



Article A Family of Transformed Difference Schemes for Nonlinear Time-Fractional Equations

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Abstract: In this paper, we present a class of finite difference methods for numerically solving fractional differential equations. Such numerical schemes are developed based on the change in variable and piecewise interpolations. Error analysis of the numerical schemes is obtained by using a Grönwall-type inequality. Numerical examples are given to confirm the theoretical results.

Keywords: time-fractional parabolic problems; change in variable; convergence; optimal error estimates; linearized schemes

1. Introduction

We aim to develop a family of effective numerical methods for solving the following nonlinear time-fractional differential equations:

$$D_*^{\alpha} y = f(t, y(t)), \quad y(0) = y_0, \quad t \in (0, T], \tag{1}$$

where $0 < \alpha < 1$, D_*^{α} is the the differential operator in sense of Caputo, given by

$$D_*^{\alpha}y = \frac{1}{\Gamma(1-\alpha)}\int_0^t (t-r)^{-\alpha}y'(r)dr.$$

The fractional differential equations provide a powerful tool to describe many natural phenomena in the fields of physics [1–3], economics [4] and biology [5].

Developing and analyzing highly effective numerical methods for fractional differential equations has been one of the hot topics. The widely used numerical methods are the L1 schemes [6–13] and L2-schemes [14–16]. Such schemes are developed by using piecewise interpolations. The optimal convergence results of L1-type schemes for time-dependent partial equations can be obtained by using the fractional Grönwall type inequalities [17–19]. Moreover, L1-type schemes can be accelerated by using sum-of-exponentials approximations [20–22]. Other widely used schemes are the so-called backward differentiation formula (BDF) convolution quadrature (CQ) methods. The CQ methods were originally proposed in [23,24] and further investigated in [25–27]. That aside, some transformed finite difference methods were constructed based on some change in variables [28–30]. More numerical schemes as well as their numerical analysis can be found in an incomplete list of references [31–33].

It is widely accepted that Equation (1) is equivalent to the following Volterra integral equation [34]:

$$y(t) = y_0 + \frac{1}{\Gamma(\alpha)} \int_0^t (t - r)^{\alpha - 1} f(r, y(r)) dr.$$
 (2)

Here and below, we always assume that the function f is Lipschitz continuous with respect to the second argument on a suitable set G and the Lipschitz constant is L. Moreover,



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Copyright: © 2023 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). suppose $f \in C^3(G)$. Then, there exists a function $\psi \in C^2[0, T]$ and some $c_1, c_2, \dots, c_J \in \mathbb{R}$ and $d_1, d_2, \dots, d_{\hat{f}} \in \mathbb{R}$ such that the solution to Equation (1) can be given by (see, e.g., [35–38])

$$y(t) = \psi(t) + \sum_{j=1}^{J} c_j t^{j\alpha} + \sum_{j=1}^{\hat{f}} d_j t^{1+j\alpha},$$
(3)

where $J := \lfloor 2/\alpha \rfloor - 1$. Clearly, the typical solution to the problem in Equation (1) has an initial layer at the beginning, and y_t blows up as $t \to 0^+$. In this paper, we aim to present effective numerical schemes to solve the fractional problems, taking the initial layer into account.

The new numerical schemes are developed based on the following changes in variables:

$$t = s^{\eta/\alpha}, \ \eta \in \mathbb{N}^+, \tag{4}$$

Equations (1) or (2) is equivalent to the following integral equation:

$$y(s^{\eta/\alpha}) = y_0 + \frac{1}{\Gamma(\alpha)} \int_0^{s^{\eta/\alpha}} (s^{\eta/\alpha} - r)^{\alpha - 1} f(r, y(r)) dr$$

$$= y_0 + \frac{1}{\Gamma(\alpha)} \int_0^s (s^{\eta/\alpha} - u^{\eta/\alpha})^{\alpha - 1} f(u^{\eta/\alpha}, y(u^{\eta/\alpha})) \frac{\eta}{\alpha} u^{\eta/\alpha - 1} du$$

$$= y_0 + \frac{\eta}{\Gamma(1 + \alpha)} \int_0^s (s^{\eta/\alpha} - r^{\eta/\alpha})^{\alpha - 1} r^{\eta/\alpha - 1} f(r^{\eta/\alpha}, y(r^{\eta/\alpha})) dr,$$
(5)

Furthermore, let

$$z(s) = y(s^{\eta/\alpha}),\tag{6}$$

Equation (5) becomes

$$z(s) = y_0 + \frac{\eta}{\Gamma(1+\alpha)} \int_0^s (s^{\eta/\alpha} - r^{\eta/\alpha})^{\alpha-1} r^{\eta/\alpha-1} f(r^{\eta/\alpha}, z(r)) dr, \ 0 \le s \le T^{\alpha/\eta}.$$
(7)

