



# Article Wavelength-Stable Metal Grating Distributed Feedback Quantum Cascade Laser Emitting at $\lambda \sim 7.2 \ \mu m$

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Abstract: In this research, we demonstrate a wavelength-stable continuous wave (CW) distributed feedback (DFB) quantum cascade laser (QCL) emitting at 7.2  $\mu$ m using a surface metal grating approach without epitaxial regrowth. The deep metal grating provides an appropriate DFB coupling coefficient and enhanced thermal extraction, resulting in improved lasing performance and the realization of impressive wavelength stability. Quantitatively, the temperature tuning coefficient of the single-mode emission is only 0.54 nm/°C from 20 °C to 70 °C, and the current tuning coefficient of the single-mode emission is 3.2 nm/A from 1.0 A to 1.6 A. A DFB-QCL with a 2 nm cavity length exhibits a low threshold current of 0.6 A and a power of 1.1 W with a slope efficiency of 1 W/A in the CW mode at 300 K. A single-mode operation with a side mode suppression ratio of 33 dB and a single-lobed far-field without beam steering is obtained in the working temperature range of 20–70 °C The improved wavelength stability using a deep surface metal grating approach promises simplified fabrication, which is meaningful for the commercial applications of QCLs.

Keywords: quantum cascade laser; mid-infrared; optoelectronic device; semiconductor laser

# 1. Introduction

Mid-infrared optoelectronic devices, such as lasers and detectors, are in high demand owing to applications involving the imaging and remote sensing of trace pollutants and industrial chemicals, where lasers with wavelengths in the range of 3 to 12  $\mu$ m are crucial [1–5]. A wavelength-tunable and stable quantum cascade laser (QCL) stands out as the preferred choice. The QCL, a semiconductor laser based on intersubband transitions, provides a devisable lasing wavelength across the entire mid-infrared range [6,7]. Moreover, the QCL exhibits the unique advantages of compact, high-power, and flexible emission wavelengths [8–10]. Considering several important gases, including CH<sub>4</sub>, N<sub>2</sub>H<sub>4</sub>, and SO<sub>2</sub>, which have characteristic absorption peaks at around 7  $\mu$ m, QCLs operating at around 7  $\mu$ m with single-mode, high-performance, ideal beam coherence, as well as superior wavelength stability, have become imperative demands [11–14].

The distributed feedback (DFB) laser is regarded as an effective technology for achieving single-mode lasing and ideal beam quality with a mature fabrication process [11,15]. So far, DFB QCLs with buried grating have demonstrated impressive lasing properties in the mid-infrared spectral range [16,17]. For instance, watt-level output power is achieved in continuous-wave (CW) mode at room temperature within the wavelength range of 4.6–7.0  $\mu$ m. The highest operation temperature reaches over 100 °C, and the side-mode suppression ratio (SMSR) exceeds 30 dB [18–21]. Compared with the buried gratings utilized in DFB laser devices, the surface metal gratings offer the advantages of avoiding



