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Honey-like Odor Meets Single-Ion Magnet: Synthesis, Crystal Structure, and Magnetism of Cobalt(II) Complex with Aromatic Trans-Cinnamic Acid

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Abstract: The hexacoordinate Co(II) complex $[\text{Co}(\text{neo})_2(\text{cin})][\text{BPh}_4] \cdot 1/2 \text{ Me}_2\text{CO}$ ($1 \cdot 1/2 \text{ Me}_2\text{CO}$) containing *trans*-cinnamic acid (Hcin) and neocuproine (neo) was prepared. The compound $1 \cdot 1/2 \text{ Me}_2\text{CO}$ was characterized via single-crystal X-ray analysis, FT-IR spectroscopy, and magnetic measurements. The coordination polyhedron of the complex cation adopts a deformed octahedron shape, and cinnamate exhibits a bidentate mode of coordination, which is unusual for mononuclear Co(II) cinnamate complexes. The analysis of DC magnetic measurements with zero-field splitting (ZFS) spin Hamiltonian revealed large magnetic anisotropy defined by the axial ZFS parameter $D = +53.2 \text{ cm}^{-1}$. AC susceptibility measurements revealed the slow relaxation of magnetization under the applied field; thus, $1 \cdot 1/2 \text{ Me}_2\text{CO}$ behaves as a field-induced single-molecule magnet. The analysis of magnetic properties was also supported by CASSCF/NEVPT2 calculations.

Keywords: cobalt(II); complex; single-molecule magnet; cinnamic acid



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1. Introduction

Single-ion magnets (SIMs), a subgroup of single-molecule magnets (SMMs) where only one metallic center is responsible for the slow relaxation of magnetization, have become increasingly attractive as more and more newly discovered SIMs have exhibited higher relaxation barriers and blocking temperatures, culminating in a series of dysprosocenium complexes [1,2] with the highest blocking temperature yet achieved of $T_B = 80 \text{ K}$ and $U_{\text{eff}} = 1541 \text{ cm}^{-1}$ for the $[(\text{Cp}^{\text{iPr5}})\text{Dy}(\text{Cp}^*)][\text{B}(\text{C}_6\text{F}_5)_4]$ complex ($\text{Cp}^{\text{iPr5}} = \text{penta(isopropyl)cyclopentadienyl}$, $\text{Cp}^* = \text{pentamethylcyclopentadienyl}$) [3].

However, as these compounds are usually highly sensitive to air and moisture, more stable alternatives, albeit generally exhibiting lower barriers, are still of much interest, such as the $[\text{Dy}(\text{bbpen})\text{X}]$ derivatives ($\text{bbpen} = N,N'$ -bis(2-hydroxybenzyl)- N,N' -bis(2-pyridylmethyl)ethylenediamine) [4,5] for Ln(III) systems or the $(\text{HNEt}_3)_2[\text{Co}(\text{bmsab})_2]$ complex ($\text{bmsab} = N,N'$ -1,2-phenylenebis(methanesulfonamide)) [6] for 3d metal systems. Out of the 3d metals, Co(II) systems are by far the most promising due to their large spin-orbit coupling and good predisposition for significant magnetic anisotropy; thus, a plethora of cobalt-based SIMs already exist [7–9]. Co(II) systems do not exhibit such high barriers but are generally easier to design compared to their Ln(III) counterparts as their coordination sphere is usually limited to 4–6 donor atoms, whereas Ln(III) may contain a much larger number of donor atoms, commonly between 8 and 12. This often leads to the unwanted coordination of either a solvent molecule or a bridging ligand, which might worsen the magnetic properties [10].

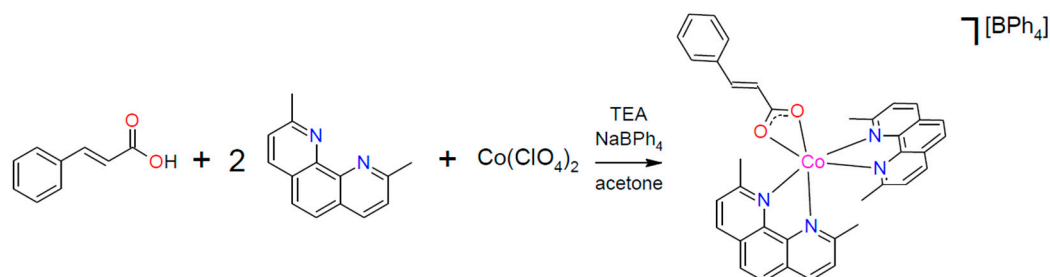
The magnetic bistability of SIMs offers a range of applications such as spintronics, quantum computing or high-density data storage [11]. Most often, switching between the spin states is performed with the external magnetic field; however, there has been much interest in the design of systems with alternative ways of effecting the form of switching

that may, for example, introduce on/off states of the slow relaxation of magnetization, provide ultrafast changes in magnetization or eliminate the impracticalities associated with external magnetic field switching focused to a scale of single molecules [12,13]. One method of design is based on the introduction of photoswitchable molecules into the metal complex structure, which would induce a change in the magnetic properties [14] upon irradiation, analogously to what has already been observed for ligand-driven light-induced spin change complexes [15,16].

Photoswitchable complexes containing azo (N=N) [17,18] or olefin (C=C) [19,20] moieties have been well known for their *cis/trans*-photoisomerization property [21] for some time, although not many have been studied for their photomagnetic properties. Their photoisomerization is generally only observed in solution as crystal packing often sterically hinders the switching action in the solid phase; however, some research groups have shown that this can be overcome, for example, by dissolving and setting the sample in a polymer matrix [22] or by utilizing layered polyoxometalate films [23].