Then, the solution to Equation (7) has the form

$$z(s) = \psi(s^{\eta/\alpha}) + \sum_{j=1}^{J} c_j s^{j\eta} + \sum_{j=1}^{\hat{J}} s^{\frac{\eta(1+j\alpha)}{\alpha}},$$
(8)

where $\psi \in C^2[0, T]$. Clearly, by the changes in variables, the initial layer will vanish, and the exact solution will become smoother. Then, the idea of developing effective numerical methods for solving Equation (1) is as follows:

- Apply the change in Equation (4) with a suitable parameter η to obtain Equation (7) and its regularity of the solution.
- Develop numerical methods based on the smoothness of the solution to Equation (7).
- Recover the numerical solution by using the simple inverse change $y(s) = z(s^{\alpha/\eta})$.

The rest of the paper is organized as follows. In Section 2, we present a family of new numerical schemes and investigate the numerical schemes' convergence results. In Section 3, we present some numerical results to confirm the theoretical findings. Finally, our conclusions are presented in Section 4.

2. Construction of the Numerical Methods

In this section, we develop some numerical methods and present the numerical results

based on the change in variable. Here and below, we always set the step size $h = \frac{T^{w, y}}{N}$, with *N* being a given integer and $s_i = ih$, $i = 0, 1, 2, \dots, N$. In what follows, we will present two numerical methods for solving Equation (7).

2.1. First-Order Accurate Methods

In this subsection, we present a numerical method based on the following change in variable:

$$t = s^{1/\alpha}$$

With Equations (7) and (8), one can check that its solution $z(s) \in C[0, T]$ with the change in variable. Therefore, we propose a first-order accurate method based on the product rectangle rule for every interval; in other words, we propose

$$\int_{s_{m-1}}^{s_m} (s_i^{1/\alpha} - r^{1/\alpha})^{\alpha - 1} r^{1/\alpha - 1} f(r^{1/\alpha}, z(r)) dr \approx f(s_m^{1/\alpha}, z(s_m)) \int_{s_{m-1}}^{s_m} (s_i^{1/\alpha} - r^{1/\alpha})^{\alpha - 1} r^{1/\alpha - 1} dr.$$

Then, it follows from Equation (7) that

$$z(s_{i}) = z_{0} + \frac{1}{\Gamma(1+\alpha)} \sum_{m=1}^{i} \int_{s_{m-1}}^{s_{m}} (s_{i}^{1/\alpha} - r^{1/\alpha})^{\alpha-1} r^{1/\alpha-1} f(r^{1/\alpha}, z(r)) dr$$

$$= z_{0} + \frac{1}{\Gamma(1+\alpha)} \sum_{m=1}^{i} f(s_{m}^{1/\alpha}, z(s_{m})) \int_{s_{m-1}}^{s_{m}} (s_{i}^{1/\alpha} - r^{1/\alpha})^{\alpha-1} r^{1/\alpha-1} dr + R_{1}^{i} \qquad (9)$$

$$= z_{0} + \frac{1}{\Gamma(1+\alpha)} \sum_{m=1}^{i} a_{m}^{i} f(s_{m}^{1/\alpha}, z(s_{m})) + R_{1}^{i},$$

where $z_0 = y_0$ and the truncation error is

$$R_{1}^{i} = \frac{1}{\Gamma(1+\alpha)} \sum_{m=1}^{i} \int_{s_{m-1}}^{s_{m}} (s_{i}^{1/\alpha} - r^{1/\alpha})^{\alpha - 1} r^{1/\alpha - 1} \Big(f(r^{1/\alpha}, z(r)) - (f(s_{m}^{1/\alpha}, z(s_{m}))) \Big) dr = \mathcal{O}(h),$$

while the coefficients are

$$a_m^i = \int_{s_{m-1}}^{s_m} (s_i^{1/\alpha} - r^{1/\alpha})^{\alpha - 1} r^{1/\alpha - 1} dr.$$
⁽¹⁰⁾

Let z_i be a numerical approximation to $z(s_i)$. By replacing $z(s_i)$ with z_i and omitting the truncation error R_i , we have the following numerical scheme:

Scheme I:
$$z_i = z_0 + \frac{1}{\Gamma(1+\alpha)} \sum_{m=1}^i a_m^i f(s_m^{1/\alpha}, z_m).$$
 (11)

Scheme I is an implicit method. At each step, iterative processes are required to solve the nonlinear equations. In order to reduce the computational cost, a linearized scheme can be developed as follows:

Scheme II:
$$z_i = z_0 + \frac{1}{\Gamma(1+\alpha)} \sum_{m=1}^l a_m^i f(s_m^{1/\alpha}, z_{m-1}).$$
 (12)

The convergence results of the proposed schemes rely heavily on the following lemmas:

