1.09	1.23	1.34	1.40	1.47
п	-0.74	-0.83	-1.04	-1.25	-1.44	-1.59	-1.69	-1.79
R^2	0.91	0.92	0.93	0.94	0.94	0.95	0.95	0.95

Table 3. The relationship d/b with m and n.

In Table 3, *m* and *n* can be approximate expressed as Equations (18) and (19), so α_D with BCs of 4S can be calculated by Equation (20)

$$m = 1.2(d/b) + 0.7 \tag{18}$$

$$n = -1.6(d/b) - 0.74 \tag{19}$$

$$\alpha_D = (1.2d/b + 0.7)\omega^2 - (1.6d/b + 0.74)\omega + 1$$
(20)

For the plates with BCs of 3S in Figure 15, the same method can be used to study the variation in α_D as the reinforcement index ω increases from 0 to 0.8 for different plate aspect ratios β from 1 to 4. It also can be seen that α_D decreases as ω is enlarged. It is noted that α_D is independent of d/b when $\beta = 2$, 3 and 4, so α_D with BCs of 3S1F can be calculated by Equation (21).

$$\alpha_{\rm D} = \begin{cases} (0.28d/b + 0.56)\omega^2 - (0.33d/b + 0.69)\omega + 1, \ \beta = 1\\ 1.18\omega^2 - 1.45\omega + 1, \ \beta = 2\\ 1.42\omega^2 - 1.73\omega + 1, \ \beta = 3\\ 1.51\omega^2 - 1.84\omega + 1, \ \beta = 4 \end{cases}$$
(21)



Figure 14. α_D vs. ω for perforated rectangular plate ($\beta = 1, 2, 3, 4$) with the BCs of 4S.



Figure 15. α_D vs. ω for perforated rectangular plate ($\beta = 1, 2, 3, 4$) with the BCs of 3S1F.

5.3. The Buckling Coefficient K_u

5.3.1. Hole Position Varying Along *x*-Axis

Figure 16a~d show the change in K_u as hole position parameter e_x/b is enlarged with the BCs of 4S, for different d/b from 0 to 0.7 and different plate aspect ratios β from 1 to 4. It can be seen that hole position parameter e_x/b has a little effect on K_u for d/b = 0.1, especially as the plate aspect ratios β enlarged. It also can be seen that all the curves are very similar in trends and values with the same d/b when e_x/b is less than 0.5. It means that when β is an integer greater than 1 and e_x/b is less than 0.5, the buckling coefficient K_u can be taken as a square plate with the same width b, the same hole d/b, and the same FRP pasting. It also can be seen from the figures that as e_x/b is enlarged from 0 to 0.5, K_u decreases but gradual slow for d/b is 0.1, and first increases and then decrease for d/b is 0.2 and 0.3, and drastically increases for d/b is from 0.4 to 0.7. It means that when d/b is bigger than 0.4 and e_x/b is less than 0.5, the buckling coefficient K_u is very unstable, so it should be avoided to place big hole too close to the loading edge. For e_x/b is bigger than 0.5, K_u shows little variation with d/b and e_x/b , except when $\beta = 2$ and d/b is bigger than 0.4. Figure 16b shows that for $\beta = 2$ and $e_x/b = 1$, K_u increases from 4.48 to 5.99 as d/b is enlarged from 0.4 to 0.7, this also proves the necessity of placing the big hole in the middle area of the plate.

Figure 17a~d show the change in K_u as hole position parameter e_x/b is enlarged with the BCs of 3S1F, for different d/b from 0 to 0.7 and different plate aspect ratios β from 1 to 4. In general, K_u of plates with BCs of 3S1F is very small, and the curves shows that K_u keeps almost constant irrespective of e_x/b when $\beta = 1$, 2, and $\beta = 3$ (or 4) with d/b less than 0.3. When $\beta = 3$ (or 4) and

 $d/b \ge 0.3$, K_u increases as the hole moves away from the loading edge towards its center along *x*-axis. Take $\beta = 3$ and d/b = 0.7 for example, as e_x/b is enlarged from 0.375 to 1.5, K_u increases from 0.41 to 0.47. Therefore, it is also recommended to place the big hole close to the middle area of the plate in *x*-axis with the BCs of 3S1F when β is bigger than 3.



Figure 16. K_u vs. e_x/b for perforated rectangular plate ($\beta = 1, 2, 3, 4$) with the BCs of 4S.