The naturally occurring and thermodynamically more stable *trans*-isomer of cinnamic acid, well known for its honey-like odor, ref. [24], has been historically one of the first widely studied olefin photoisomerisable systems. Dating back to the early 1900s, the first irradiation studies [25,26] showed the possibility of the isolation of the separate geometric isomers, and B. K. Vaidya later reported [27] on a more robust procedure for the isolation of the *cis*-isomer by simply irradiating a methanolic solution of a *trans*-cinnamic acid derivative with UV light and subsequently separating the isomers via fractional crystallization or, alternatively, vacuum distillation [28]. A new easier methodology for the preparation of *cis*-isomers of various cinnamic acid derivatives was introduced by M. L. Salum and co-workers, which involves the irradiation of their ionic liquids in acetonitrile, during which less soluble *cis*-isomer precipitates out of the solution [29]. Compared to most of its azobenzene counterparts, *cis*-cinnamic acid shows much better thermal stability and only readily reverts to the *trans*-isomer when heated. Interestingly, only a handful of mononuclear Co(II) complexes containing cinnamic acid as a ligand have been prepared, and in all cases, cinnamic acid coordinates monodentately with the single carboxylate oxygen to the metal center [30–33].

Based on the aforementioned insights, we decided to utilize *trans*-cinnamic acid (Hcin) as a co-ligand for the preparation of a new single-ion magnet with a photoswitchable moiety. The cobalt(II) bis(neocuproine) complex seemed to be an excellent initial building block, as a more recent paper by J. Vallejo and co-workers utilized this exact system with the benzoic acid co-ligand to yield a complex exhibiting a slow relaxation of magnetization [34]. Moreover, the neocuproine ligand was beneficial for the preparation of other Co(II) SMMs in the past [35–37]. We were successful in the preparation of the title compound $[\text{Co}(\text{neo})_2(\text{cin})][\text{BPh}_4] \cdot 1/2\text{Me}_2\text{CO}$ (**1**·1/2Me₂CO) (Scheme 1), a novel single-ion magnet with a light-switchable moiety, and herein report on its synthesis, crystal structure, and magnetic data measurements.



Scheme 1. Synthetic procedure for preparation of title compound **1**·1/2Me₂CO.

2. Materials and Methods

2.1. General Considerations and Instrumentation

All chemicals and solvents were purchased from available commercial sources and were used without further purification.

C/H/N elemental analysis was performed on a Thermo Scientific Flash 2000 analyzer. Infrared spectra were obtained on Jasco FT/IR-4700 via the ATR technique. Thermal stability was studied using a Discovery SDT 650 thermal analyzer (TA Instruments) for simultaneous thermogravimetry (TG) and differential scanning calorimetry (DSC) with a 5 °C/min gradient of temperature and a dynamic air atmosphere (50 mL/min).

The shape of the coordination polyhedron was analyzed with Shape 2.1 software [38,39]. The crystal data and computational data were analyzed and visualized with the help of Mercury [40], Diamond [41], and Vesta 3 [42] software.

2.2. Synthesis of Compound $1 \cdot 1/2\text{Me}_2\text{CO}$

To a stirred solution of 45.8 mg (125 μmol) of $\text{Co}(\text{ClO}_4)_2 \cdot 6\text{H}_2\text{O}$ in 5 mL of acetone, 54.3 mg (250 μmol) of neocuproine (neo) hemihydrate was added. Subsequently, 18.5 mg (125 μmol) of *trans*-cinnamic acid (Hcin) was added, and after its dissolution, 17.3 μL (125 μmol) of triethylamine was added. To this mixture, a filtered solution of 85.5 mg (250 μmol) of sodium tetraphenylborate in 3 mL of acetone was poured in, and the reaction mixture was put into a fridge overnight to yield pink–red crystals of the $[\text{Co}(\text{neo})_2(\text{cin})][\text{BPh}_4] \cdot 1/2\text{Me}_2\text{CO}$ product, which were filtered off, washed with acetone and dried in a vacuum desiccator. The yield was 70.2 mg (60%). Large crystals were obtained by letting the reaction mixture stand undisturbed at room temperature for several days. Anal. Calc. values for $\text{C}_{62.5}\text{H}_{54}\text{BCoN}_4\text{O}_{2.5}$ are as follows: C, 77.32; H, 5.61; N, 5.77. The findings were as follows: C, 77.20; H, 5.65; N, 5.74. FT-IR (ATR, cm^{-1}): 486 w, 549 w, 591 w, 607 w, 654 w, 683 w, 702 s, 729 m, 742 w, 774 w, 813 w, 851 m, 980 w, 1030 w, 1067 w, 1100 w, 1152 w, 1213 w, 1249 w, 1293 w, 1357 m, 1417 s, 1454 w, 1479 w, 1498 m, 1537 w, 1563 w, 1579 w, 1592 w, 1622 w, 1641 m, 1715 w, and 3055 m.

2.3. X-ray Diffraction Analysis

A suitable single crystal of compound **1** was used for the X-ray diffraction experiment using a Rigaku XtaLAB Synergy-I diffractometer with a microfocused RTG-source PhotonJet-i (Cu) and a HyPix Bantam detector. The structure was solved using the SHELXT [43] program and refined via the full matrix least-squares procedure with Olex2.refine [44] in OLEX2 (version 1.5) [45]. The multi-scan absorption corrections were applied using the program CrysAlisPro 1.171.40.82a [46]. The crystal structure of compound **1** was determined at two different temperatures. Initially, the measurement was conducted at room temperature (**1**@293K). The quality of structure refinement was sufficient; however, despite observing a small residual density, we were not able to confirm the presence of the co-crystallized acetone molecule. Upon measuring the crystal structure at 100K, we confirmed the presence of half of an acetone molecule per asymmetric unit in the crystal structure of **1**@100K. The quality of refinement for **1**@100K was significantly better than that of **1**@293K; therefore, the structure of **1**@100K was used for the discussion of the crystal structure.