Lemma 1 ([39]). Let x_i , $0 \le i \le N$ be a sequence of non-negative real numbers. If

$$x_i \leq \psi_i + Mh^{\sigma+1-lphaeta} \sum_{j=0}^{i-1} \frac{j^{\sigma} x_j}{(i^{eta} - j^{eta})^{lpha}}, \quad 0 \leq i \leq N,$$

where $0 < \alpha < 1$, $1 \le \beta \le \sigma + 1$, $\sigma \ge 0$, *M* is a positive constant, and ψ_i , $0 \le i \le N$, is a monotonic increasing sequence of non-negative real numbers, then

$$x_{i} \leq \psi_{i} \sum_{n=0}^{\infty} \left(\frac{M(ih)^{\sigma+1-\alpha\beta}}{\beta}\right)^{n} \hat{B}_{n}(\alpha,\beta,\sigma), \quad 0 \leq i \leq N.$$
$$\hat{B}_{n}(\alpha,\beta,\sigma) = \begin{cases} 1, & n=0, \\ \prod_{l=1}^{n} B\left(\frac{l}{\beta}(\sigma+1-\alpha\beta)+\alpha,(1-\alpha)\right), & n \geq 1, \end{cases}$$

where *B* is the Beta function.

When $\sigma + 1 - \beta = 0$, the following is true:

$$x_i \leq \psi_i E_{1-lpha}(rac{M\Gamma(1-lpha)}{eta}(ih)^{eta(1-lpha)}), \ 0 \leq i \leq N,$$

where $E_{\alpha}(z) = \sum_{k=0}^{\infty} \frac{z^k}{\Gamma(1+k\alpha)}$ is the Mittag-Leffler function.

Now, we have the following convergence results:

Theorem 1. Suppose $f \in C^3(G)$. Then, the schemes in Equations (11) and (12) are first-order accurate.

Proof. Let $e_i = z(s_i) - z_i$. Subtracting Equation (11) from Equation (9) gives

$$|e_{i}| = \left|e_{0} + \frac{1}{\Gamma(1+\alpha)} \sum_{m=1}^{i} a_{m}^{i}(f(s_{m}^{1/\alpha}, z(s_{m})) - f(s_{m}^{1/\alpha}, z_{m})) + R_{1}^{i}\right|$$

$$\leq |e_{0}| + \frac{L}{\Gamma(1+\alpha)} \sum_{m=1}^{i} a_{m}^{i}|e_{i}| + |R_{1}^{i}|.$$
(13)

It follows from Equation (10) that, for $1 \le m \le i - 1$, we have

$$a_{m}^{i} = \int_{s_{m-1}}^{s_{m}} \frac{r^{\frac{1}{\alpha}-1}}{(s_{i}^{\frac{1}{\alpha}}-r^{\frac{1}{\alpha}})^{1-\alpha}} = h \frac{\zeta^{\frac{1}{\alpha}-1}}{(s_{i}^{\frac{1}{\alpha}}-\zeta^{\frac{1}{\alpha}})^{1-\alpha}}$$
(14)

$$\leq h \frac{s_m^{\frac{1}{\alpha}-1}}{(s_i^{\frac{1}{\alpha}} - s_m^{\frac{1}{\alpha}})^{1-\alpha}} = h \frac{m^{\frac{1}{\alpha}-1}}{(i^{\frac{1}{\alpha}} - m^{\frac{1}{\alpha}})^{\alpha-1}},$$
(15)

where ζ belongs to the interval $[s_{m-1}, s_m]$ and the mean value theorem of the integral is used. Noting that

$$a_i^i = (s_i^{1/\alpha} - s_{i-1}^{1/\alpha})^\alpha = \xi^{1/\alpha - 1}h, \quad (\xi \in (s_{i-1}, s_i)).$$
(16)

then $a_i^i |e^i|$ can be absorbed by the left-hand side of the inequality when *h* is sufficiently small. By combining Equations (14) and (15) with the truncation error R_1^i , we can obtain

$$\begin{aligned} |e_i| &\leq |e_0| + \frac{C_0 L}{\Gamma(1+\alpha)} \sum_{m=1}^{i-1} h \frac{m^{\frac{1}{\alpha}-1}}{(i^{\frac{1}{\alpha}} - m^{\frac{1}{\alpha}})^{1-\alpha}} + \frac{L}{\Gamma(1+\alpha)} \xi^{\frac{1}{\alpha}-1} h + C_2 h \\ &\leq C_3 h + \frac{C_0 L}{\Gamma(1+\alpha)} \sum_{m=1}^{i-1} h \frac{m^{\frac{1}{\alpha}-1}}{(i^{\frac{1}{\alpha}} - m^{\frac{1}{\alpha}})^{1-\alpha}}, \end{aligned}$$

where C_2 and C_3 are two constants independent on *h*. When applying Lemma 1, the above inequality yields that

$$e_i \leq C_3 h E_{1-\alpha}(\frac{\alpha C_0 \Gamma(1-\alpha)}{\Gamma(1+\alpha)}(ih)^{\frac{1-\alpha}{\alpha}}).$$

