

## SUPPLEMENTARY MATERIALS

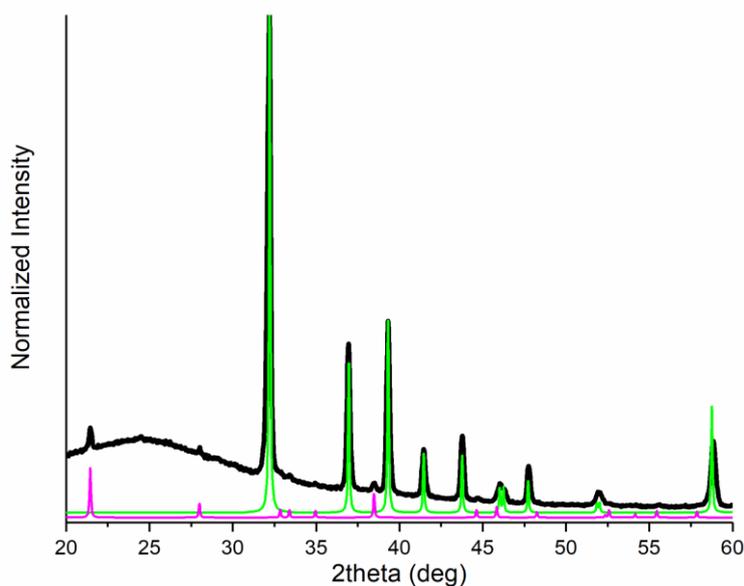
# The Jahn-Teller Distortion at High Pressure: The Case of Copper Difluoride

Dominik Kurzydłowski

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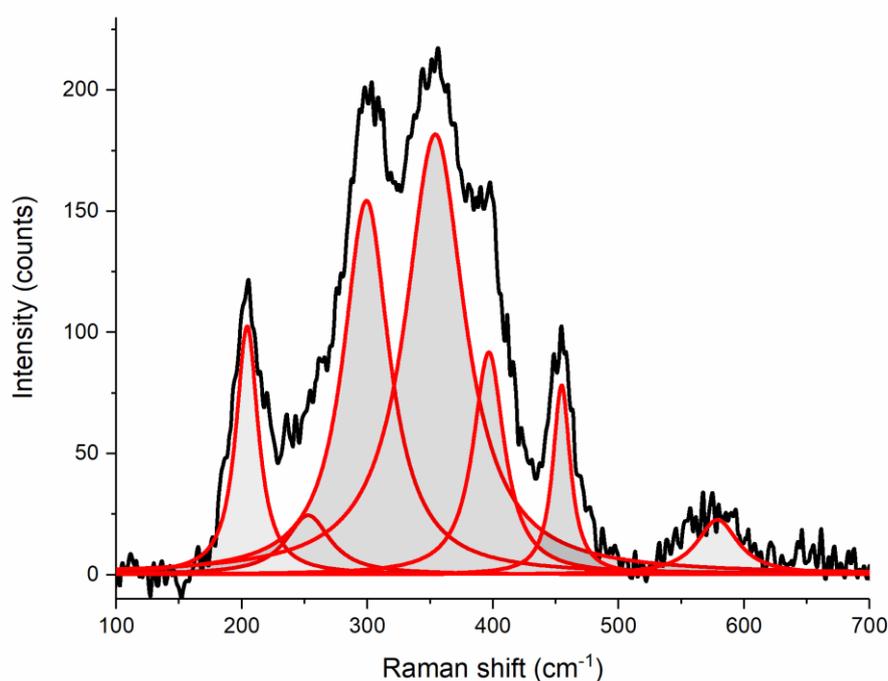
## 1. Ambient-Pressure X-Ray Diffraction

The crystallinity and purity of  $\text{CuF}_2$  was investigated with the use of powder X-ray diffraction. Figure S1 shows a powder diffractogram obtained with the X'Pert Pro diffractometer (Panalytical, Almeo, Holand) using a cobalt X-ray source ( $I(\text{CoK}_{\alpha 1})/I(\text{CoK}_{\alpha 2}) = 2$ ; angle range from  $10^\circ$  to  $60^\circ$ ; acquisition time 5 h). As can be clearly seen the sample is mainly composed of  $\text{CuF}_2$  with only a minor contamination (less than 7 vol%) identified as hydrated copper difluoride ( $\text{CuF}_2 \cdot 2\text{H}_2\text{O}$ ).



