

# Article

# Experimental Analyses on Multiscale Structural and Mechanical Properties of $\epsilon$ -Si/GeSi/C-Si Materials

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**Abstract:** Strained silicon ( $\varepsilon$ -Si) is a promising material that could extend Moore's law by enhancing electron mobility. A  $\varepsilon$ -Si material is usually composed of multiscale, multilayer heterostructures, where the strained-silicon film or strap is tens-of-nanometers thick, and its buffer layers are of the micrometer scale. The structural properties determine the electrical performance and reliability of  $\varepsilon$ -Si-based devices. Inhomogeneous residual stress is induced during the preparation, which induces  $\varepsilon$ -Si structure failure. In this work, biaxial strained-silicon films that contain graded and relaxed germanium-silicon buffer layers were prepared on monocrystalline silicon wafers through reduced-pressure chemical-vapor epitaxy. The layer components and thicknesses were measured using energy-dispersive spectroscopy and scanning-electron microscopy. Crystal and lattice characters were observed by using high-resolution transmission-electron microscopy and micro-Raman spectroscopy. The residual stress distribution along cross-sections of the  $\varepsilon$ -Si multilayer structures was examined by using micro-Raman mapping. The experimental results showed that, with a gradual increase in germanium concentration, the increasing residual stress was suppressed owing to dislocation networks and dislocation loops inside the buffer layers, which favored the practical application.

**Keywords:** strained silicon (ε-Si); multiscale; structural property; mechanical property; micro-Raman spectroscopy; cross-section; dislocation

# 1. Introduction

Strained silicon ( $\varepsilon$ -Si) is regarded as a promising material that could extend Moore's law [1]. Because introduced strain leads to an increased carrier mobility and reduced resistance [2], the performance of  $\varepsilon$ -Si-based devices is improved significantly [3]. In recent years, with continuous developments in  $\varepsilon$ -Si technology, semiconductor chips have become smaller, faster, and more energy efficient [4]. Strained-silicon technology has been applied widely in core chips inside mobile phones, personal computers, laptops, tablets and even televisions [5].

The  $\varepsilon$ -Si structure complexity determines the difficulty of the manufacturing process. In general,  $\varepsilon$ -Si is achieved by manufacturing a set of multiscale, multilayer heterogeneous twoor three-dimensional structures [6]. The manufacture of such a complex structure requires a strict and well-designed control of process parameters to ensure that each processing procedure step



may yield a material/structure that meets the design requirements of geometrical dimensions and structural-mechanical properties [7,8]. Furthermore, the design, manufacture, performance evaluation and control technologies for microelectronics with non-uniform and local-singular stress distribution is regarded as the future of the microelectronics industry [9].

The (equal biaxial or uniaxial) strain state and its magnitude determine the optical/electrical properties of  $\varepsilon$ -Si-based devices [10]. The strain state and magnitude of the  $\varepsilon$ -Si depend heavily on the buffer-layer properties (such as the element composition, thickness, lattice quality, defects and residual stress) and interfacial properties (such as mismatch and shear stress) [11]. Size miniaturization of a  $\varepsilon$ -Si-based device leads to increased defect densities and residual stress, which results in a degraded optoelectronic performance and structure reliability of the entire integrated device [12].

Experimental studies of multilayer heterogeneous structures, such as strained-silicon, face two major challenges, namely, "multiscale" and "multiproperty coupling". A strained-silicon layer is usually several to tens-of-nanometers thick, whereas the buffer-layer thickness tends to be of the micrometer scale, and the substrate is typically several hundreds of microns thick. Unique key physical and mechanical behaviors that interact exist at each spatial scale. Cross-scale interactions between the multiscale physical–mechanical properties increase the complexity of experimental studies of strain silicon.