Now, we conclude that Scheme I is first-order accurate. The convergence of Scheme II can be obtained in a similar fashion. $\hfill\square$

2.2. Second-Order Accurate Methods

In this subsection, we present the numerical methods based on the following change in variable:

$$t = s^{2/\alpha}$$

Again, through Equations (7) and (8), one can check that its solution $z(s) \in C^2[0, T]$ with the change in variable. Therefore, we propose a second-order accurate method based on the product trapezoidal quadrature rule; in other words, we have

$$\int_{s_m}^{s_{m+1}} (s_i^{2/\alpha} - r^{2/\alpha})^{\alpha - 1} r^{2/\alpha - 1} f(r^{2/\alpha}, z(r)) dr$$

$$\approx \int_{s_m}^{s_{m+1}} (s_i^{2/\alpha} - r^{2/\alpha})^{\alpha - 1} r^{2/\alpha - 1} \left(\frac{s_{m+1} - r}{h} f(s_m^{2/\alpha}, z(s_m)) + \frac{r - s_m}{h} f(s_{m+1}^{2/\alpha}, z(s_m+1)) \right) dr.$$

Then, it follows from Equation (7) again that

$$\begin{aligned} z(s_i) &= z_0 + \frac{2}{\Gamma(1+\alpha)} \sum_{m=0}^{i-1} \int_{s_m}^{s_{m+1}} (s_i^{2/\alpha} - r^{2/\alpha})^{\alpha - 1} r^{2/\alpha - 1} f(r^{2/\alpha}, z(r)) dr \\ &= z_0 + \frac{2}{\Gamma(1+\alpha)} \sum_{m=0}^{i-1} \int_{s_m}^{s_{m+1}} (s_i^{2/\alpha} - r^{2/\alpha})^{\alpha - 1} r^{2/\alpha - 1} \left(\frac{s_{m+1} - r}{h} f(s_m^{2/\alpha}, z(s_m)) + (17) \right) \\ &\quad \frac{r - s_m}{h} f(s_{m+1}^{2/\alpha}, z(s_m+1)) dr + R_2^i \\ &= z_0 + \frac{2}{\Gamma(1+\alpha)} \sum_{m=0}^{i} b_m^i f(s_m^{2/\alpha}, z(s_m)) + R_2^i, \end{aligned}$$

where the truncation error is

$$\begin{aligned} R_2^i &= \frac{2}{\Gamma(1+\alpha)} \sum_{m=0}^{i-1} \int_{s_m}^{s_{m+1}} (s_i^{2/\alpha} - r^{2/\alpha})^{\alpha-1} r^{2/\alpha-1} \Big(f(r^{2/\alpha}, z(r)) - \frac{s_{m+1} - r}{h} f(s_m^{2/\alpha}, z(s_m)) - \frac{r - s_m}{h} f(s_{m+1}^{2/\alpha}, z(s_m+1)) \Big) dr = \mathcal{O}(h^2), \end{aligned}$$

and the coefficient is

$$\begin{split} b_0^i &= \int_{s_0}^{s_1} (s_i^{2/\alpha} - r^{2/\alpha})^{\alpha - 1} r^{2/\alpha - 1} \frac{s_1 - r}{h} dr = \frac{1}{2} s_i^2 - \frac{1}{2} (s_i^{2/\alpha} - s_1^{2/\alpha})^{\alpha} - \frac{1}{h} \int_{s_0}^{s_1} (s_i^{2/\alpha} - r^{2/\alpha})^{\alpha - 1} r^{2/\alpha} dr \\ &= \frac{1}{2} s_i^2 - \frac{1}{2} (s_i^{2/\alpha} - s_1^{2/\alpha})^{\alpha} - \frac{\alpha}{2h} \int_{s_0^{2/\alpha}}^{s_1^{2/\alpha}} (s_i^{2/\alpha} - u)^{\alpha - 1} u^{\alpha/2} du \\ &= \frac{1}{2} s_i^2 - \frac{1}{2} (s_i^{2/\alpha} - s_1^{2/\alpha})^{\alpha} - \frac{\alpha s_i^{2+1}}{2h} \int_0^{(s_1/s_i)^{2/\alpha}} (1 - r)^{\alpha - 1} r^{\alpha/2} dr \\ &= \frac{1}{2} s_i^2 - \frac{1}{2} (s_i^{2/\alpha} - s_1^{2/\alpha})^{\alpha} - \frac{\alpha s_i^{2+1}}{2h} \mathbf{B} (i^{-2/\alpha}, \alpha/2 + 1, \alpha), \end{split}$$