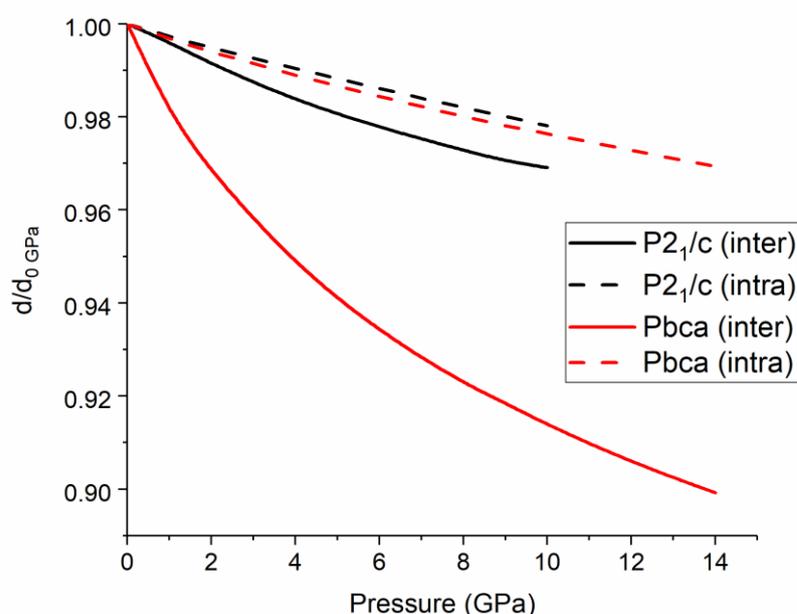
**Figure S1** The experimental powder X-ray diffraction pattern of a sample of  $\text{CuF}_2$  (black line) together with patterns simulated for the  $\text{CuF}_2$  crystal (green) and  $\text{CuF}_2 \cdot 2\text{H}_2\text{O}$  (magenta).

## 2. Deconvolution of the Raman Spectrum of $\text{CuF}_2$ at 19.6 GPa



**Figure S2** The experimental Raman spectrum of  $\text{CuF}_2$  (black line) together with the deconvolution into Lorentzian profiles (red lines).

### 3. Inter-Sheet vs. Intra-Sheet Compressibility of $P2_1/c$ and $Pbca$



**Figure S3** Comparison of the compressibility of inter-sheet (full lines) and intra-sheet (dashed lines) Ag-Ag distances in  $P2_1/c$  (black) and  $Pbca$  (red).

Note that from Figure S3 it is evident that  $Pbca$  exhibits much softer behavior of the inter-sheet contacts compared with  $P2_1/c$ , which is in line with its smaller bulk modulus.

### 4. Comparison of Raman-Active Modes of $ZnF_2$ and $CuF_2$

Zinc difluoride ( $ZnF_2$ ) adopts the undistorted rutile structure ( $P4_2/mnm$ ) which exhibits four Raman-active modes:  $B_{1g}$ ,  $E_g$ ,  $A_{1g}$ , and  $B_{2g}$  (Table S1). Upon the symmetry lowering to the  $P2_1/c$  space group (adopted by  $CuF_2$ ) the  $E_g$  mode splits into two non-degenerate modes ( $B_g$  and  $A_g$ ), while the silent  $A_{2g}$  mode transforms to a Raman-active  $B_g$  mode (Table S1). As a result  $CuF_2$  exhibits six Raman-active modes.

**Table S1.** Comparison of  $\Gamma$ -point Raman-active modes of  $ZnF_2$  and  $CuF_2$ . Experimental frequencies obtained at ambient pressure ( $\omega_0$ ) are given in  $cm^{-1}$ .

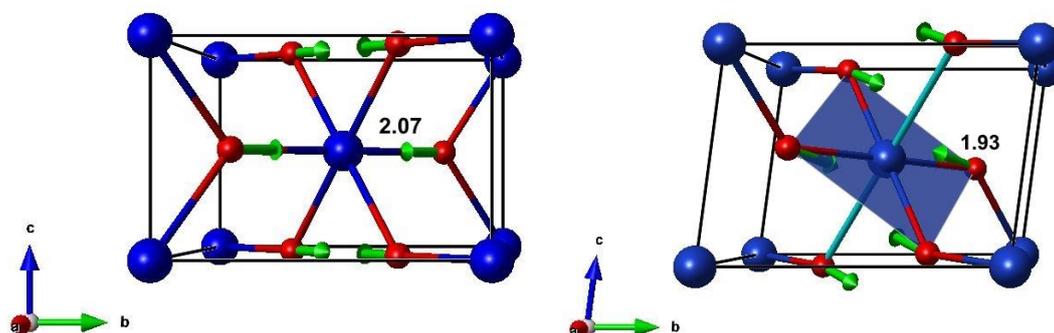
$ZnF_2$			$CuF_2$	
$\omega_0^1$	Symmetry		Symmetry	$\omega_0$
69	$B_{1g}$	$\rightarrow$	$A_g$	n.d.
240	$E_g$	$\rightarrow$	$B_g$	221
		$\rightarrow$	$A_g$	254
n.d. <sup>2</sup>	$A_{2g}^2$	$\rightarrow$	$B_g$	293
347	$A_{1g}$	$\rightarrow$	$A_g$	355
523	$B_{2g}$	$\rightarrow$	$B_g$	566