A number of methods at different spatial scales have been applied to experimental studies on multilayer heterogeneous structures [13–16]. The profiler has been used to measure the average residual stress from the mesoscale [17]. By using (scanning- and transmission-) electron microscopy, the structural morphology and dislocation distribution in the microscale have been determined [18]. At the same scale, nano-indentation has been used to measure the modulus [19], whereas Raman microscopy has been applied to characterize the lattice quality and residual stress [20–22]. High-resolution transmission-electron microscopy (TEM) has been used to analyze the lattice structure and local strain of materials at the nanoscale [23].

However, most widely used experimental methods can measure or characterize only one or several material parameters at their respectively applicable scales. These techniques do not provide sufficient understanding of the multiscale structure properties, such as strained silicon, and it is difficult to discover the interaction between the different structural and mechanical properties at different spatial scales. Therefore, there is a need for collaborative analysis using methods at different spatial scales to investigate the interaction mechanisms of multiproperty coupling.

In this work, a biaxial strained-silicon material with germanium–silicon as a buffer layer was prepared and studied experimentally. The structural properties, including the components and thickness of each layer at the microscale and the crystal and lattice characteristics and quality at the nanoscale, was characterized by energy-dispersive spectroscopy (EDS), scanning-electron microscopy (SEM), micro-Raman spectroscopy (MRS) and TEM. The residual stress distribution at different depths was analyzed by using micro-Raman mapping on the sample cross-section.

#### 2. Materials and Methods

The  $\varepsilon$ -Si/GeSi/C-Si materials were prepared through reduced pressure chemical vapor epitaxy (CVE) [24]. SiH<sub>2</sub>Cl<sub>2</sub> was used as the Si precursor, GeCl<sub>4</sub> was the Ge precursor, and H<sub>2</sub> was the carrier gas. The monocrystalline Si (100) wafer was pre-treated using an H<sub>2</sub>SO<sub>4</sub>:H<sub>2</sub>O<sub>2</sub> = 3:1 solution for 10 min to remove contaminants and then dipped in 10% HF solution to remove the oxide layer. The CVE growth temperature was 900 °C, and the pressure was  $1.33 \times 10^4$  Pa. In the beginning, the Ge component increased linearly to form a graded germanium-silicon (Ge<sub>x</sub>Si<sub>1-x</sub>, x = 0→0.2) buffer layer of ~3 µm thickness. The growth rate slowed and the Ge concentration was fixed at 20% to form a Ge<sub>0.2</sub>Si<sub>0.8</sub> buffer layer of ~0.8 µm thickness. A 10 nm strained-silicon cap layer was deposited at 750 °C and 1.33 × 10<sup>4</sup> Pa. Figure 1a provides a schematic structure of the prepared  $\varepsilon$ -Si/GeSi/C-Si material.

The cross-sectional samples were prepared as follows. A Si-based  $\epsilon$ -Si/GeSi wafer was cut to approximately 10 mm  $\times$  10 mm square pieces. Two wafers were pasted by using epoxy resin with

the  $\varepsilon$ -Si films face-to-face. Then, two C-Si pieces, cut from another (100) silicon wafer with a 430 µm thickness were pasted outside the two  $\varepsilon$ -Si/GeSi wafers, respectively. Then, this piece group was cut to some thin slices about a ~1 mm thickness. The two cross-section surfaces of each thin slice were polished by using the chemical-mechanical polishing technique (CMP). Samples for TEM analysis were thinned further by using sandpaper (3M, Maplewood, Minnesota, USA), lapping film (Allied High Tech Products, Rancho Dominguez, CA) and a precise ion-polishing machine (Gatan, Pleasanton, CA, USA).

A ZEISS EVO MA15 system and a JEOL 2100 system were used for the SEM and TEM tests, respectively. Raman experiments for the cross-sectional samples were conducted by using a laser confocal micro-Raman system (InVia Reflex, Renishaw, Gloucestershire, UK) with a 532 nm laser exciting light. A  $50 \times$  Leica objective (N.A. 0.8) was used, and hence, the spot size of the incident laser was about 1  $\mu$ m in diameter. Besides, Raman tests on the untreated surface of the strained-silicon samples used the same system, but with a 325-nm laser and a  $40 \times$  Leica ultraviolet (UV) objective (N.A. 0.5). The sampling spot of the incident laser was about 5  $\mu$ m in diameter.