where **B** $(i^{-2/\alpha}, \alpha/2 + 1, \alpha)$ is the product of an incomplete beta function and beta function. For $1 \le m \le i - 1$, we have

$$\begin{split} b_{m}^{i} &= \int_{s_{m}}^{s_{m+1}} (s_{i}^{2/\alpha} - r^{2/\alpha})^{\alpha - 1} r^{2/\alpha - 1} \left(\frac{s_{m+1} - r}{h}\right) dr + \int_{s_{m-1}}^{s_{m}} (s_{i}^{2/\alpha} - r^{2/\alpha})^{\alpha - 1} r^{2/\alpha - 1} \left(\frac{r - s_{m-1}}{h}\right) dr \\ &= \frac{m+1}{2} \left((s_{i}^{2/\alpha} - s_{m}^{2/\alpha})^{\alpha} - (s_{i}^{2/\alpha} - s_{m+1}^{2/\alpha})^{\alpha} \right) - \frac{m-1}{2} \left((s_{i}^{2/\alpha} - s_{m-1}^{2/\alpha})^{\alpha} - (s_{i}^{2/\alpha} - s_{m}^{2/\alpha})^{\alpha} \right) \\ &- \frac{\alpha s_{i}^{2+1}}{2h} \int_{(s_{m}/s_{i})^{2/\alpha}}^{(s_{m+1}/s_{i})^{2/\alpha}} (1 - r)^{\alpha - 1} r^{\alpha/2} dr + \frac{\alpha s_{i}^{2+1}}{2h} \int_{(s_{m-1}/s_{i})^{2/\alpha}}^{(s_{m}/s_{i})^{2/\alpha}} (1 - r)^{\alpha - 1} r^{\alpha/2} dr \\ &= \frac{m+1}{2} \left((s_{i}^{2/\alpha} - s_{m}^{2/\alpha})^{\alpha} - (s_{i}^{2/\alpha} - s_{m+1}^{2/\alpha})^{\alpha} \right) - \frac{m-1}{2} \left((s_{i}^{2/\alpha} - s_{m-1}^{2/\alpha})^{\alpha} - (s_{i}^{2/\alpha} - s_{m}^{2/\alpha})^{\alpha} \right) \\ &- \frac{\alpha s_{i}^{2+1}}{2h} \mathbf{B}((m+1/i)^{2/\alpha}, \alpha/2 + 1, \alpha) + 2 \frac{\alpha s_{i}^{2+1}}{2h} \mathbf{B}((m/i)^{2/\alpha}, \alpha/2 + 1, \alpha) \\ &- \frac{\alpha s_{i}^{2+1}}{2h} \mathbf{B}(((m-1)/i)^{2/\alpha}, \alpha/2 + 1, \alpha). \end{split}$$

In addition, we have

$$\begin{split} b_i^i &= \int_{s_{i-1}}^{s_i} (s_i^{2/\alpha} - r^{2/\alpha})^{\alpha - 1} r^{2/\alpha - 1} \frac{r - s_{i-1}}{h} dr \\ &= \frac{\alpha s_i^{2+1}}{2h} \mathbf{B}(1, \frac{\alpha}{2} + 1, \alpha) - \frac{\alpha s_i^{2+1}}{2h} \mathbf{B}(((i-1)/i)^{2/\alpha}, \frac{\alpha}{2} + 1, \alpha) - \frac{i - 1}{2} (s_i^{2/\alpha} - s_{i-1}^{2/\alpha})^{\alpha}. \end{split}$$

Again, by replacing $z(s_i)$ with z_i and omitting the truncation error R_i in Equation (9), we have the following numerical scheme:

Scheme III:
$$z_i = z_0 + \frac{2}{\Gamma(1+\alpha)} \sum_{m=0}^i b_m^i f(s_m^{2/\alpha}, z_m).$$
 (18)

By applying the Newton linearized method to approximate the nonlinear term, we have the following linearized scheme:

Scheme IV:
$$z_i = z_0 + \frac{2}{\Gamma(1+\alpha)} \sum_{m=0}^{i-1} b_m^i f(s_m^{2/\alpha}, z_m) + \frac{2}{\Gamma(1+\alpha)} b_i^i \Big(f(s_i^{2/\alpha}, z_{i-1}) + f_2(s_i^{2/\alpha}, z_{i-1})(z_i - z_{i-1}) \Big).$$
 (19)

where $f_2(s_i^{2/\alpha}, z_{i-1}) = \frac{\partial}{\partial y} f(s, y) \Big|_{s=s_i^{2/\alpha}, y=z_{i-1}}$.