The non-routine aspects of refinement were as follows. In each cavity within the crystal structure of the studied compound, there resides one acetone molecule. This molecule exhibits disorder due to a combination of positional and disorder involving a special position. Attempts to model this disorder did not yield a satisfactory model with physically meaningful bond lengths and angles for the acetone molecule. Consequently, solvent density was removed using a solvent masking procedure [47] incorporated in OLEX2. As a result, 66 electrons were found within the 358 \AA^3 of a void volume per unit cell. This is consistent with the presence of 0.5 acetone per asymmetric unit cell (equivalent to 16 electrons per asymmetric unit, and 64 electrons per unit cell).

X-ray powder diffraction was measured using Rigaku MiniFlex600.

2.4. Magnetic Measurements

Both static and dynamic magnetic field measurements were performed on a Quantum Design MPMS[®]3 SQUID magnetometer. The temperature-dependent DC magnetic data were acquired at $B = 0.2$ T, and field-dependent isothermal data were measured in the range $B = 0\text{--}7$ T for $T = 2$ and 5 K. The magnetic data were corrected for diamagnetic susceptibility and the signal of the eicosane and sample holder. Diamagnetic susceptibility was estimated using Pascal's constants [48]. The AC susceptibility data were acquired for frequencies 1–1000 Hz for a temperature range up to 6 K, and for a magnetic field of up to 0.5 T.

2.5. Theoretical Calculations

Theoretical calculations were performed with the ORCA 5.0.4 software package [49,50]. First, hydrogen atomic positions were optimized using density functional theory (DFT) on a complex cation fragment extracted from the crystal structure using a BP86 functional [51]. The resolution of identity (RI) approximation [52] with the def2/J auxiliary basis set [53] for Coulomb fitting was employed, and the def2-TZVP basis set for all atoms except for H and C, where less expensive def2-SVP was chosen, was employed [54]. Next, the state-average complete active space self-consistent field (CASSCF) calculations [55] supplemented with N -electron valence second-order perturbation theory (NEVPT2) [56,57] were performed with the active space defined by 7 electrons in 5 d-orbitals, CAS(7e,5o). The same basis sets were used, and additionally, a def2-TZVP/C auxiliary basis set, for correlation fitting [58], and chain-of-sphere approximation (RIJCOSX) [59] were utilized.

3. Results and Discussion

3.1. Synthesis and General Characterization

The synthesis of the title compound was performed in a one-step fashion and did not require an inert atmosphere or anhydrous conditions (Scheme 1). First, the solution of the bis(neocuproine) cobalt(II) precursor complex was prepared by dissolving cobalt perchlorate in acetone, into which two equivalents of neocuproine hemihydrate were added. To this solution, one equivalent of *trans*-cinnamic acid was added and subsequently neutralized with one equivalent of triethylamine. The perchlorate anion was then replaced with tetraphenylborate by adding a solution of two equivalents of sodium tetraphenylborate in acetone. The bulk crystalline product then crystallized in a fridge overnight. Big monocrystals could be obtained via slow and undisturbed isothermal crystallization at room temperature.

The purity of the bulk sample was confirmed via elemental analysis and X-ray powder diffraction (Figure S1), and was in agreement with the crystal structure determined from a single crystal. As is typical for coordination compounds, the FT-IR spectrum (Figure S2) shows weak and downfield-shifted stretching vibrations of the coordinated cinnamate carboxylate group at 1641 cm^{-1} . Intense C=C/C=N stretching vibrations of the aromatic rings are observed around 1417 cm^{-1} owing to the high abundance of aromatic rings in the structure. Bending out-of-plane vibrations of C-H groups and aromatic rings can then be observed at 851 , 729 and 702 cm^{-1} . C-H stretching vibrations of the methyl groups of neocuproine are present at 3055 cm^{-1} . Additionally, the weak vibration at 1715 cm^{-1} suggests a presence of acetone in the crystal structure, which is supported by elemental analysis and thermogravimetric measurements (Figure S6), where the gradual loss of acetone can be observed up to $215\text{ }^{\circ}\text{C}$ followed by the decomposition of the complex. Single-crystal X-ray diffraction analysis confirmed that the title compound, compound **1**, forms an acetone hemisolvate.

Moreover, the photochemical studies were performed in the solid phase on powdered samples utilizing infrared spectroscopy, as a distinct change in FT-IR spectra is observed for *cis*- and *trans*-isomers of cinnamic acid [60]. The individual FT-IR spectra were measured before and after 15 min of simultaneous irradiation with two 6 W UV lamps (254 and 365 nm). Experiments were performed on compound **1**· $\frac{1}{2}\text{Me}_2\text{CO}$ and also dried compound **1** (dried for 30 min at $160\text{ }^{\circ}\text{C}$). No change in spectra was observed after UV irradiation (Figure S7)

for both $1 \cdot 1/2\text{Me}_2\text{CO}$ and **1**, and we therefore conclude that no photoisomerization occurs in the aforementioned conditions. It is most likely that the crystal packing blocks the switching action even after the removal of the acetone molecule.