**Figure 1.** (a) Schematic structure of prepared  $\varepsilon$ -Si/GeSi/C-Si material, (b) scanning-electron microscopy (SEM) image of multilayer vicinity of sample cross-section. CVE, chemical vapor epitaxy.

## 3. Results and Discussion

#### 3.1. Microscale Structural Properties

Figure 1b shows an SEM image of the multilayer vicinity of the cross-sectional sample. Because of differences in the electrical resistance of each layer, distinctive monocrystalline silicon,  $Ge_xSi_{1-x}$ ,  $Ge_{0.2}Si_{0.8}$  and epoxy resin layers were observed. A narrow black strip of ~1 µm thickness is the adhesive layer of epoxy resin that adheres the two multilayer wafers. The region with uniform and relatively deeper greyscale, at the left of the image and far from the adhesive layer, is the C-Si substrate. From the C-Si substrate to the wafer edge (~3.6 µm thickness), the grayscale becomes brighter gradually, which is the graded and relaxed buffer layers. The C-Si/Ge<sub>x</sub>Si<sub>1-x</sub> interface is clear and straight, but that of the Ge<sub>x</sub>Si<sub>1-x</sub>/Ge<sub>0.2</sub>Si<sub>0.8</sub> interface is barely distinguishable visible in this SEM image. The strained-silicon layer (~10 nm thick) is not recognized in the SEM image, because of the resolution limitation of the SEM system.

The SEM image also shows that there is no visible fluctuation on the surface of the sample cross-section. The cross-sectional surface roughness was measured by using an atomic-force microscope (AFM, Dimension Icon, Bruker, Billerica, Massachusetts, USA) and a white-light interferometer (WLI, NewView 8300, Zygo, Middlefield, Connecticut, USA). Figure 2a gives the height data of a random 500 nm  $\times$  500 nm area on the cross-sectional surface detected by AFM. The root-mean-square deviation of the height data (Rq) is 0.43 nm, and the arithmetical mean deviation of the height data (Ra) is 0.33 nm.

Figure 2b gives the height data of a random 21  $\mu$ m × 21  $\mu$ m area on the cross-sectional surface detected by using a white-light interferometer with a spatial resolution of 0.34  $\mu$ m in the *X*/*Y* direction and 0.01 nm in the *Z* direction (root-mean-square repeatability). R<sub>q</sub> and R<sub>a</sub> are essential surface-roughness parameters. The experimental results showed that the surface roughness of the sample cross-section was less than 0.5 nm, regardless of the nano- or microscale area.



**Figure 2.** The surface roughness of cross-sectional surface measured by (**a**) atomic-force microscope (AFM) and (**b**) white-light interferometer.

The elemental distribution along the depth direction of the sample cross-section is shown in Figure 3 as determined by EDS (Oxford Instruments, Oxford, UK) with a 42-nm step length. When the EDS sampling point on the cross-section surface approaches the wafer edge, the elemental silicon content exhibits an approximately linear decrease and that of elemental germanium shows a linear increase. Such linear variations of the elemental components that are sustained inside a region of ~2.6 µm in thickness, corresponding to the graded buffer layer (via the Ge<sub>x</sub>Si<sub>1-x</sub> layer).



**Figure 3.** Elemental distribution along the depth direction achieved by using energy-dispersive spectroscopy (EDS).

The content of elemental silicon increased slightly and gradually until the sampling point reached the wafer edge. The opposite behavior resulted in the content of elemental germanium. The thickness of this region was ~1  $\mu$ m, which corresponds to the relaxed buffer layer (viz. the Ge<sub>0.2</sub>Si<sub>0.8</sub> layer). In this layer, the actual distributions of elemental Si and Ge did not agree with, but were not distinct

from, the design shown in Figure 1. The EDS result showed that it was difficult to recognize the interface between the two buffer layers, and the exact thickness of the two buffer layers. The sudden drop of the Si component at a 1  $\mu$ m depth may have been caused by parameter alternation during the ultra-high-vacuum chemical-vapor epitaxy, which helps to locate the Ge<sub>x</sub>Si<sub>1-x</sub>/Ge<sub>0.2</sub>Si<sub>0.8</sub> interface.