Figure 17. Cont.



Figure 17. K_u vs. e_x/b for perforated rectangular plate ($\beta = 1, 2, 3, 4$) with the BCs of 3S1F.

5.3.2. Hole Position Varying Along *y*-Axis

Figure 18 shows the change in K_u as hole position parameter e_y/b is enlarged from 0.5 to 0.9 with the BCs of 4S, for different d/b from 0 to 0.7, different plate aspect ratios β from 1 to 4, and different $e_x/b = 0.5, 1.0, 1.5, and 2.0$. The center of the hole is halfway between nodal cross lines when $e_x/b = 0.5$ and 1.5, and the center of the hole is at the nodal cross line when $e_x/b = 1$ and 2. It could be noticed that for a small hole d/b (e.g., d/b = 0.1 and 0.2), there are little effect of the hole location e_y/b on K_u , and the effect becomes more pronounced as hole size d/b increases. When $d/b \ge 0.2$ in Figure 18a~d, it can be seen that the curves are very similar irrespective of β , that is: with the increasing of e_y/b , K_u will decrease dramatically as e_y/b is enlarged, especially when $d/b \ge 0.4$. Take $\beta = 3$, $e_x/b = 0.5$ and d/b = 0.6 for example, as e_y/b is enlarged from 0.5 to 0.65, K_u decreases from 3.28 to 2.20, this decline is very severe. The curves with $e_x/b = 1$ in Figure 18e~g are also very similar, and the law of changes are similar to curves with $e_x/b = 0.5$ (Figure 18a~d). What's different is that for curves in Figure 18e~g, when e_y/b is less than 0.575, K_u is bigger than 4 which is the K_u of plate without a hole. The law of changes in Figure 18h~j are between figures discussed above.

Figure 19 shows the change in K_u as hole position parameter e_y/b is enlarged from 0.05 to 0.95 with the BCs of 3S1F, for different d/b from 0 to 0.7 and different plate aspect ratios β from 1 to 4. It can be seen that the K_u of plates with BCs of 3S1F is much less than those with BCs of 4S, especially when β is bigger than 1. The figure also shows that when the hole size d/b is smaller than a specific value (for $\beta = 1$: 0.1; $\beta = 2$: 0.1; $\beta = 3$: 0.3; $\beta = 4$: 0.4), there is a little effect of the hole position e_y/b on buckling coefficient K_u and the d/b effect becomes more pronounced as d/b increases.

Over all, for a small hole size, there is little effect of the hole position e_y/b on buckling coefficient K_u regardless of the BCs (4S or 3S1F), and becomes more pronounced as d/b increases. The plate's stability improves as the hole places near the middle of the plate for 4S and the simple support edge for 3S1F in *y*-axis.



Figure 18. Cont.

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Figure 18. K_u vs. e_y/b for perforated rect38angular plate ($\beta = 1, 2, 3, 4$) with the BCs of 4S.



Figure 19. Cont.





Figure 19. K_u vs. e_u/b for perforated rectangular plate ($\beta = 1, 2, 3, 4$) with the BCs of 3S1F.

5.3.3. Hole Size *d*/*b*

Figure 20 show the change in K_{μ} as d/b is enlarged from 0 to 0.7 with the BCs of 4S. When e_x/b is 1 or 2, the center of opening is in the node point. And when e_x/b is 0.5 or 1.5, the center of opening is in the middle of the node points. As shown Figure 20, two different behaviors can be seen depending on the hole in the node point or in the middle of the node points. When the hole is at the node point (e_x/b) is 1 or 2), it is evident from the figure that K_u show an increasing trend as d/b is enlarged, for $\beta = 2$, K_u will increase from 4.02 to 5.99 as d/b is enlarged from 0 to 0.7, it is much better than that of all other cases. It also can be seen some discontinuities in the curves. Such discontinuities can be explained by the buckling mode shapes changing of the plate under compression, take $\beta = 2$ and $e_x/b = 1$ for example, when d/b increases from 0 to 0.7, the buckling mode has changed from 2 half-waves to 3 half-waves (seen from Figure 10). On the other hand, when the hole is in the middle of the node points $(e_x/b \text{ is } 0.5 \text{ or } 1.5), K_u$ of plates with $e_x/b = 1.5$ first decrease slowly and then increase as d/b is enlarged, and decreases gradually as d/b is enlarged for plates with $e_x/b = 0.5$. At the same time, we found that the K_u of plates with $e_x/b = 1.5$ are always bigger than those of plates with $e_x/b = 0.5$. For a plate under uniaxial compression, the nodal cross-lines (the dotted line in the figures) act as imaginary supports for the plate in the lateral direction and the resistance to buckling is provided by the bending stiffness of the plate parts between these nodal cross lines, when the center of opening is in the node point, the bending stiffness will show little reduction, and when the hole is in the middle of the node points, it will reduce its bending stiffness and then reduce its buckling load and buckling coefficient, and K_u will further reduce as d/b increases, this can explain the phenomenon of K_u with $e_x/b = 1$ and 2 are bigger than those of plates with $e_x/b = 0.5$ and 1.5. Therefore, it is recommend to put the hole in the node point in design.