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**Copyright:** © 2023 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). waveguide regrowth and simplifying the manufacturing process [22–25]. Due to the benefits of a simple fabrication process and low facility requirements [26], various surface metal grating DFB lasers have been explored in recent years [27-31]. Since the investigation of the surface metal grating DFB QCL in 2006 [22], the performance of DFB QCL devices with metal gratings has shown continuous improvement [32]. A record high output power single-mode QCL with a 30 dB SMSR emitting at  $\lambda \sim 4.8 \,\mu\text{m}$  has been realized [23,24]. The progress in metal grating DFB QCL shows the potential of approaching reliable singlemode mid-infrared lasing with a brief process. However, metal grating DFB QCLs have a weakness: the unstable wavelength. Specifically, due to the thick cladding waveguide of mid-infrared QCLs, the surface metal grating is positioned far from the active region, resulting in a weak coupling coefficient ( $\kappa$ ) for metal grating QCLs. Consequently, general metal grating DFB QCLs require a long cavity length to achieve sufficient coupling strength  $\kappa L$  [24,33]. However, the long cavity length of the DFB laser not only induces spectral distortion, such as photon hole burning, but also enhances thermal accumulation in the laser cavity, leading to diminished performance and a significant wavelength shift [32,34]. At present, for surface metal grating DFB QCLs, the reported temperature tuning coefficient of single-mode emissions could reach  $0.9 \text{ nm}/^{\circ}\text{C}$ , and the current tuning coefficient of single-mode emission is 11.8 nm/A in CW mode [24,25]. Nevertheless, in the practical operation of mid-infrared sensing, avoiding wavelength shifts is crucial. This is because a wavelength shift can cause a mismatch between the accurate lasing wavelength and the fingerprint-like absorption peaks of target gases, leading to distorted sensing results [35]. Consequently, the reported metal grating DFB QCLs fail to meet the demands of commercial trace gas sensing. Hence, improving the wavelength stability of metal grating DFB QCLs is imperative.

In this work, deep surface grating with thick capping metal is put forward to increase the coupling coefficient and improve thermal extraction. The grating depth is etched to approximately 500 nm, and the innovative fabrication process of the deep surface grating achieves a coupling strength of 2.6 with a 2 mm device cavity length. Herein, we present a high-performance metal grating DFB QCL emitting at  $\lambda \sim 7.2 \mu m$  at room temperature in CW mode. The watt-level output power device shows improved single-mode stability within the temperature range of 20–70 °C and a current range of 1.0–1.8 A. The realizations of wavelength-stable metal grating DFB QCL with outstanding lasing properties open up new possibilities for mid-infrared lasers.

# 2. Material Growth and Device Fabrication

#### 2.1. Material Epitaxy of QCL

The QCL wafer utilized in this work is grown on an n-InP (Si-doped,  $1 \times 10^{17}$  cm<sup>-3</sup>) substrate, with the active region, serving as the core of the whole laser structure, sandwiched between two n-doping InP cladding layers. The active structure follows a bound-tocontinuum design, comprising 50 repetitions of periodic multi-epitaxial layer groups grown by molecular beam epitaxy (MBE) with strain-balanced In<sub>0.43</sub>Al<sub>0.57</sub>As/In<sub>0.6</sub>Ga<sub>0.4</sub>As material compositions (see Figure 1a). The layer sequence for one period, in angstroms, starting from one In<sub>0.43</sub>Al<sub>0.57</sub>As barrier, is **18.1**/32.3/**17.2**/27.9/**19.3**/<u>27.1</u>/<u>27.7</u>/25.8/**39.6**/17.1/**9.2**/50.6/ 8.9/47.1/10.4/38.8, where the In<sub>0.43</sub>Al<sub>0.57</sub>As barrier layers are in bold, In<sub>0.6</sub>Ga<sub>0.4</sub>As well layers are in regular font, and two n-doped layers ( $1.5 \times 10^{17} \text{ cm}^{-3}$ ) are underlined. Based on this active region structure design, the calculated wavefunctions of relevant energy levels in the conduction band are plotted in Figure 1b. The QCL operating mechanism is based on an optical transition design, where the photon is generated from the carrier transition between the upper energy level (red) and the lower energy level (blue, and Cambridge blue). According to the simulation result, the transition energy is calculated as 178 meV, corresponding to a lasing wavelength of 7.2 µm. During the MBE growth of the active region, optimized growth parameters were employed. InAlAs tend to grow better at high substrate temperatures (510-535 °C) and a low V/III ratio, while InGaAs grow better at low substrate temperatures (480–515 °C) and a high V/III ratio. A compromise

was reached by using the same substrate temperature and arsenic overpressure for both materials. Thus, the substrate temperature of 510 °C and the V/III ratio of 15–19 in growth were selected. The growth rate used was approximately 0.7  $\mu$ m/h and 0.98  $\mu$ m/h for In<sub>0.6</sub>Ga<sub>0.4</sub>As and In<sub>0.43</sub>Al<sub>0.57</sub>As, respectively. The temperature stability of group III effusion cells should be calibrated better than 0.2 °C. The group III flux stability was adjusted to better than 1.5%. Due to the sensitivity of the InGaAs layer to the substrate temperature, the operation stability of the substrate heater at the growth temperature was confirmed to be better than 0.3 °C.