By applying the extrapolation to approximate the nonlinear term, we obtain the following linearized scheme:

Scheme V:
$$z_0 = z_0$$
,
 $z_1 = z_0 + \frac{2}{\Gamma(1+\alpha)} b_0^1 f(s_0^{2/\alpha}, z_0) + \frac{2}{\Gamma(1+\alpha)} b_1^1 \Big(f(s_1^{2/\alpha}, z_0) + f_2(s_1^{2/\alpha}, z_0)(z_1 - z_0) \Big),$ (20)
 $z_i = z_0 + \frac{2}{\Gamma(1+\alpha)} \sum_{m=0}^{i-1} b_m^i f(s_m^{2/\alpha}, z_m) + \frac{2}{\Gamma(1+\alpha)} b_i^i f(s_i^{2/\alpha}, 2z_{i-1} - z_{i-2}), \quad i \ge 2.$

Next, we have the following convergence results:

Theorem 2. Suppose $f \in C^3(G)$. Then, the schemes in Equations (18)–(20) are second-order accurate.

Proof. Let $e_i = z(s_i) - z_i$. Subtracting Equation (11) from Equation (9) gives

$$|e_{i}| = \left|e_{0} + \frac{1}{\Gamma(1+\alpha)} \sum_{m=1}^{i} b_{m}^{i}(f(s_{m}^{1/\alpha}, z_{m}) - f(s_{m}^{1/\alpha}, z_{m})) + R_{2}^{i}\right|$$

$$\leq |e_{0}| + \frac{L}{\Gamma(1+\alpha)} \sum_{m=1}^{i} b_{m}^{i}|e_{i}| + |R_{2}^{i}|.$$
(21)

Now, we present some estimates for the coefficients b_m^i . First, it holds that

$$b_0^i = \int_{s_0}^{s_1} (s_i^{2/\alpha} - r^{2/\alpha})^{\alpha - 1} r^{2/\alpha - 1} \frac{s_1 - r}{h} dr = \frac{1}{2} s_i^2 - \frac{1}{2} (s_i^{2/\alpha} - s_1^{2/\alpha})^{\alpha} \le C_1 h^2$$

For $1 \le m \le i - 1$, it holds that

$$\begin{split} b_m^i &= \int_{s_m}^{s_{m+1}} (s_i^{2/\alpha} - r^{2/\alpha})^{\alpha - 1} r^{2/\alpha - 1} \left(\frac{s_{m+1} - r}{h}\right) dr + \int_{s_{m-1}}^{s_m} (s_i^{2/\alpha} - r^{2/\alpha})^{\alpha - 1} r^{2/\alpha - 1} \left(\frac{r - s_{m-1}}{h}\right) dr \\ &= (s_i^{2/\alpha} - \xi_m^{2/\alpha})^{\alpha - 1} \xi_m^{2/\alpha - 1} (s_{m+1} - \xi_m) + (s_i^{2/\alpha} - \xi_m^{2/\alpha})^{\alpha - 1} \xi_m^{2/\alpha - 1} (\xi_m - s_{m-1}) \\ &\leq C_2 h^2 (i^{2/\alpha} - m^{2/\alpha})^{\alpha - 1} m^{2/\alpha - 1}, \end{split}$$

where C_2 is a constant independent of h, $\xi_m \in (s_m, s_{m+1})$, and $\xi_m \in (s_{m-1}, s_m)$. Noting that

$$b_{i}^{i} = \int_{s_{i-1}}^{s_{i}} (s_{i}^{2/\alpha} - r^{2/\alpha})^{\alpha - 1} r^{2/\alpha - 1} \frac{r - s_{i-1}}{h} dr \le h (s_{i}^{2/\alpha} - s_{i-1}^{2/\alpha})^{\alpha - 1} s_{i}^{2/\alpha - 1}$$

then $b_i^i |e^i|$ can be absorbed by the left-hand side of the inequality when *h* is sufficiently small.

Now, together with the estimates for b_m^i and Lemma 1, we conclude that Scheme III is second-order accurate. The rest of the results can be obtained in a similar manner. \Box

3. Applications

In this section, several numerical examples are given to illustrate the convergence results, and the L^{∞} norm of the error is computed with different α . All numerical examples are calculated by using the software MATLAB, and T = 1.

Example 1. We consider the nonlinear time fractional ODEs as follows:

$$D_*^{\alpha}u - (u^2 - u) = g, t \in (0, T],$$
(22)

where g(t) satisfies the exact solution $u = t + t^{\alpha}$, $0 < \alpha < 1$.