Both elements decreased to zero because the EDS sampling point entered the epoxy resin layer of the sample cross-section. The EDS data did not show the signal from the  $\varepsilon$ -Si layer, which should be because the  $\varepsilon$ -Si layer thickness (~10 nm) was less than the EDS sampling step length.

Figure 4a shows the wavenumber image of the Si-Si band (~520 cm<sup>-1</sup>) that was obtained by Raman mapping at a 25  $\mu$ m × 24  $\mu$ m region near the multilayer interfaces on the sample cross-section. The step lengths for Raman mapping along the parallel and vertical direction to the interfaces were 1.2  $\mu$ m and 0.2  $\mu$ m, respectively. Figure 4b shows the mean values of the wavenumbers at different locations along the depth direction based on Figure 4a. Figure 4a,b shows that the wavenumber at ~520 cm<sup>-1</sup> changes slightly inside the Si substrate, acutely in the buffer layers, and continuously at the interfaces. When the scanning spot originated from the Si substrate into the buffer layers, the wavenumber showed a sharp and linear decrease with a speed of  $-4.97 \pm 0.07$  cm<sup>-1</sup>/ $\mu$ m. However, this rapidly decreasing trend stopped within a region of ~1  $\mu$ m next to the epoxy layer (compared with the top surface of the strained-silicon wafer), where the wavenumber decreased gently with a speed of  $-0.83 \pm 0.28$  cm<sup>-1</sup>/ $\mu$ m. Hence, from the wavenumber image, the thickness of the rapidly decreased wavenumber region is ~2.6–2.8  $\mu$ m, and that of the gently decreased wavenumber region is ~0.8–1  $\mu$ m. The former region is the Ge<sub>x</sub>Si<sub>1-x</sub> layer and the latter is the Ge<sub>0.2</sub>Si<sub>0.8</sub> layer. The strain distribution can be achieved by using the wavenumber results, as discussed below.



**Figure 4.** (**a**) Si-Si band wavenumber image obtained by Raman mapping on the sample cross-section, (**b**) wavenumber distribution along the depth direction.

Figure 5a shows the Raman intensity image of the same region in Figure 4a. Figure 5b shows the mean values at different locations along the depth direction based on Figure 5a. In Figure 5, when the sampling spot for Raman mapping from the Si substrate crossed the buffer layer, the Raman intensity of the Si-Si band decreased first and then increased. The Raman-intensity-decreased region in Figure 5 is almost the same as that of the wavenumber rapid-decreased region in Figure 4, so the Raman intensity-increased region to the wavenumber gently-decreased region. This result indicates that the quantity of elemental Si in the monocrystalline stage decreased with a continuous enhancement of elemental Ge in the graded-buffer ( $Ge_xSi_{1-x}$ ) layer, and then improved gradually

inside the relaxed-buffer ( $Ge_{0.2}Si_{0.8}$ ) layer even though the elemental content of the materials did not change any more.

Figure 6a shows the full-width-at-half-maximum (FWHM) image of the same region of Figure 4a. Figure 6b gives the mean values at different locations along the depth direction based on Figure 6a. In Figure 6, the FWHM of the Si-Si band broadened first and then narrowed when the sampling spot of Raman mapping from the Si substrate crossed the buffer layer. The FWHM broadened region in Figure 6 is different from (smaller than) the wavenumber rapid-decreased region in Figure 4 and the Raman-intensity-decreased region in Figure 5. Therefore, with a continuous enhancement of elemental Ge, the average lattice quality of the Si element in the monocrystalline stage in the graded-buffer ( $Ge_xSi_{1-x}$ ) layer decreased to an extreme and then improved continuously. This tendency was retained in the relaxed-buffer ( $Ge_0.2Si_{0.8}$ ) layer.