**Figure 1.** (a) Structure sequence of the QCL active region. (b) Schematic conduction band diagram of the QCL active region under an electric field of 65 kV/cm.

The growth of thick InP performed run to run every day using solid source MBE with a valved cracking phosphorus effusion cell will lead to the inevitable residual of white phosphorus in the growth chamber. Because of the limited cracking efficiency of  $P_4$  to  $P_2$ , the unintentional release of  $P_4$  is unavoidable. The naturally enhanced vapor pressure of white phosphorus may introduce potential "pollution" during the growth of the InAlAs/InGaAs active region. To mitigate this issue, the InP cladding layers were grown using low-pressure (100 mbar) metal-organic chemical vapor deposition (MOCVD). The precursors were trimethyl indium and purified phosphine. A growth temperature of 600–650 °C and a growth rate of 0.3–0.8 nm/s were used. The lower cladding consists of a 4.3 µm-thick n-InP layer [graded Si doping,  $(3 - 0.3) \times 10^{17} \text{ cm}^{-3}$ ]. The upper cladding is composed of a 4 µm-thick n-InP layer [graded Si doping,  $(5 \times 10^{18} \text{ cm}^{-3})$ .

#### 2.2. Design and Fabrication of Surface Metal Grating

In this work, a standard periodic distributed feedback (DFB) type with anti-reflection (AR) coating deposited on the front facet was employed to achieve single-mode lasing. Due to the need for a systematic grating simulation to obtain an optimal optical field distribution and a suitable coupling coefficient value, a surface–plasmon coupling mechanism-based simulation was explored for the grating structure design. The preliminary fabrication process (holographic exposure and inductively coupled plasma dry etching) yielded an ideal grating shape with a 35% duty cycle and a 1050 nm period width. These parameters were retained in the simulation process, with attention focused on enhancing the coupling coefficient by adjusting the grating depth and capping the metal thickness. Under the condition of infinite capping metal thickness, the relationship between the coupling coefficient and grating depth is displayed in Figure 2a. With the rise of grating depth, the

coupling coefficient steadily increases when grating is deeper than 200 nm. This shows that increasing the grating depth is an effective method to enhance the coupling coefficient. Therefore, a grating depth of 500 nm was determined, as it was capable of providing a high coupling coefficient. Next, the influence of the capping metal thickness was considered, and the simulated curve of the coupling coefficient value  $\kappa$  versus capping metal thickness is presented in Figure 2b. The coupling coefficient  $\kappa$  reaches its peak when the capping metal thickness is increased to 350 nm, and then decreases with further thickness increase. Finally, the coupling coefficient  $\kappa$  gradually stabilizes around 6.5 cm<sup>-1</sup> once the metal thickness exceeds 550 nm. Given that the capping metal used in this work is deposited by micron-order gold (Aurum, Au) with a metal thickness exceeding 550 nm, the expected value of coupling coefficient  $\kappa$  in this research is 6.5 cm<sup>-1</sup>. Meanwhile, as the cavity length (L) of the laser devices is cleaved to 2 mm, and the back facet of the devices is coated with high-reflectivity (HR) coating, the simulated value of the designed coupling strength  $\kappa$ L is calculated as 2.6.



**Figure 2.** (a) The simulation result of the relationship between the grating depth and coupling coefficient under the condition of infinite capping metal. (b) The curve of the simulated coupling coefficient versus capping metal thickness. (c) The scanning electron microscope (SEM) picture of the grating after etching (before metal deposition). Information about the grating width, period length, and grating depth is shown in this picture.