We solve Equation (22) by using Schemes I–V. To verify the numerical errors and convergence orders, we use the temporal step sizes $d_s = 1/1000, 1/2000, 1/3000, 1/4000$ with different α in the first-order scheme. The results presented in Tables 1 and 2 indicate that the convergence order was one, which coincided with our theoretical results. In Tables 3–5, we give the maximum error and the convergence orders for the Newton iterative and Newton linearized methods and the extrapolation skills, respectively. In all these cases, the time steps chosen were $d_s = 1/100, 1/200, 1/400, 1/800$, and we can see from Table 3 that the convergence order was two, which is consistent with the theoretical findings. The results illustrated in Tables 4 and 5 indicate that the numerical experiment performed better than the theoretical conclusions. Here, we also compared our methods with some classical ones, and the results in Table 6 indicate that these methods' orders were α . Moreover, we present the evolution of the maximum norm of the error, and the results found in Figure 1 indicate that our method performed well at the beginning.

	lpha=0.4		$\alpha = 0.6$		$\alpha = 0.8$	
N	Errors	Orders	Errors	Orders	Errors	Orders
1000	3.4639×10^{-1}	*	6.7655×10^{-3}	*	$2.2863 imes 10^{-3}$	*
2000	$1.3343 imes10^{-1}$	1.3763	3.3359×10^{-3}	1.0201	$1.8669 imes 10^{-3}$	1.0731
3000	8.3082×10^{-2}	1.1685	$2.2133 imes10^{-3}$	1.012	$7.0564 imes10^{-4}$	1.0649
4000	$6.0391 imes 10^{-2}$	1.1088	1.6559×10^{-3}	1.0086	5.2011×10^{-4}	1.0643

Table 1. Errors and orders in temporal direction for Scheme I (Example 1).

 Table 2. Errors and orders in temporal direction for Scheme II (Example 1).

	$\alpha = 0.4$		$\alpha = 0$	$\alpha = 0.6$.8
N	Errors	Orders	Errors	Orders	Errors	Orders
1000	4.3014×10^{-1}	*	2.7604×10^{-2}	*	7.1714×10^{-3}	*
2000	2.7749×10^{-1}	0.6324	1.4006×10^{-2}	0.9788	$3.5936 imes 10^{-3}$	0.9968
3000	$2.0582 imes10^{-1}$	0.7368	$9.3831 imes10^{-3}$	0.9879	$2.3972 imes10^{-3}$	0.9985
4000	1.6379×10^{-1}	0.7941	1.6559×10^{-3}	1.0086	$1.7983 imes 10^{-3}$	0.9991

Table 3. Errors and orders in temporal direction for Scheme III (Example 1).

	lpha=0.4		$\alpha = 0$	$\alpha = 0.6$.8
N	Errors	Orders	Errors	Orders	Errors	Orders
100	1.8761×10^{-2}	*	7.2753×10^{-5}	*	2.6841×10^{-5}	*
200	$3.7660 imes10^{-3}$	2.3167	$1.8112 imes10^{-5}$	2.0060	$6.6615 imes10^{-6}$	2.0105
400	$9.0495 imes10^{-4}$	2.0571	4.5335×10^{-6}	1.9983	$1.6600 imes 10^{-6}$	2.0046
800	2.2606×10^{-4}	2.0012	1.1354×10^{-6}	1.9974	$4.1439 imes 10^{-6}$	2.0021

Table 4. Errors and orders in temporal direction for Scheme IV (Example 1).

	lpha=0.4		$\alpha = 0$	$\alpha = 0.6$		$\alpha = 0.8$	
N	Errors	Orders	Errors	Orders	Errors	Orders	
100	1.5948×10^{-1}	*	$2.2079 imes 10^{-3}$	*	2.4690×10^{-4}	*	
200	3.1524×10^{-2}	2.3388	$3.6106 imes10^{-4}$	2.6123	$3.8513 imes10^{-5}$	2.6805	
400	$5.7871 imes10^{-3}$	2.4455	$5.8291 imes10^{-5}$	2.6309	$6.2528 imes10^{-6}$	2.6228	
800	1.0401×10^{-3}	2.4761	$9.2534 imes10^{-6}$	2.6552	1.0753×10^{-6}	2.5998	

Table 5. Errors and orders in temporal direction for Scheme V (Example 1).

	$\alpha = 0.4$		$\alpha = 0$	$\alpha = 0.6$		$\alpha = 0.8$	
N	Errors	Orders	Errors	Orders	Errors	Orders	
100	$8.7281 imes 10^{-2}$	*	$1.9757 imes 10^{-3}$	*	2.0034×10^{-4}	*	
200	2.7388×10^{-2}	1.6721	$3.2896 imes10^{-4}$	2.5864	$3.1694 imes10^{-5}$	2.6602	
400	$6.2592 imes 10^{-3}$	2.1295	$5.3119 imes10^{-5}$	2.6302	$5.2614 imes10^{-6}$	2.5907	
800	1.2079×10^{-3}	2.3735	$8.3960 imes 10^{-6}$	2.6615	9.3199×10^{-7}	2.4971	