**Figure 5.** (a) Si-Si band Raman-intensity image obtained by Raman mapping on the sample cross-section, (b) Raman-intensity distribution along the depth direction.



**Figure 6.** (**a**) Si-Si band full-width-at-half-maximum (FWHM) image obtained by Raman mapping on the sample cross-section, (**b**) FWHM distribution along the depth direction.

### 3.2. Nanoscale Structural Properties

The structural properties at the nanoscale were investigated by using TEM. Figures 7 and 8 give high- and low-resolution TEM images of the sample cross-section, respectively. Figure 7a shows a high-resolution TEM image, including the  $\varepsilon$ -Si and Ge<sub>0.2</sub>Si<sub>0.8</sub> layers. Three regions exist, each of which has a uniform contrast grade, and which corresponds to a dissimilar layer. The left lower region, which is brightest, is an epoxy resin layer of the sample that serves as the adhesive for the two wafers. The image shows no regular lattice structure in this region. The region adjacent to the epoxy is the  $\varepsilon$ -Si layer. The largest region in the image is the Ge<sub>0.2</sub>Si<sub>0.8</sub> layer. Inside the  $\varepsilon$ -Si layer, the lattice structure is complete and clear, and shows a regular mesh of the monocrystalline state. The growth surface is in the same atomic sheet. No visible defects exist in this layer. Moreover, the interface of the  $\varepsilon$ -Si layer thickness is relatively uniform at ~8.67 nm. Figure 7 also shows that the cross-section is a (100) crystal plane.



**Figure 7.** High-resolution transmission-electron microscopy (TEM) images of sample cross-section (**a**) at the  $\varepsilon$ -Si and Ge<sub>0.2</sub>Si<sub>0.8</sub> layers, (**b**) inside the Ge<sub>0.2</sub>Si<sub>0.8</sub> layer, (**c**) at the Ge<sub>0.2</sub>Si<sub>0.8</sub>/Ge<sub>x</sub>Si<sub>1-x</sub> interface and (**d**) inside the Ge<sub>x</sub>Si<sub>1-x</sub> layer.



**Figure 8.** Low-resolution TEM images of sample cross-section (**a**) around  $\text{Ge}_{0.2}\text{Si}_{0.8}$  layer, (**b**,**c**) at  $\text{Ge}_x\text{Si}_{1-x}/\text{C-Si}$  interface and (**d**) inside  $\text{Ge}_x\text{Si}_{1-x}$  layer with a dislocation loop.

Figure 8a shows a low-resolution TEM image around the  $Ge_{0.2}Si_{0.8}$  layer, and Figure 7b shows a high-resolution enlarged image of spot **A** in Figure 8a, which is inside the  $Ge_{0.2}Si_{0.8}$  layer. In the low-resolution TEM image, no obvious defect exists inside the  $Ge_{0.2}Si_{0.8}$  layer. The interface between the  $Ge_xSi_{1-x}$  and  $Ge_{0.2}Si_{0.8}$  layers is clearly recognizable, which shows that the  $Ge_{0.2}Si_{0.8}$  layer thickness is uniform, at ~0.96 µm (~1µm). According to the SEM, EDS and Raman results, as shown in Figures 1 and 3–7, the total thickness of the two GeSi buffer layers is 3.6 µm, and hence, that of the graded layer ( $Ge_xSi_{1-x}$ ) is ~2.6 µm.

In the high-resolution TEM images Figure 7a,b, the lattice phenomena in the  $Ge_{0.2}Si_{0.8}$  layer are analogous to those in the  $\varepsilon$ -Si layer. The two layers have the same lattice structure and dimensions and can only be distinguished according to differences in the contrast grades. The lattice maintained a good and similar crystal structure, which indicates that the stress in the  $\varepsilon$ -Si and  $Ge_{0.2}Si_{0.8}$  layers is not released by defects or dislocations.