3.2. Description of the Crystal Structure

Prepared single crystals of $1 \cdot 1/2\text{Me}_2\text{CO}$ (Table 1) belong to the monoclinic system with the space group $P2_1/n$. The asymmetric unit (Figure 1a) contains a complex cation, $[\text{Co}(\text{neo})_2(\text{cin})]^+$, and tetraphenylborate anion, $[\text{BPh}_4]^-$. The cobalt(II) ion is hexacoordinate with the $\{\text{N}_4\text{O}_2\}$ donor set. The geometry of the coordination polyhedron was analyzed with Shape 2.1 software, which determined the shape to be closest to the octahedron (O_h) with a score of 3.150. Inspection of the structure revealed that the deviation from a perfect octahedron is caused mainly by the steric hindrance of the neocuproine methyl groups (Figure 1b). The carboxylate group of the cinnamic acid was coordinated with the metal center with both oxygen atoms almost equidistantly and with bond lengths $d(\text{Co1-O1}) = 2.134(1) \text{ \AA}$ and $d(\text{Co1-O2}) = 2.193(1) \text{ \AA}$, and they form an acute angle of $\angle(\text{O1-Co1-O2}) = 61.12(4)^\circ$, which additionally contributes to the deformation of the coordination polyhedron. To our knowledge, this is the first resolved structure of a mononuclear Co(II) complex where the cinnamate ligand is coordinated bidentately to a single cobalt center. Bond lengths and angles of the coordination polyhedron are summarized in Table 2.

Table 1. Crystallographic data and details of the structure refinement of compound **1**.

Compound	1@100K	1@293K
Empirical formula	$\text{C}_{61}\text{H}_{51}\text{BCoN}_4\text{O}_2$	$\text{C}_{61}\text{H}_{51}\text{BCoN}_4\text{O}_2$
Formula weight	941.79	941.79
<i>T</i> /K	100.0(1)	293(2)
Crystal system, space group	Monoclinic, $P2_1/n$	Monoclinic, $P2_1/n$
Unit cell dimensions		
<i>a</i> /Å	19.2108(2)	19.1252(2)
<i>b</i> /Å	11.24860(10)	11.49370(10)
<i>c</i> /Å	23.3407(2)	23.4892(2)
α (°)	90	90
β (°)	94.0510(10)	93.9610(10)
γ (°)	90	90
<i>V</i> /Å ³	5031.20(8)	5151.05(8)
<i>Z</i> , <i>D_c</i> /g·cm ^{−3}	4, 1.243	4, 1.214
Absorption coefficient (mm ^{−1})	3.042	2.971
<i>F</i> (000)	1972	1972
Reflections collected/unique (<i>R</i> _{int})	32,069/9403 (0.0259)	9354/7614 (0.0248)
Data/restraints/parameters	9403/0/626	9354/0/626
Goodness-of-fit on <i>F</i> ²	1.045	1.048
Final <i>R</i> indices [<i>I</i> > 2σ(<i>I</i>)]	<i>R</i> ₁ = 0.0359, <i>wR</i> ₂ = 0.0932	<i>R</i> ₁ = 0.0608, <i>wR</i> ₂ = 0.1851
<i>R</i> indices (all data)	<i>R</i> ₁ = 0.0401, <i>wR</i> ₂ = 0.0956	<i>R</i> ₁ = 0.0718, <i>wR</i> ₂ = 0.1959
CCDC	2299689	2306545

Despite the abundant presence of aromatic rings, not many of them form any notable π – π interactions, except for stacking contact between the two neocuproine rings of adjacent complex cations of a unit cell (Figure 2a) with a distance between their centroids of 4.030 Å. The closest distance between the cobalt centers is 10.1549(5) Å, which basically rules out potential magnetic exchange interactions.

The cavity, calculated using a probe radius of 1.2 Å, is situated adjacent to the cinnamate fragment and occupies 6.6% of the unit cell volume (Figure 2b). As previously discussed in Section 3.1, each cavity is occupied by one acetone molecule. The precise modeling of their positions was not carried out due to their involvement in intricate positional disorder (for detailed information, please consult the experimental section).

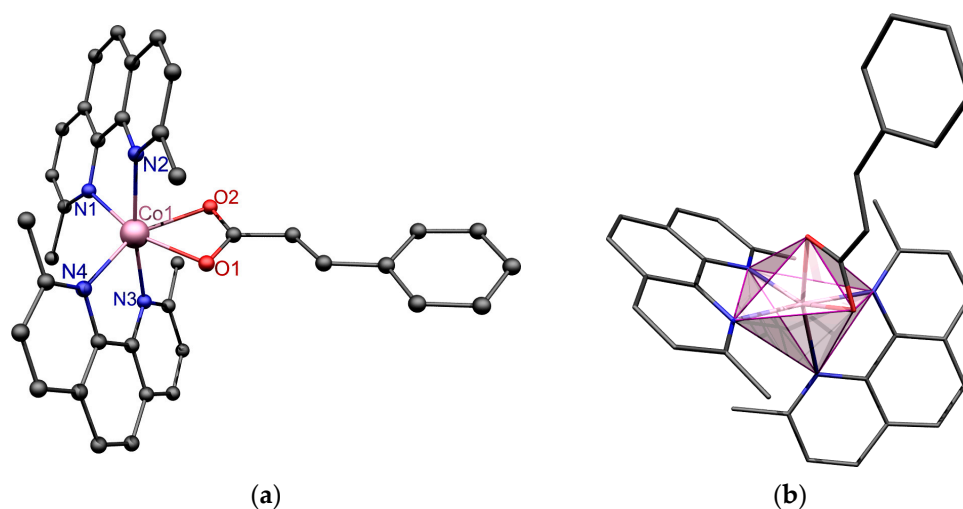


Figure 1. (a) Complex cation of **1** with atom labels of coordination sphere; (b) complex cation with highlighted coordination polyhedron. Hydrogen atoms omitted for clarity.

Table 2. Selected interatomic parameters (\AA , $^\circ$) of the coordination polyhedron.