		$\alpha = 0.4$		$\alpha = 0.0$	6	$\alpha = 0.8$	
	N	Errors	Orders	Errors	Orders	Errors	Orders
Euler	1000	5.8753×10^{-3}	*	2.5839×10^{-3}	*	5.9704×10^{-4}	*
	2000	4.9689×10^{-3}	0.2418	1.7779×10^{-3}	0.5394	3.5089×10^{-4}	0.7668
	4000	$4.0802 imes10^{-3}$	0.2843	$1.2407 imes10^{-3}$	0.5614	$2.0410 imes10^{-4}$	0.7817
	8000	3.2811×10^{-3}	0.3145	8.0860×10^{-4}	0.5752	1.1806×10^{-4}	0.7898
L1	1000	$8.3161 imes 10^{-3}$	*	$2.8395 imes 10^{-3}$	*	6.2471×10^{-4}	*
	2000	$6.4074 imes10^{-3}$	0.3762	$1.8880 imes 10^{-3}$	0.5888	3.5977×10^{-4}	0.7961
	4000	$4.9231 imes10^{-3}$	0.3802	$1.2522 imes10^{-3}$	0.5924	2.0696×10^{-4}	0.7977
	8000	$3.7727 imes 10^{-3}$	0.3840	8.2909×10^{-4}	0.5949	1.1897×10^{-4}	0.7987
CQBDI	F1 10	$6.6267 imes 10^{-3}$	*	$1.5370 imes 10^{-3}$	*	1.8332×10^{-4}	*
	20	$5.1092 imes10^{-3}$	0.3752	$1.0393 imes10^{-3}$	0.5646	$1.1208 imes10^{-4}$	0.7099
	40	$3.9323 imes10^{-3}$	0.3810	$6.9808 imes10^{-4}$	0.5741	6.7741×10^{-4}	0.7264
	80	3.0031×10^{-3}	0.3856	$4.6667 imes 10^{-4}$	0.5810	$4.0584 imes10^{-5}$	0.7391

Table 6. Errors and orders in temporal direction.



Figure 1. Evolution of maximum errors for different methods.

Example 2. We consider the nonlinear time fractional Allen–Cahn equation

$$D_*^{\alpha}u - u_{xx} - (u - u^3) = g_t(x, t) \in \Omega \times (0, T],$$
(23)

where $\Omega = [0, \pi]$ and g(x, t) satisfies the exact solution is $u(x, t) = (t + t^{\alpha}) \sin x$.

Similarly, we solve the time-fractional Allen–Cahn equation by using Scheme III based on a variable transform. We take M = 1000 with N = 8, 16, 32, 64 to find the maximum of the errors and orders in the temporal direction. Moreover, we consider different spatial step sizes $d_x = \frac{\pi}{M}$, where M = 8, 16, 32, 64 with N = 1000 for different α . The numerical errors are shown in Tables 7 and 8, respectively, where it can clearly be seen that the convergence orders in the temporal and spatial directions are both two.

Table 7. Errors and orders in temporal direction for Scheme III (Example 2).

	$\alpha = 0.4$		$\alpha = 0$	$\alpha = 0.6$		$\alpha = 0.8$	
N	Errors	Orders	Errors	Orders	Errors	Orders	
8	1.5343×10^{-3}	*	3.4241×10^{-4}	*	7.3544×10^{-4}	*	
16	$4.3389 imes10^{-4}$	1.8222	$8.7884 imes10^{-5}$	1.9459	$1.6378 imes10^{-4}$	2.1669	
32	1.1709×10^{-4}	1.8897	$2.2818 imes10^{-5}$	1.9616	3.8450×10^{-5}	2.0907	
64	$3.0965 imes 10^{-5}$	1.9189	5.8943×10^{-6}	1.9527	9.2644×10^{-6}	2.0532	

	$\alpha = 0.4$		$\alpha = 0$	$\alpha = 0.6$.8
M	Errors	Orders	Errors	Orders	Errors	Orders
8	$3.4258 imes 10^{-3}$	*	$3.3368 imes 10^{-3}$	*	$3.2843 imes 10^{-3}$	*
16	$8.6220 imes10^{-4}$	1.9903	8.4019×10^{-4}	1.9897	8.2728×10^{-4}	1.9891
32	$2.1600 imes10^{-4}$	1.9970	$2.1043 imes10^{-4}$	1.9974	$2.0720 imes10^{-4}$	1.9974
64	5.4114×10^{-5}	1.9969	5.2643×10^{-5}	1.9990	5.1812×10^{-5}	1.9997

Table 8. Numerical spatial accuracy of Scheme III for different α (Example 2).

4. Conclusions

In this paper, we presented a family of transformed finite difference methods for numerically solving fractional differential equations while taking the initial singularities of the solutions into account. The convergence results were obtained by using a fractional Grönwall-type inequality. The numerical results were given to illustrate the theoretical results.

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Abbreviations

The following abbreviations are used in this manuscript:

- MDPI Multidisciplinary Digital Publishing Institute
- DOAJ Directory of Open Access Journals
- TLA Three-letter acronym
- LD Linear dichroism

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