Afterward, holographic lithography combined with standard inductively coupled plasma (ICP) dry etching was employed to produce the DFB grating on the highly doped capping InGaAs layer situated on top of the wafer. ICP dry etching of high-depth gratings poses challenges due to the instability of the gas etching rate with increasing depth, making it difficult to guarantee the morphology of the deep grating. To achieve deep grating with ideal morphology, the fabrication processes of holographic lithography and ICP dry etching were optimized. In holographic lithography, a high-precision and highly photosensitive coupling photoresist (S1805) was adopted to achieve the flatness of the  $SiO_2$  hard mask, reducing the serrated morphology of the grating sidewall after dry etching of InGaAs. During the ICP dry etching process, the etching temperature was elevated to 240 °C; the aim was to use high-temperature dry etching to prevent the accumulation of etching products on the surface of the grating. Simultaneously, optimized menus and parameters were applied in the ICP etching process, with the etching gas adopting the chlorine gas to argon gas ratio of 1.5:40, and the power settings for the ICP RF and Table RF corresponding to 300 W and 100 W, respectively. The optimized ICP process aims to ensure the etched grating sidewall is vertical and the bottom is flat.

The actual morphology of the DFB grating can be observed through scanning electron microscope (SEM) pictures, as shown in Figure 2c. The period width and etching depth of the grating are measured at 1047.97 nm and 509.79 nm, respectively, resulting in a calculated grating duty cycle of approximately 36%. The actual shape of the DFB grating after processing closely aligns with the expected structure design, featuring a 1050 nm period width, 500 nm grating depth, and 35% duty cycle. Therefore, the actual coupling strength  $\kappa$ L is expected to be close to the simulated value of 2.6, indicating a slightly overcoupled state. As a strong coupling strength contributes to confining the laser to the center of devices and facilitates the production of single-mode operation, the structure of the grating and the design of the corresponding coupling strength are anticipated to achieve a high-performance single-mode mid-infrared lasing operation.

## 2.3. Fabrication of QCL Devices

After finishing the fabrication of metal grating on the surface, the wafer was transformed into metal grating DFB QCL devices with a double-channel ridge waveguide. The double-channel ridge waveguide, with a 10 µm core active region width, was created through ICP dry etching. Subsequently, wet etching was employed to refine the roughness induced by the uneven dry etching hard mask formed by the DFB grating on the top of the waveguide, resulting in a stripe-like rough sidewall. Following the deposition of the standard insulation layer (SiO<sub>2</sub>) and window opening, 30 nm Ti and 120 nm Au were deposited and annealed to establish ohmic contact metallization. A micron-order thick Au metal layer was then plated on the surface to fill and flatten the DFB grating and channel. The fabricated sample was cleaved into independent devices with a 2 mm cavity length and optical coating. The back facet of the devices received a high-reflecting (HR) coating consisting of 165 nm  $Al_2O_3/20$  nm Ti/50 nm Au/25 nm Ti/120 nm  $Al_2O_3$ , while the front facet was coated with an anti-reflective (AR) coating comprising periodic 135 nm  $Al_2O_3/45$  nm Ge. The measured reflectivity values for HR and AR coatings are 98.8% and 0.3%, respectively. Finally, the lasers were flip-chip mounted epi-side down on diamond heat sinks with indium solder, and then subsequently bonded to copper heat sinks. The fabrication process of QCL devices in this research adheres to the standard fabrication and packaging procedures for semiconductor optoelectronic devices, ensuring repeatability and verifiability.