Co1-N1	2.119(1)
Co1-N2	2.161(1)
Co1-N3	2.137(1)
Co1-N4	2.126(1)
Co1-O1	2.134(1)
Co1-O2	2.193(1)
N1-Co1-N2	78.88(5)
N3-Co1-N4	79.60(5)
O1-Co1-O2	61.12(4)

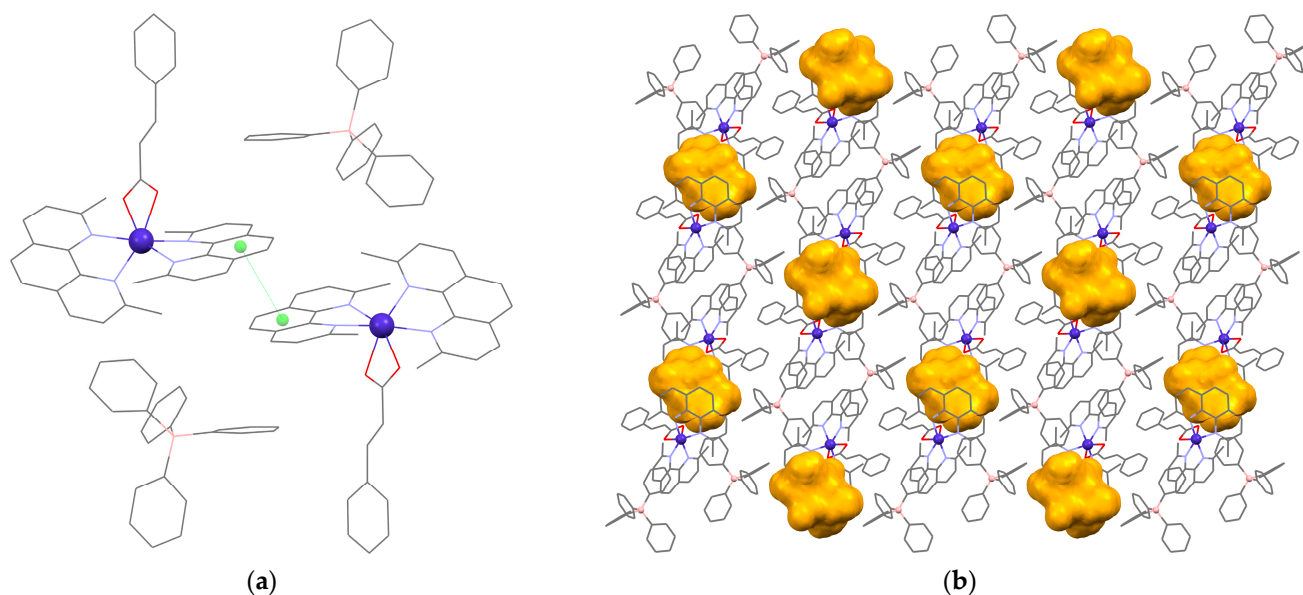


Figure 2. (a) π - π stacking interaction present between the neocuproine rings; centroids represented as green spheres; (b) voids, displayed as yellow surfaces, present in the structure ($3 \times 3 \times 3$ packing; view alongside crystallographic b-axis). Hydrogen atoms omitted for clarity.

3.3. Magnetic Properties

3.3.1. Static Magnetic Properties

The temperature- and field-dependent DC magnetic data of $1.1/2\text{Me}_2\text{CO}$ are presented in Figure 3. The effective magnetic moment at room temperature was determined to be $\mu_{\text{eff}}/\mu_{\text{B}} = 5.12$, a value much higher than the calculated spin-only moment of $\mu_{\text{eff}}/\mu_{\text{B}} = g\sqrt{S(S+1)} = 3.87$ for Co(II) with $S = 3/2$ and $g = 2.0$. The gradual decrease in μ_{eff} was then observed as the sample was cooled down to 1.8 K with $\mu_{\text{eff}}/\mu_{\text{B}} = 3.81$. This strongly suggests that complex $1.1/2\text{Me}_2\text{CO}$ possesses large magnetic anisotropy due to the zero-field splitting [61]. Isothermal molar magnetization was saturated to the value of $M_{\text{mol}}/(N_{\text{A}}\mu_{\text{B}}) = 2.35$ at 2 K and 7 T, which also reflects large magnetic anisotropy because the expected theoretical value is $M_{\text{mol}}/(N_{\text{A}}\mu_{\text{B}}) = g \cdot S = 3$ for $g = 2.0$. Therefore, the simultaneous fit of both temperature- and field-dependent magnetization data was performed using the following spin Hamiltonian (Equation (1)):

$$\hat{H} = D(\hat{S}_z^2 - \hat{S}/3) + E(\hat{S}_x^2 - \hat{S}_y^2) + \mu_{\text{B}}Bg\hat{S} \quad (1)$$

where D and E are axial and rhombic zero-field splitting parameters, respectively, and the last term is the Zeeman term [44]. The analysis was performed for both positive and negative D -parameters; however, only the positive D -parameter was able to reproduce all experimental data properly. The best-fit parameters are $D = +53.2 \text{ cm}^{-1}$, $g_{\text{xy}} = 2.59$, $g_z = 2.0$ (fixed) and $\chi_{\text{TIP}} = 21.3 \times 10^{-9} \text{ m}^3\text{mol}^{-1}$ (Figure 3), where χ_{TIP} represents temperature-independent paramagnetism [62]. Such a large value of D means that magnetic levels originating from $S = 3/2$ are separated by $\Delta = 2D = 106 \text{ cm}^{-1}$. It needs to be mentioned that the fitting procedure was not sensitive to the E parameter.