Figure 8b,c presents low-resolution TEM images around the  $Ge_xSi_{1-x}$  layer. Figure 7c,d presents high-resolution enlarged images at the  $Ge_xSi_{1-x}/Ge_{0.2}Si_{0.8}$  interface and at spot **B** in Figure 8b, respectively, where spot **B** is inside the  $Ge_xSi_{1-x}$  layer. In the low-resolution TEM image in Figure 8b,c, the interface between the  $Ge_xSi_{1-x}$  layer and C-Si substrate appears clear and straight. Numbers of dislocations exist in the  $Ge_xSi_{1-x}$  layer, especially close to the  $Ge_xSi_{1-x}/Si$  interface. Some appear as dislocation loops, such as the one in Figure 8d. Most dislocations connect and form dislocation networks inside the layer. Some vertical threading dislocation lines reach the  $Ge_xSi_{1-x}/C$ -Si interface.

Dislocation lines can only be terminated on the inner surface of the interface or on the free crystal surface, but they may not stop inside the crystal and appear in the form of isolated sections. Therefore, inside the crystal, any dislocation must be closed to form a dislocation ring or be connected to form a dislocation network, or move toward a surface/edge to form a threading dislocation.

All TEM images of the  $\varepsilon$ -Si layer, Ge<sub>0.2</sub>Si<sub>0.8</sub> layer and Si substrate of the sample are clean and flat, without visible defects/dislocations. In contrast, dislocations are concentrated in the Ge<sub>x</sub>Si<sub>1-x</sub> layer with the highest dislocation intensity near the Ge<sub>x</sub>Si<sub>1-x</sub>/C-Si interface. The dislocations first arose in the beginning of the Ge<sub>x</sub>Si<sub>1-x</sub> epitaxial growth. With an increase in thickness during the epitaxial growth, some (existing and/or new born) dislocation lines connected with others, to form dislocation networks, and remained within the crystal. Some connected and became dislocation in the Ge<sub>x</sub>Si<sub>1-x</sub>/C-Si interface of dislocation in the Ge<sub>x</sub>Si<sub>1-x</sub>/C-Si interface, the average lattice quality of monocrystalline silicon deteriorated rapidly in the beginning of the graded buffer layer, which leads the FWHM of the Raman peak to broaden (as shown in Figure 6). With the gradual reduction in dislocations, the lattice quality improved.

Each type of dislocation helps to relax internal/residual stresses that are induced by the change in lattice dimension [25,26]. However, the threading dislocation is harmful to the structural reliability and performance of semiconductor devices, so it is necessary to avoid the appearance of threading dislocations. By using the manufacturing process to produce a strained silicon, dislocation

lines mainly form rings or networks inside the crystal structure, which avoids the generation of threading dislocations.

#### 3.3. Distribution of Residual Stress with Depth

The residual stress inside the samples was analyzed, using MRS, which was regarded as an effective method for residual stress analysis of any semiconductor material, such as silicon. The basic Raman-mechanical model for crystal silicon (C-Si) follows the secular equation of the lattice dynamics [21,27]. According to existing investigations, the Raman-mechanical relationship of C-Si depends on the crystal plane and the stress state for measurement. In this work, the UV Raman tests were performed on the surface of the strained-silicon film, whose crystal plane was (100). Figure 9a shows a typical spectrum from the UV Raman tests. The 325-nm UV laser can only penetrate the crystal silicon to a depth of ~10 nm, which is similar to the thickness of the strained silicon. Hence, the wavenumber detected by the 325-nm laser achieves a mean wavenumber of 511.98 cm<sup>-1</sup> based on results from 66 different sampling spots, which originated solely from the  $\varepsilon$ -Si layer and were characterized by the average strain and stress of this layer. The residual stress in the  $\varepsilon$ -Si layer can be treated as a biaxial stress state, with stress components  $\sigma_x = \sigma_y = \sigma$  and strain components  $\varepsilon_x = \varepsilon_y = \varepsilon$ , where *x* and *y* denote the in-plane directions. According to the Raman-mechanical relationship of C-Si with the (100) crystal plane [28],

$$\varepsilon = -1.204 \times 10^{-3} \Delta w$$
  

$$\sigma = -217.28 \Delta w \text{ (MPa)}$$
(1)