<sup>1</sup> Data from ref [24]; <sup>2</sup> silent mode.

The largest differences between frequency of symmetry-related modes is found for the highest energy  $B_g$  mode of  $CuF_2$  which is shifted up by  $43\text{ cm}^{-1}$  compared to the  $B_{2g}$  mode of  $ZnF_2$ . One might speculate that this difference is a result of the JT distortion of the  $Cu^{2+}$  octahedron.

The comparison of the two modes indicates that the both involve stretching of a pair of Zn-F/ $Cu$ -F bonds (Figure S4). For the  $CuF_2$  this is a pair of equatorial bonds which are JT-shortened compared

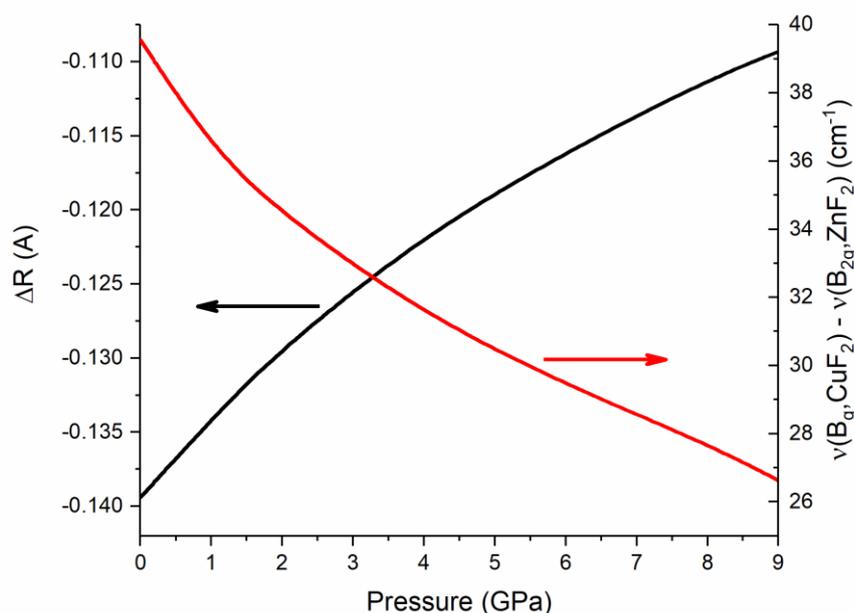
to analogous bonds found in  $\text{ZnF}_2$  by about  $0.14 \text{ \AA}$  at ambient pressure. Thus, as expected, the higher frequency of the  $B_g$  mode stems from the distortion induced by the JT effect.



**Figure S4** Comparison of the eigenvectors of the  $B_{2g}$  mode of  $\text{ZnF}_2$  (**left**), and the symmetry-related  $B_g$  mode of  $\text{CuF}_2$  (**right**). The calculated distances of the Zn-F/Cu-F bonds which undergo stretching along these modes are given in  $\text{\AA}$ .