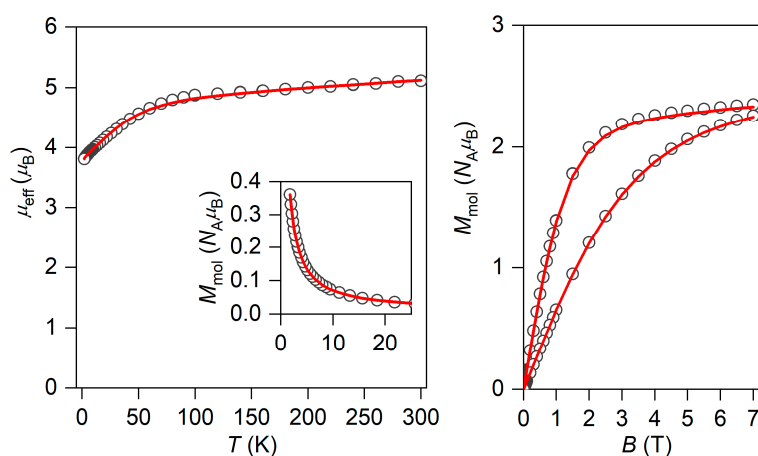


Figure 3. Temperature dependence of the effective magnetic moment, the molar magnetization measured at $B = 0.2 \text{ T}$ (inset), and the isothermal magnetizations measured at $T = 2$ and 5 K for $1.1/2\text{Me}_2\text{CO}$. Empty circles—experimental data; full lines—calculated data with parameters in the text.

3.3.2. Dynamic Magnetic Properties

The AC susceptibility measurements were performed to investigate the possibility of the slow relaxation of the magnetization in compound $1.1/2\text{Me}_2\text{CO}$. Therefore, first, AC susceptibility was measured at $T = 2 \text{ K}$ with a varying static magnetic field as depicted in Figure 4. Evidently, the application of the static magnetic field resulted in an increase in out-of-phase AC susceptibility, and it was possible to analyze experimental data within the range of $B_{\text{dc}} = 0.015$ to 0.2 T with the one-component Debye model in accordance with Equation (2),

$$\chi(2\pi f) = \chi_s + \frac{\chi_T - \chi_s}{1 + (i2\pi f\tau)^{1-\alpha}} \quad (2)$$

which led to isothermal (χ_T) and adiabatic (χ_S) susceptibilities, relaxation times (τ) and distribution parameters (α)—see Table S1. The low-frequency relaxation channel started to appear above 0.2 T and was not analyzed further.

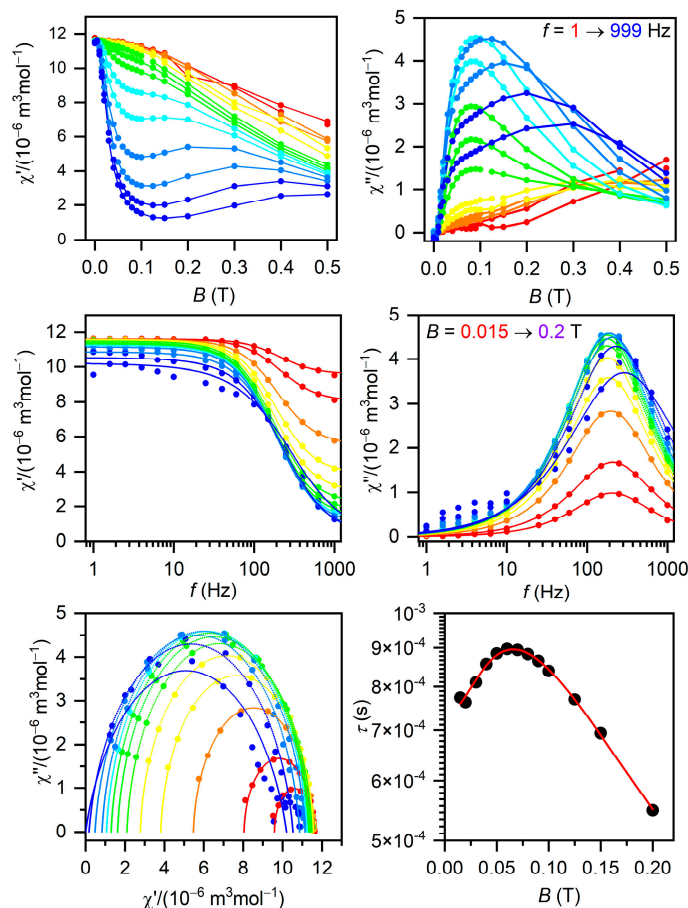


Figure 4. Field-dependent AC susceptibility data for $1.5/2\text{Me}_2\text{CO}$. **Top:** in-phase χ' and out-of-phase χ'' molar susceptibilities at $T = 2$ K (full lines are only guides for the eyes). **Middle:** frequency dependence of in-phase χ' and out-of-phase χ'' molar susceptibilities fitted with one-component Debye's model using Equation (2) (full lines). **Bottom:** the Argand (Cole–Cole) plot with full lines fitted with Equation (2) and, on the right, the fit of resulting relaxation times, τ , with the combination of quantum tunneling and Raman relaxation processes (red line) using Equation (3).

Next, the field-dependence of τ vs. B was subsequently fitted with Equation (3),

$$\frac{1}{\tau} = \frac{b_1}{1 + b_2 B^2} + d \frac{1 + e B^2}{1 + f B^2} T^n \quad (3)$$

which comprises quantum tunnelling (parameters b_1 and b_2) and the Raman process (parameters d , e , f and n). As the expected value of n for Kramers ions is nine, it was fixed during fitting [63,64]. The resulting parameters are $b_1 = 548(67) \text{ s}^{-1}$, $b_2 = 486(136) \text{ T}^{-2}$, $d = 1.60(15) \text{ s}^{-1} \text{ K}^{-9}$, $e = 39(14) \text{ T}^{-2}$, and $f = 4.4(4.4) \text{ T}^{-2}$ —see Figure 4.