According to Equation (1),  $\varepsilon = -1.204 \times 10^{-3} \times (511.98-519.78) = 9.39\% \approx 1\%$  and  $\sigma = -217.28 \times (511.98-519.78) \approx 1.69$  GPa, where wavenumber 519.78 cm<sup>-1</sup> was calibrated on a silicon wafer of the same production batch with the  $\varepsilon$ -Si/GeSi/C-Si samples and on the same day of the measurements above. The results showed that the real in-plane strain in the  $\varepsilon$ -Si layer was ~1%.



Figure 9. Raman spectra of (a) strained silicon (surface) and (b) germanium-silicon (cross-section).

For the Raman measurement on the cross-section with a (110) crystal plane, the applicable Raman-mechanical relationship of (110) C-Si was given as Equation (2), where the residual stress is treated as a uniaxial stress state, because the normal component on the cross-section surface (denoted as  $\sigma_x$ ) was released during the sample preparation,  $\sigma_x = 0$ ,  $\sigma_y = \sigma$  and strain components  $\varepsilon_x = \varepsilon_y = \varepsilon$ , where *x* and *y* denote the in-plane directions [28].

$$\varepsilon = -2.58 \times 10^{-3} \Delta w$$
  

$$\sigma = -434.58 \Delta w \text{ (MPa)}$$
(2)

The cross-sectional stress analyses for the Raman data from the sampling spot inside the C-Si layer can use the Raman-mechanical relationship shown as Equation (2), whereas those from the GeSi

and a Ge-Ge band at ~300 cm<sup>-1</sup>. If elemental Si dominated the Ge-Si alloy, just like the buffer layers in this work, the Raman peak of the Si-Si band would be strong and sufficiently narrow (shown in Figures 5 and 6) for mechanical measurement. However, Equation (2) is not applicable directly for the Si-Si band of germanium silicon alloy. As shown in Figure 4, the Raman wavenumber decreases sharply inside the GeSi layers, whose increment is up to 14 cm<sup>-1</sup>, denoting more than 6 GPa by using Equation (2). This is inconsistent with the reality of C-Si or GeSi strength. In fact, due to the influence of phonon confinement, the change in the wavenumber of the Si-Si band in the Raman spectrum of germanium silicon is not only induced by the strain, but also by the influence of phonon confinement. Tsang [29] presented an empirical relationship for the wavenumber  $w_{Si-Si}$ , the Ge content *x* and the strain  $\varepsilon$ , which is regarded as suitable for the germanium-silicon alloy with 0 < *x* < 0.25. Meanwhile, the elastic parameters of the germanium-silicon alloy followed the modulus-content linear model proposed by Kasper et al. [30].

$$C_{ij} = xC_{ij,\text{Ge}} + (1-x)C_{ij,\text{Si}}, \quad i, j = 1-6,$$
(3)

where  $C_{ij}$ ,  $C_{ij,Ge}$  and  $C_{ij,Si}$  were components of the elastic coefficient tensors of the Ge-Si alloy, germanium and silicon, respectively. Hence the Raman-mechanical relationship suitable for the stress analysis inside the GeSi layers on the (110) cross-section in this work is given by Equation (4) [11].

$$\sigma = (100x - 435)(\Delta w + 62x) \text{ (MPa)}.$$
(4)