# 3. Results and Discussion

#### 3.1. Electroluminescence Spectra of QCL

In order to acquire the internal energy level information from the epitaxial laser material, the majority of the grown QCL wafer was processed into surface metal grating DFB QCL devices. Additionally, a portion of the sample was fabricated into a Fabry–Perot (FP) laser to obtain electroluminescence (EL) spectra. The spectrum of the  $20-\mu$ m-wide/3mm-long FP device was subsequently measured using Fourier-transform infrared (FTIR) and a mercury cadmium telluride (MCT) detector under room temperature, using different current levels in the pulsed injection mode (5  $\mu s$  pulse width, 20 kHz repetition rate, 10% duty cycle). As shown in Figure 3, with the resolution of 0.5 cm<sup>-1</sup> in FTIR rapid scan mode, the EL results of the QCL device were obtained with the progressively pumped current ranging from 0.2 A to 1 A. The main emission peak consistently appeared at around 7.2 µm, which is consistent with the calculated value derived from band structure simulation. A weaker peak at around 6  $\mu$ m arose from the carrier transmission from a higher parasitic energy level to a low energy level, as indicated by the orange color in the simulation result (Figure 1b). The EL peak exhibited a behavior akin to being "currentpinned," wherein a fivefold increase in the driving current resulted in minimal change in the peak. Consequently, it is conceivable that the lasing spectrum of the QCL device is also insensitive to changes in the driving current. This characteristic suggests a promising avenue for achieving substantial wavelength stability in these QCL devices.



**Figure 3.** The electroluminescence spectra of a 20-µm-wide and 3-mm-long FP device operation under pulsed injection current, 5 µs in width, at a repetition rate of 20 kHz.

# 3.2. Lasing Performance of Metal Grating DFB QCL

Upon completing the wafer epitaxy and laser device fabrication, measurements were conducted to comprehensively evaluate the performance of QCL devices featuring a 2 mm cavity length and a 10  $\mu$ m ridge width. The typical light current–voltage (LIV) results, obtained under room temperature conditions (300 K) with the continuous wave (CW) mode injection current, are presented in Figure 4a. The laser device shows excellent operational characteristics, boasting a 600 mA threshold current and a high slope efficiency of 1 W/A. As the injection current increases, the output power could surpass 1.1 W.



**Figure 4.** (a) Light current–voltage (LIV) experiment of a 10- $\mu$ m-wide/2-mm-long sample measured under 300 K room temperature and CW mode. (b) Temperature-dependent LIV result measured in the temperature range from 20 °C to 60 °C with 5 °C steps. The inset figure shows the relationship between slope efficiency and operation temperature.

Subsequently, employing a thermoelectric cooler (TEC) for temperature control, the temperature-dependent output power of the QCL was measured across the range of 20 °C to 60 °C, with the results shown in Figure 4b. With the increasing operating temperature, the slope efficiency and threshold current of the laser exhibit slight changes under continuous

wave driving, while the threshold current increases marginally from 600 mA to 665 mA, and the calculated slope efficiency decreases to approximately 0.94 W/A from 1 W/A. Thus, the QCL device in this research demonstrates not only an impressive slope efficiency but also remarkable thermal stability of output power.

#### 3.3. Wavelength-Stable Metal Grating DFB QCL

As mentioned in the preceding section of this paper, wavelength stability stands as a crucial criterion for DFB laser devices, and the imperative need for the low wavelength shift in DFB QCLs becomes apparent for practical applications. Therefore, the lasing spectrum of the devices was measured by Fourier-transform infrared (FTIR) and a deuterated triglycine sulfate (DGTS) detector, with a resolution of 0.06 cm<sup>-1</sup>. The wavelength shift was then calculated from the obtained spectrum results.

