# Palladium (III) Fluoride Bulk and PdF<sub>3</sub>/Ga<sub>2</sub>O<sub>3</sub>/PdF<sub>3</sub> Magnetic Tunnel Junction: Multiple Spin-Gapless Semiconducting, Perfect Spin Filtering, and High Tunnel Magnetoresistance

### **Computational Methods**

In our study, the structural, electronic, and magnetic properties for rhombohedral-type PdF<sub>3</sub> bulk were calculated based on the density functional theory (DFT), as implemented in the Nanodcal package [1]. The Perdew-Burke-Ernzerhof (PBE) exchange-correlation functional was applied [2] for the parameterization of general gradient approximation (GGA) [3]. The DFT+U function was considered to examine the electronic structures of PdF<sub>3</sub>.

The spin-transport calculations were performed by the nanodcal package, which used the DFT combined with the non-equilibrium Green's Function method [4]. In this calculation, the Monkhorst-Pack k-meshes for the electrons and the centers were 9×9×100 and 9×9×1, respectively. The self-consistent calculation is limited to 10<sup>-5</sup> Hartree tolerance. All these parameters were evaluated and found to be sufficient to obtain accurate results. The spin-dependent current is calculated by the Landauer-Buttiker formula:

$$I^{\uparrow(\downarrow)} = \frac{e}{h} \int_{-\infty}^{\infty} \left\{ T^{\uparrow(\downarrow)}(E, V_b) \left[ f_L(E - \mu_L) - f_R(E - \mu_R) \right] \right\} dE$$

where  $f_{L/R}(E - \mu_{L/R})$  is the Fermi-Dirac distribution for the left (right) electrode, and  $\mu_{L/R} = E_F \pm eV/2$  is the corresponding electrochemical potential.  $T^{\uparrow(\downarrow)}(E, V_b)$  is the spin-dependent transmission coefficient:

$$T^{\uparrow(\downarrow)}(E,V_b) = Tr \Big[ \Gamma_L G^R \Gamma_R G^A \Big]^{\uparrow(\downarrow)},$$

where  $G^{R(A)}$  is the retarded (advanced) Green's function of the central region, and  $\Gamma_{L(R)}$  is the coupling matrix of the left (right) electrode.

## Models of bulk Palladium (III) fluoride and PdF<sub>3</sub>/Ga<sub>2</sub>O<sub>3</sub>/PdF<sub>3</sub> MTJ

The crystal structure of the rhombohedral-type PdF<sub>3</sub> (space group:  $R\overline{3}c$ , No. 167, ICSD ID: 16675), as shown in Figure S1. The crystal structure has been fully relaxed and detailed structure optimization methods can be found in The Materials Project database [5]. The obtained lattice constants are a = b = c = 5.618 Å,  $\alpha = \beta = \gamma = 54.956^{\circ}$ , respectively. The lattice parameters obtained by theory are consistent with the experimental values [6]. According to the database [5], one can see that the calculated formation energy is -1.592 eV and negative value of formation energy shows the possible stability of bulk Palladium (III) fluoride. The magnetic ground state of PdF<sub>3</sub> is ferromagnetic with a magnetic moment of 0.779 µB.



Figure S1. (a) Front and (b) side views of lattice structure of rhombohedral-type PdF<sub>3</sub> bulk.

In our work, the MTJ device model is periodic along the x- and y-axes, while the transport direction is along the z-axis. The crystals of Ga<sub>2</sub>O<sub>3</sub> can be found here [7]. Ga<sub>2</sub>O<sub>3</sub> with rhombohedral-type is a semiconductor and no virtual frequency in the phonon spectrum of Ga<sub>2</sub>O<sub>3</sub> guarantee the stability of this material. The optimized distance between PdF<sub>3</sub> layer with F-terminated interface and O layer was found to be 2.61 Å. The lattice mismatch between PdF<sub>3</sub> and Ga<sub>2</sub>O<sub>3</sub> is 3.72%.



Figure S2. Geometric structure of PdF<sub>3</sub>-Ga<sub>2</sub>O<sub>3</sub>-PdF<sub>3</sub> magnetic tunnel junction (MTJ).



**Figure S3.** Band structures of PdF<sub>3</sub> bulk calculated based on the PBE+U method. Here, we assessed the influence brought by Hubbard U from 1 eV to 4 eV for Pd-4*d* orbits.

# Band structure under MBJ potential

Since the GGA potential sometimes underestimates potential band gaps, here we check the band structure of  $PdF_3$  by using a combination of modified Becke-Johnson exchange potential with GGA (MBJGGA) [8]. As shown in Figure S4, we can observe that the novel multiple linear-type spin-gapless semiconducting band structures in  $PdF_3$  retain under the MBJ calculation.



Figure S4. Band structure of PdF3 bulk under MBJGGA calculation within VASP software.



Figure S5. Band structure of  $PdF_3$  bulk calculated based on the PBE functional with consideration of the SOC effect.



Figure S6. Band structures of  $PdF_3$  bulk under different pressures. These results were calculated based on the PBE functional.



**Figure S7.** (a) Equilibrium-state transmission spectrum in parallel configuration for PdF<sub>3</sub>-based MTJ. (b) Equilibrium-state transmission spectrum in anti-parallel configuration for PdF<sub>3</sub>-based MTJ.



**Figure S8.** Non-equilibrium transmission spectrum versus electron energy at a fixed bias voltage for PdF<sub>3</sub>-based MTJ.

### References

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