As the field dependence of τ showed a maximum in the range 0.05–0.1 T, the static magnetic field was then fixed to $B_{\text{DC}} = 0.09$ T and the AC susceptibility was measured from 1.8 to 6 K—see Figure 5. Herein, the experimental data were again analyzed with a one-component Debye model (Equation (2)) in the range 1.8–5.0 K, and the resulting parameters are listed in Table S2. The temperature dependence of τ was analyzed with Equation (4),

$$\frac{1}{\tau} = CT^n \quad (4)$$

for the Raman process, which resulted in $C = 129(2) \text{ s}^{-1} \text{ K}^{-n}$ and $n = 3.16(2)$ —see Figure 5. The Raman exponent has smaller value than the expected value of 9 for Kramers ions, but when optical and acoustic phonons are considered, values in the range 1–6 are acceptable [33,65–68].

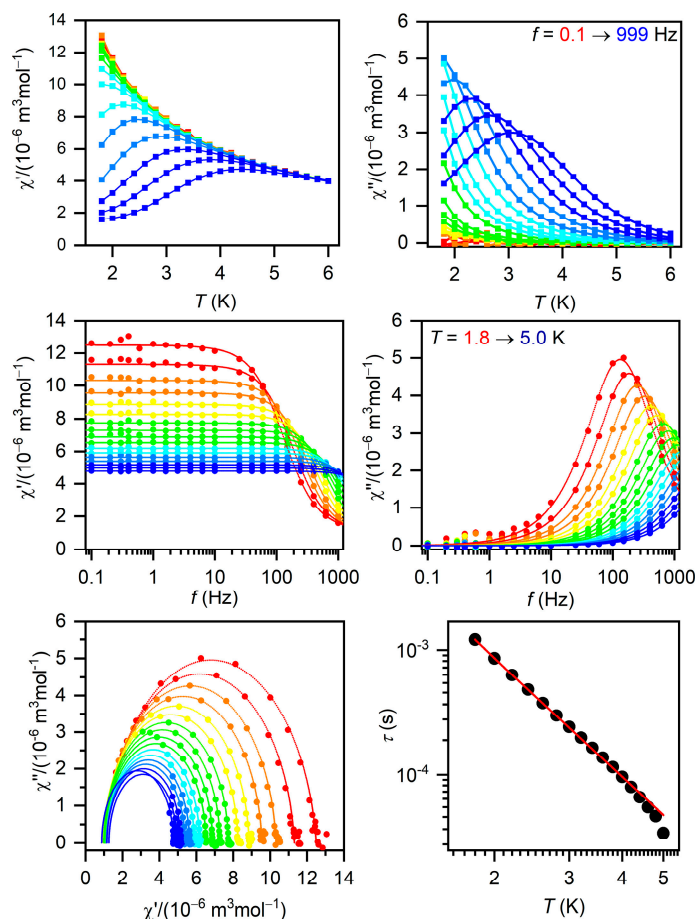


Figure 5. The temperature-dependent AC susceptibility data for **1**. **Top:** in-phase χ' and out-of-phase χ'' molar susceptibilities at $B_{DC} = 0.09 \text{ T}$ (full lines are only guides for the eyes). **Middle:** frequency dependence of in-phase χ' and out-of-phase χ'' molar susceptibilities fitted with the one-component Debye model using Equation (2) (full lines). **Bottom:** the Argand (Cole–Cole) plot with full lines fitted with Equation (2) and, on the right, the fit of resulting relaxation times, τ , with the Raman relaxation processes (red line) using Equation (4).

3.4. Theoretical Calculations

With the help of ab initio ligand field theory (AILFT) [69,70], the energies of d-orbitals were calculated and are depicted in Figure 6. The assignment of the corresponding d-orbitals is shown in Figure S3. The splitting of d-orbitals is typical of pseudo-octahedral Co^{II} complexes, in which t_{2g} and e_g orbitals within ideal O_h symmetry are split due to the lower symmetry of the real complex under study. Hence, the ground ligand field term, T_{1g} , is also split into three terms within the range 0–1500 cm^{-1} (Figure 6, middle). Finally, the spin–orbit interactions result in ligand field multiplets, which are shown in Figure 6 (right), where the energy separation of two lowest Kramers doublets is $\Delta = 125 \text{ cm}^{-1}$, and this value is close to that derived from the fitting of the experimental magnetic data. The further analysis resulted in the zero-field splitting parameters $D = 60.7 \text{ cm}^{-1}$, $E/D = 0.149$, $g_x = 2.453$, $g_y = 2.666$ and $g_z = 2.004$. The actual D -tensor is depicted in Figure S4. As the E/D ratio is non-zero, the resulting magnetic anisotropy can be either of an easy-plane or easy-axis type [71]. Hence, the ground state Kramers doublet was analyzed with the effective spin 1/2, which resulted in $g_1 = 2.010$, $g_2 = 3.753$, $g_3 = 6.316$ and $g_{\text{ave}} = 4.026$. Thus, it can be

concluded that it holds $g_1, g_2 < g_{ave}$ and $g_{ave} < g_3$, which means there is an easy-axis type of the magnetic anisotropy in compound **1** according to CASSCF/NEVPT2 calculations. This is in agreement with the calculated 3D magnetization depicted in Figure S5 ($T = 2$ K and $B = 0.1$ T), which shows a shape characteristic of an easy-axis type of anisotropy. Therefore, further analysis was conducted with the SINGLE_ANISO module [72], in order to compute the ab initio magnetization blocking barrier displayed in Figure 7. The corresponding matrix element of the transversal magnetic moment between ground states with opposite magnetization is equal to 0.96; thus, it is larger than 0.1, which suggests a large predisposition for the quantum tunneling of magnetization, which is in line with the experimental AC data.