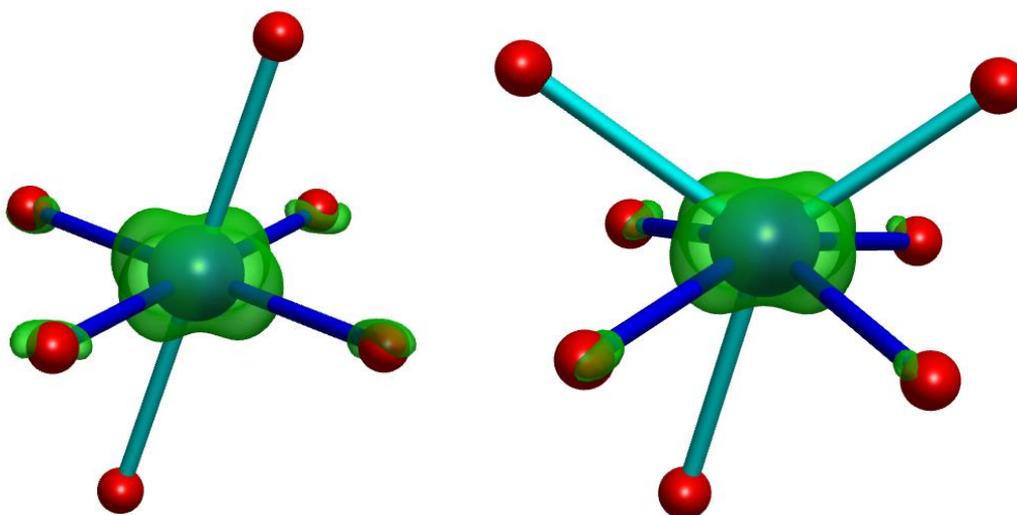
Comparison of the calculated pressure evolution of the abovementioned bonds indicates that upon compression the difference between these bonds decreases to about  $0.11 \text{ \AA}$  at 9 GPa (Figure S5). This is in line with the reduction of the JT effect in the  $P2_1/c$  structure of  $\text{CuF}_2$  at higher pressure (see Figure 7a).

This reduction should lead to a decrease in the difference between the frequencies of the  $B_g$  and  $B_{2g}$  modes. Unfortunately the pressure dependence of the  $B_{2g}$  mode of  $\text{ZnF}_2$  has not been established experimentally, therefore, we compared those calculated theoretically (Figure S5). As can be seen the difference becomes smaller upon compression ( $40 \text{ cm}^{-1}$  at 0 GPa to  $27 \text{ cm}^{-1}$  at 9 GPa).



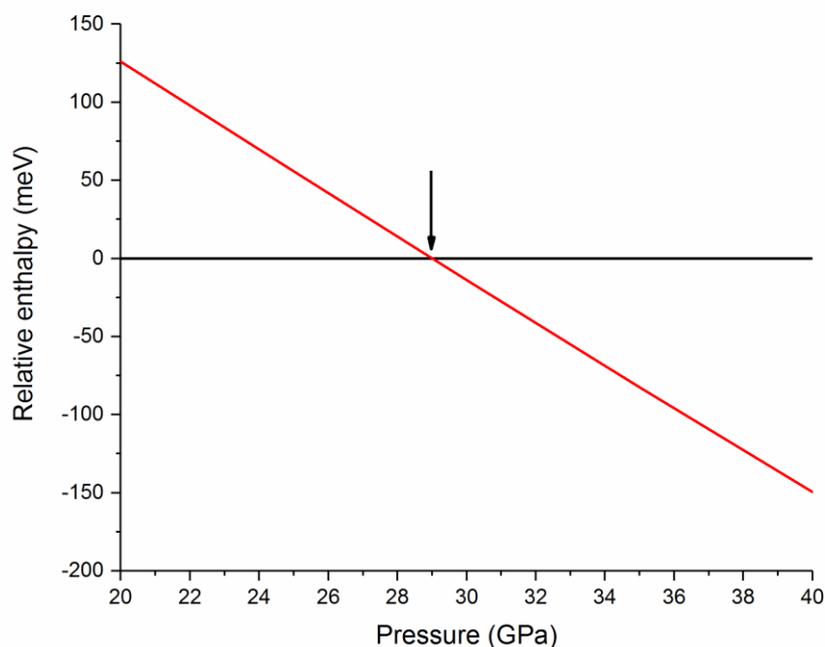
**Figure S5** Calculated pressure evolution of the difference between the Zn-F/Cu-F bonds highlighted in Figure S4 (black line) together with the predicted differences in the frequencies of the highest  $B_g$  mode of  $\text{CuF}_2$  and the  $B_{2g}$  mode of  $\text{ZnF}_2$  (red line).

### 5. Spin Density in *Pbca* at 30 GPa and *Pnma* at 100 GPa



**Figure S6** Comparison of the spin-density (green surface) calculated for *Pbca* at 30 GPa (**left**) and *Pnma* at 100 GPa (**right**). The isosurface is drawn at  $0.2 \text{ \AA}^{-3}$ .

### 6. Relative Enthalpies of the Cotunnite and HP-PdF<sub>2</sub> Phases of ZnF<sub>2</sub>



**Figure S7** The pressure dependence of the relative enthalpy of the cotunnite phase of ZnF<sub>2</sub> (red line) referenced to that of the HP-PdF<sub>2</sub> phase (black line). The arrow marks a predicted phase transition at 29 GPa.