Figure 21 show the change in K_u as d/b is enlarged from 0 to 0.7 with the BCs of 3S1F. It is obviously that the buckling coefficient K_u of all the plates decrease as d/b is enlarged, especially for $e_x/b = 0.5$ and $\beta = 1$, K_u decreases from 1.42 to 0.55 as d/b is enlarged from 0 to 0.7. Unlike the BCs of 4S, where the change rules of K_u relies on β and e_x/b , K_u decreases rapidly as d/b increases irrespective of β and e_x/b for plates with BCs of 3S1F, This phenomenon also can be explained as the buckling mode shapes are the same in all cases (seen from Figure 13). Meanwhile, by comparing Figure 21 with Figure 20, it is clear that K_u of plates with BCs of 3S1F are far less than those of 4S.





Figure 20. K_u vs. d/b with BCs of 4S. (a): $e_x/b = 1$ and 2; (b): $e_x/b = 0.5$ and 1.5.



Figure 21. K_u vs. d/b with BCs of 3S1F. (a): $e_x/b = 1$ and 2; (b): $e_x/b = 0.5$ and 1.5.

6. An Example for the Usage of the Proposed Method

Example calculation for the ultimate capacity of a perforated plate strengthened with FRP under uniaxial compression, according to the winter equations in (22)~(24), where ρ is a slenderness reduction factor, λ is the plate slenderness, f_y is the yield stress of the steel, and f_{cr} is the critical stress. f_{cr} can be calculated using this method of this paper.

$$\rho = \frac{\lambda - 0.22}{\lambda^2} \le 1 \tag{22}$$

$$\lambda = \sqrt{\frac{f_y}{f_{cr}}} \tag{23}$$

$$P_u = \rho f_y b t_t \tag{24}$$

The basic parameters of the example (Table 2, No. 7):

a = b = 240 mm, d/b = 0.1, $e_x/b = 0.5$, $e_x/b = 0.5$, $t_s = 2.61$ mm, $t_{f0} = 0.167$ mm, $n_f = 1$, $E_s = 2.0 \times 10^5$ MPa, $E_f = 2.6 \times 10^5$ MPa, $v_s = 0.3$, $v_f = 0.25$, $f_y = 287.8$ MPa, BCs: 4S.

The critical stress *f*_{crCal}:

 $t_{t} = 2.61 + 0.167 = 2.777 \text{ (mm)}, \alpha = \frac{E_{f}/(1-v_{f}^{2})}{E_{s}/(1-v_{s}^{2})} = 1.225, \beta = \frac{t_{f}}{t_{s}} = 0.064 < 0.2,$ $D_{0} = C_{0} - B_{0}^{2}/A_{0} = E_{s}t_{s}^{3}/12(1-v_{s}^{2}) = 3.35E5, D_{t}/D_{0} = 1 + 3\alpha\beta = 1.24, D_{t} = 4.14E5,$ $e_{x}/b = 0.5, \beta = 1, d/b = 0.1 \text{ from Figure 20}, K_{u} = 3.77,$ $\omega = \frac{n_{f}E_{f}t_{f0}}{E_{s}t_{s}} = 0.0784, \text{ from Equation (20)}, \alpha_{D} = (1.2d/b + 0.7)\omega^{2} - (1.6d/b + 0.74)\omega + 1 = 0.947,$ $f_{crCal} = \alpha_{D}K_{u}\frac{\pi^{2}}{t_{t}b^{2}}D_{t} = 90.13 \text{ (MPa)}.$