The residual stress distribution along the deepness direction of this multilayer structure is given as Figure 10, where the stress inside the C-Si substrate was achieved by using Equation (2), and that inside the germanium-silicon buffer layers was achieved by using Equation (4). Figure 10 shows that there exists serious residual stress inside the buffer layers. The residual stress increases rapidly in the beginning of the epitaxial growth. The mean stress gradient during the first 2  $\mu$ m of the graded buffer layer was 777.5 MPa/ $\mu$ m. With further growth of the Ge<sub>x</sub>Si<sub>1-x</sub> layer, the residual stress stopped increasing and then began to decrease, which reduced the stress gradients and even turned them negative, until the beginning of the relaxed buffer layer. In the Ge<sub>0.2</sub>Si<sub>0.8</sub> layer, the residual stress was stable compared with that in the Ge<sub>x</sub>Si<sub>1-x</sub> layer. The mean residual stress in the relaxed buffer layer was ~1.01 GPa.



**Figure 10.** (**a**) Distribution image of residual stress on sample cross-section, (**b**) mean residual stress along the depth direction.

#### 4. Discussion

The stress distribution achieved by using Raman mapping showed that the residual stress did not increase continuously with the gradual rise in germanium content. The diameter of elemental germanium is larger than that of the silicon, and so is the lattice size of the Ge crystal. In the germanium–silicon layer prepared by reduced-pressure CVE, most silicon maintained its monocrystalline stage according to the Raman FWHM results (shown in Figure 7) and high-resolution TEM image (shown in Figure 7b,d). Therefore, some strain, and hence residual stress, was introduced into the buffer layer. The maximum residual tensile stress introduced by doping Ge in Si was 2.2 GPa, which did not induce any visible break or crack in each layer or interface, because of the birth and propagation of the dislocation, which was coincident with the GeSi alloy growth. Dislocation networks and dislocation loops helped to release the residual stress induced by the change in lattice dimension during the growth of a graded buffer layer where the germanium content increased gradually, which was positive for the reliability and stability of the  $\varepsilon$ -Si structure during its application.

In the Ge<sub>x</sub>Si<sub>1-x</sub> layer, the germanium component started from zero at the Ge<sub>x</sub>Si<sub>1-x</sub>/C-Si interface and increased with a linear gradient (shown in Figure 3a), which resulted in a gradual change in lattice constant, avoided a sudden difference in lattice constant at the interface and reduced the mismatch strain and dislocation nucleation density compared with other growth methods. When the epitaxial Ge<sub>x</sub>Si<sub>1-x</sub> layer grew accompanied by a linear increase in Ge content, the mismatch strain relaxed gradually, and slowed down the extent of dislocation lines and the formation of dislocation networks. The dislocation lines stopped when the strain energy was released, which made the dislocation networks limited inside the Ge<sub>x</sub>Si<sub>1-x</sub> layer and unable to extend into the Ge<sub>0.2</sub>Si<sub>0.8</sub> and  $\varepsilon$ -Si. Thus, the growth of the gradient-graded component can be effective for confining dislocations within the buffer layer.

As shown by the Raman results, the residual stress remained stable and the lattice quality improved inside the relaxed buffer layer. No obvious dislocation existed in this layer in the TEM images. These results proved that it was helpful to reduce the negative effect of germanium, and to improve the liability and performance of the microdevice based on strained-silicon technology.

## 5. Conclusions

In this work, multiscale structural and mechanical properties of a  $\varepsilon$ -Si/GeSi/C-Si material were investigated by several experimental methods. Based on sample cross-sections, the components and thickness of each layer were measured by EDS and SEM and the crystal and lattice characteristics and quality were observed by TEM and MRS. The residual stress distribution inside this multilayer material was achieved through micro-Raman mapping. The experimental results showed that the residual stress introduced during the heteroepitaxial growth was effectively controlled inside the buffer layers within a safe magnitude. The twin (graded plus relaxed) GeSi alloy structure provided a stable and reliable buffer for the manufacture of a strained-silicon film/ribbon in microelectronic devices.

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