The spectrum results, representing the variation in injection current, are illustrated in Figure 5a. In this depiction, the 2 mm  $\times$  10 µm QCL device is driven by the CW mode current ranging from 1.0 A to 1.8 A under room temperature. The peak of the emission spectrum only experiences a minute increase of 0.5 cm<sup>-1</sup> in wavenumber, transitioning from 1392.7 cm<sup>-1</sup> (1.0 A) to 1393.2 cm<sup>-1</sup> (1.8 A). Consequently, the calculated wavelength shift is only 2.58 nm. The alterations in the side-mode suppression ratio (SMSR) and lasing wavelength under varying injection currents are plotted in Figure 5b. As the injection current increases from 1.0 A to 1.6 A, the SMSR exhibits an ascending trend, reaching its peak value of 33 dB before gradually declining as the current continues to increase. Moreover, the wavelength shift is clearly observed in Figure 5b, with the shift ratio under current modulation measuring only around 3.2 nm/A in the range of 1.0 A to 1.8 A, and notably, no mode hopping is observed.



**Figure 5.** (a) The spectrum under various injection currents, measured with the CW mode injection current and room temperature. The current is set in the range of 1.0 A to 1.8 A with 100 mA steps. (b) The curve of the wavelength (wavenumber) versus the injection current (blue), and the curve of the SMSR under different injection currents (red).

The subsequent measurement in this work aimed to assess the lasing wavelength stability of the DFB laser device under varying operational temperatures. The temperature range was set from 20 °C to 70 °C with a 10 °C temperature step, and the measurement was conducted at a constant injection current of 1.6 A in continuous wave (CW) mode. The temperature-dependent lasing spectrum is shown in Figure 6a, revealing the conventional red shift of the operational wavelength with increasing temperature. No observed mode hopping occurred during this process. In detail, at a 20 °C operation temperature, the lasing wavelength of the QCL device was 7.17  $\mu$ m (1394.7 cm<sup>-1</sup> wavenumber), and it changed to 7.197  $\mu$ m (1389.5 cm<sup>-1</sup> wavenumber) when the temperature increased to 70 °C.



Throughout this measurement, the normalized intensity of the side mode remained below  $10^{-2}$ , signifying that the side-mode suppression ratio (SMSR) remained higher than 20 dB.

**Figure 6.** (a) Temperature-dependent spectrum measurement of metal grating DFB QCL in an operation temperature range from 20 °C to 70 °C with 5 °C steps. The spectra measured under room temperature and CW mode are presented inset. (b) The curve of the wavelength (wavenumber) versus operation temperature (blue), and the curve of the SMSR under varying temperatures, from 20 °C to 70 °C with 5 °C steps (red).

In Figure 6b, the spectrum measured under room temperature (300 K) and the CW mode is presented inset as an individual figure, demonstrating the SMSR exceeding 30 dB (33 dB) under the given measurement conditions. Further analysis of the temperature-dependent results is presented in Figure 6b. The QCL lasing wavelength uniformly increased with the increasing operational temperature, achieving a remarkable wavelength shift ratio of only 0.54 nm/°C in the temperature range of 20 °C to 70 °C. Additionally, despite the decreasing trend in the SMSR value with the increasing temperature, the DFB laser continued to operate with a 26 dB SMSR under 70 °C. The temperature-dependent measurement manifested the impressive high-temperature stability of the DFB QCL device reported in this work.

In addition to lasing properties, such as output power and wavelength, the quality of the light beam is a meaningful evaluation of laser devices, influencing beam coherence and transmission loss in practical applications. The beam quality of the QCL device was assessed through far-field profile measurements along the ridge-width direction in this work. Figure 7 shows the divergence angle measured by the far-field profile with varying continuous wave (CW) injection currents under room temperature.

In the current range of 1.0 A to 1.6 A, the increase in current does not alter the singlelobe distribution of the divergence angle. The full-width half-maximum (FWHM) at 1.0 A is 28°, increasing to 36° with a 1.6 A injection current. The persistence of a single-lobe distribution and the low FWHM indicate the fundamental transverse mode and wellcoupled light beam, standing for the high beam quality of the reported DFB QCL. The pyroelectric camera image of the light beam spot is attached in the top-left corner of Figure 7, showing a lasing spot shape close to a circle with an aspect ratio of around 5:4. This spot picture confirms that the lasing mode of DFB lasers is the fundamental transverse mode, emphasizing excellent light beam quality.