The ultimate capacity *P*_{uCal}:

$$f_y = 287.8MPa, \ \lambda = \sqrt{\frac{f_y}{f_{cr}}} = 1.79, \ \rho = \frac{\lambda - 0.22}{\lambda^2} = 0.491 \le 1, \ P_{uCal} = \rho f_y bt_t = 94.12 \ (KN).$$

The experimental values of critical stress f_{crExp} and ultimate load P_{uExp} are 92.87 MPa and 96.64 KN, it is obvious that the values of f_{crCal} and P_{uCal} are very close to f_{crExp} and P_{uExp} , respectively. Other plates are also calculated using this method and plotted in Table 2.

7. Conclusions and Recommendations

In order to study the elastic stability and propose a calculation method of perforated plates strengthened with FRP under uniaxial compression, series of FE models are built and several parameters are considered in the study: material's geometrical and mechanical properties, boundary conditions, plate aspect ratio, hole sizes, and hole position. A calculation method of buckling stress is proposed and modified based on the theory of composite plate and the numerical results. In addition to the detailed discussion and conclusions given in the previous section, some general conclusions and recommendations can be summarized as follows:

- The theory of composite plate proposed by Pister is available to calculate the flexural rigidity of the pure steel plate and the perfectly bonded layered plate, but it should be modified when calculating the flexural rigidity of the perfectly bonded layered plate due to not considered the orthotropic properties of FRP.
- The buckling coefficient *k* needs to be modified as K_u which considers the change of plate aspect ratio β , hole sizes d/b and the hole position e_x (or e_y).
- For BCs of 4S, when e_x/b is an integer, the buckling mode shapes change as d/b is enlarged, which cause some discontinuities in K_u vs. d/b curves (Figure 20). And for BCs of 3S1F, the buckling mode shapes is unchanged irrespective of d/b.
- For BCs of 4S, as e_x/b is enlarged from 0 to 0.5, K_u decreases but gradual slow for d/b is 0.1, and first increases and then decrease for d/b is 0.2 and 0.3, and drastically increases for d/b is enlarged from 0.4 to 0.7. It means that when d/b is bigger than 0.4 and e_x/b is less than 0.5, the buckling coefficient K_u is very unstable, so it should be avoided to place big hole too close to the loading edge in *x*-axis, especially when d/b is bigger than 0.4.
- For BCs of 3S1F, K_u is very small and stable irrespective of e_x/b when $\beta = 1, 2, \text{ and } \beta = 3$ (or 4) with d/b less than 0.4. When $\beta = 3$ or 4 and d/b is bigger than 0.4, K_u increases as the hole moves away from the loading edge towards its center along *x*-axis. Therefore, it is recommended to place the big hole ($d/b \ge 0.4$) close to the middle area of the plate in *x*-axis, especially when d/b is bigger than 0.4.

- For a small hole size, there is little effect of the hole position e_y/b on buckling coefficient K_u regardless of the BCs (4S or 3S1F), and that effect becomes more pronounced as d/b increases. So it is recommended to put the holes near the middle of the plate for 4S and the simple support edge for 3S1F in *y*-axis.
- As shown in Table 2, the buckling stresses and ultimate loads calculated by the proposed method of this paper combining with the winter equation are very close to the experimental values.

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Glossary/Nomenclature/Abbreviations

a,b	the lengths in the longitudinal (transverse) directions
t_s	the thicknesses of the steel plate
t_f	the total thicknesses of FRP sheet
n _f	FRP Layers' number
t_{f0}	single layer of FRP sheet's thickness
d	diameter of the hole
$e_{x(y)}$	the distance along $x(y)$ -axis from the center of the hole to the left(bottom) edge
E_S, E_f	the elastic modulus of steel and FRP
v_s, v_f	the Poisson ratio of steel and FRP
β	Plate aspect ratio, equal to b/a
k	the buckling coefficient
D_0, D_t	the flexural rigidity of plate without FRP and strengthened with FRP
<i>f</i> _{cr}	the critical stress of the plate
P _{cr}	the critical load of the plate
P_u	the ultimate load of the plate
α	reinforcement index of FRP, calculated by Equation (8)
γ	reinforcement index of FRP, calculated by Equation (9)
ω	reinforcement index of FRP, calculated by Equation (16)
K _u	the modified buckling coefficient

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