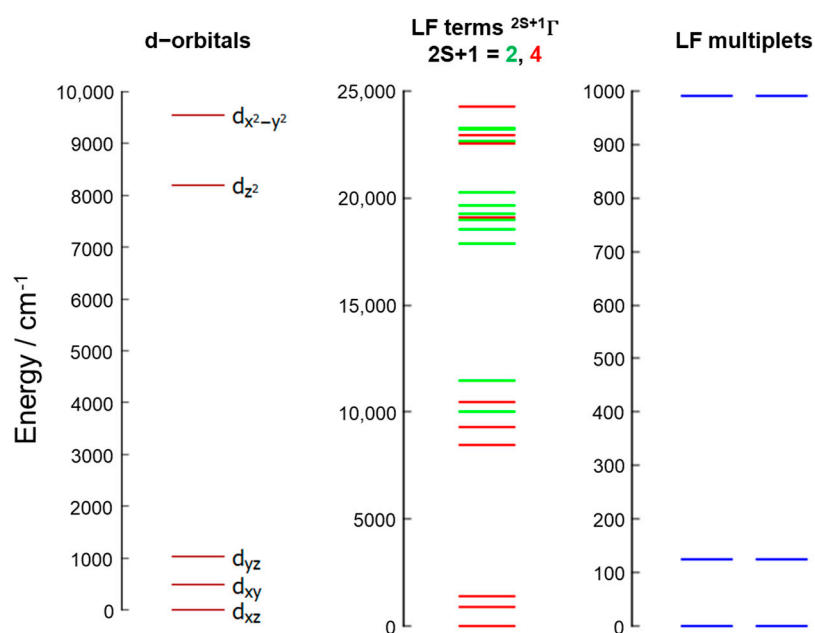


Figure 6. Energy levels of d-orbitals (**left**) calculated using ab initio ligand field theory, and ligand field terms (**middle**) and multiplets (**right**) calculated using CASSCF/NEVPT2 for compound **1**.

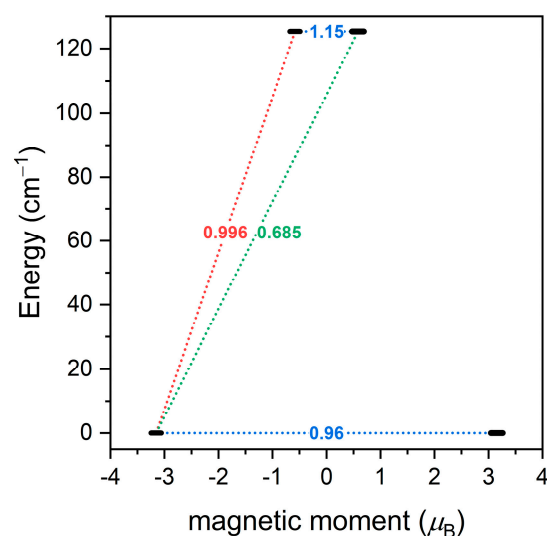


Figure 7. The results of the SINGLE_ANISO analysis of CASSCF/NEVPT2 calculations for compound **1**. The numbers presented in the plots represent the corresponding matrix element of the transversal magnetic moment (for values larger than 0.1, an efficient relaxation mechanism is expected).

4. Conclusions

The crystal structure shows an unusual bidentate mode of coordination of cinnamate into a single Co(II) center, contributing to the deformation of the coordination polyhedron. Single-crystal measurement at room temperature most likely led to the formation of voids due to solvent loss.

Dynamic magnetic measurements revealed that the title compound, compound $1 \cdot 1/2\text{Me}_2\text{CO}$, exhibits a slow relaxation of magnetization in an applied external magnetic field with a predominant Raman process as the relaxation mechanism. Static magnetic measurements, in combination with theoretical calculations, showed that compound $1 \cdot 1/2\text{Me}_2\text{CO}$ possesses large magnetic anisotropy of the axial type due to zero-field splitting with a positive D value of $+53.2\text{ cm}^{-1}$. Despite the evident discrepancy between fitted and calculated rhombicity, the experimental data are in good agreement with CASSCF/NEVPT2 results. It is well known that static magnetic measurements usually cannot determine rhombicity with high precision, and more sensitive methods, such as EPR, would have to be utilized to resolve both E and g -factors more accurately.

To conclude, a Co(II) field-induced single-molecule magnet with a potentially photoisomerizable ligand was prepared and, this opens up an avenue for the further investigation of similar compounds, which would be interesting not only in the field of molecular magnetism, but also in the field of photomagnetism.

Supplementary Materials: The following supporting information can be downloaded at <https://www.mdpi.com/article/10.3390/magnetochemistry9110229/s1>. Figure S1. X-ray powder diffraction of compound $1 \cdot 1/2\text{Me}_2\text{CO}$ measured at room temperature and compared to the simulation calculated from SC-XRD data measured at 293 K. Figure S2. FT-IR spectrum of compound $1 \cdot 1/2\text{Me}_2\text{CO}$. Figure S3. Assignment of energy levels of d-orbitals of compound **1**. Figure S4. D-tensor (NEVPT2) of compound **1**. Axes: x—red; y—green; z—blue. Figure S5. Calculated 3D magnetization of compound **1** at $T = 2\text{ K}$ and $B = 0.1\text{ T}$. Figure S6. The results of simultaneous TG/DSC thermal analysis of $1 \cdot 1/2\text{Me}_2\text{CO}$ ($5\text{ }^\circ\text{C/min}$, 50 mL/min air atmosphere) depicted as TG (in blue) and DSC (green) curves; TG = thermogravimetry, and DSC = differential scanning calorimetry. Figure S7. Top—compound $1 \cdot 1/2\text{Me}_2\text{CO}$ before (blue) and after (red) irradiation. Bottom—compound **1** before (blue) and after (red) irradiation. Table S1. The parameters of the one-component Debye model used to analyze the field-dependent AC susceptibility data of compound $1 \cdot 1/2\text{Me}_2\text{CO}$; Table S2. The parameters of the one-component Debye model used to analyze the temperature-dependent AC susceptibility data of compound $1 \cdot 1/2\text{Me}_2\text{CO}$.

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