**Figure 7.** The divergence angle measured by the far-field profile along the ridge-width direction with different injection currents, from 1 A to 1.6 A. The attached figure shows the light beam spot of the QCL operation after collimating. The spot is measured by a pyroelectric camera placed 60 cm away from the laser. The results in Figure 7 are measured under room temperature and c.w. mode.

#### 3.4. Discussion

The foundation of this research is the realization of the deep surface metal grating structure, and several challenges, such as the rough etched morphology, were addressed through the optimization of the grating fabrication process. The morphology of the SiO<sub>2</sub> hard mask after holographic exposure and reactive ion etching was analyzed and improved to minimize the serrated shape of the fabricated grating. Furthermore, the optimization of gas flow and RF power in the ICP dry etching process was employed to achieve suitable roughness on the grating sidewall, even with a dry depth as high as 500 nm. In this work, the realization of a 500 nm depth surface metal grating structure resulted in a coupling coefficient as high as  $6.5 \text{ cm}^{-1}$ . This achievement represents a significant improvement compared to previous works, where the reported results only reached a 200 nm depth and a coupling coefficient of  $3.1 \text{ cm}^{-1}$  [23,25]. The increased grating depth effectively enhances the coupling coefficient, offering more flexibility in the design of distributed feedback (DFB) structures and contributing to improved performance in DFB laser devices.

In the research domain of surface metal grating distributed feedback (DFB), a high coupling coefficient signifies the attainment of high surface coupling efficiency for the light field, leading to low-threshold single-mode lasing [36]. DFB lasers with shorter cavity lengths can achieve the necessary coupling intensity for single-mode light emission, mitigating the influences of thermal accumulation and space hole burning associated with longer cavity lengths [37]. Consequently, a high coupling coefficient is crucial for minimizing wavelength drift due to thermal effects and realizing high single-mode characteristics in DFB lasers. Therefore, the deep surface metal grating DFB quantum cascade laser (QCL) is anticipated to yield high-performance DFB QCL with excellent wavelength temperature stability.

The measurement results in this study align with theoretical expectations. As demonstrated by the experimental results, the DFB QCL device fabricated with a deep metal grating structure shows impressive lasing performance with a well-defined light beam. This outcome features the reasonable design of the DFB structure and the feasibility of DFB QCL with a surface metal grating. Furthermore, it is noteworthy that the DFB QCL in this research exhibits outstanding wavelength stability, characterized by an extremely low wavelength shift ratio and the absence of mode hopping under both current and temperature variations. The achievement of wavelength-stable metal grating DFB QCL is attributed to the utilization of a deep metal grating design and the successful implementation of the challenging fabrication process. With an increased coupling coefficient, the contractible cavity effectively alleviates the thermal challenges of metal grating DFB QCL.

## 4. Conclusions

In conclusion, this work reports on a DFB QCL that emits at around 7.2  $\mu$ m, demonstrating a record wavelength stability. The wavelength shift ratio for the 2-mm-long/10- $\mu$ m width device is remarkably low, measuring only 3.2 nm/A under continuous wave (CW) injection current turning and 0.54 nm/°C during temperature variations. The device exhibits stable single-mode operation with a maintained signal-to-noise ratio (SMSR) of over 20 dB, without any occurrence of mode hopping, across the injection current range of 1.0 A to 1.6 A and temperature range of 20 °C to 70 °C. Furthermore, the DFB QCL achieves an output power of 1.1 W with a high slope efficiency of 1 W/A under room temperature and CW mode, accompanied by an SMSR value for the output lasing reaching 33 dB. The realization of a wavelength-stable DFB QCL with such high performance holds promise for applications in mid-infrared sensing, where a stable light source is crucial for precisely matching the absorption spectra of target substances.

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