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On Plowing Frictional Behavior during Scratch Testing: A Comparison between Experimental and Theoretical/Numerical Results

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Received: 26 November 2018; Accepted: 8 January 2019; Published: 11 January 2019



Abstract: Scratch testing is a contact mechanics based nondestructive testing method that, if correctly evaluated, can give a lot of information about the material and tribological behavior of a material. In contrast to the situation with another contact-based method, indentation testing, wear characteristics can also be investigated, for example. In order to get results of practical importance from a scratch test, it is necessary to have evaluation formulae available. Indeed, such formulae exist for scratch testing but can be substantially influenced by frictional effects. For this reason, closed-form analytical relations have been suggested for the purpose of accounting for such effects during scratching and in particular the plowing frictional effect. As a major benefit, these relations can also be of assistance during material characterization through scratch testing. However, the proposed existing relations are based solely on theoretical/numerical analyses and, remembering that the scratch test of course is an experimental approach, verification by experiments is a necessity. Such a task is performed in the present study and it is shown that, based on standard contact global properties, the relations are accurate for most polymeric materials but could also be used for some metallic ones.

Keywords: scratch testing; frictional behavior; plowing friction; mechanical analysis; polymeric materials

1. Introduction

Scratch testing, as well as normal indentation, is frequently used for characterizing materials constitutively, in particular for thin film materials and at small length scales. When it comes to analyses of these two contact mechanics based methods, normal indentation is by far the most used. For the purpose of material characterization, the analyses by Tabor [1] and Johnson [2,3] are classical works devoted to classical elastic-plastic materials. In these studies global indentation properties, i.e., normal hardness, H_n (average contact pressure), normal indentation load, P_n , and indentation depth curve, h , are related to elastic-plastic material properties. It should also be mentioned that another important contact parameter is the actual area of contact necessary in order to calculate the hardness.

The analysis by Johnson [2,3] defined the fundamentals for a general analysis of indentation testing. For sharp indentation on classical Mises elastoplastic materials, this general correlation parameter, schematically shown in Figure 1, could be defined:

$$\Lambda = E \tan \beta / ((1 - \nu^2) \sigma_{\text{rep}}) \quad (1)$$

In Equation (1), E is the Young's modulus, ν is Poisson's ratio, β is the angle between the (sharp) indenter and the undeformed surface, and σ_{rep} is the material flow stress at a representative value

of the effective (accumulated) plastic strain ε_p . Based on the parameter Λ , three regions or levels of contact behavior can be defined. These levels, as shown in Figure 1, are:

- (1) Level I: Dominating elastic deformations, i.e., low indentation load, where elastic contact analysis is sufficient.
- (2) Level II: Elastic and plastic deformations are of equal magnitude.
- (3) Level III: Plastic deformations are dominant in the contact region.

In level II, it was determined by Johnson [2,3] that in this region

$$H = \ln \Lambda, \quad (2)$$

while in level III, Tabor [1] showed that at a sharp indentation

$$H = C\sigma_{rep}. \quad (3)$$

For a Vickers indenter $C = 3$ and $\varepsilon_p \approx 0.08$ (Tabor [1]) and for a cone indenter $C = 2.55$ and $\varepsilon_p = 0.11$ (Atkins and Tabor [4]). The latter case is specified for a cone indenter with an angle of 22° between the indenter and the undeformed surface.

In a general situation, the simple formulae above might not be sufficient for high-accuracy results. This concerns, for example, the case when the material strain-hardening cannot be described by a regular power law hardening—cf. [5–7], where advanced finite element methods were utilized. Indeed, remembering the complexity of the indentation problem finite element methods is an essential part of the solution procedure, as specified in, for example, [5–12].

Finite element methods are obviously also a very important part of an analysis of scratch testing (visualized in Figure 2). This has been performed in [13–21] (just to mention a few). An important conclusion from these results is that the Johnson parameter [2,3], in Equation (1), also governed correlation during scratch problems. This was further clarified in [22,23], where explicit formulae for this correlation were presented for a number of global scratch quantities, accounting for the difference in scratch behavior at elastoplastic and rigid plastic scratching. One of these global scratch quantities investigated was the so-called apparent coefficient of friction, μ_0 , and this quantity will be further investigated presently. It goes almost without saying that experimental verification of these scratch relations is of greatest importance and this will be attempted here for the apparent coefficient of friction with predictions from [23]. These predictions will be compared to corresponding experimental results originally presented by Wredenberg and Larsson [19]. The experimental results are pertinent to both polymeric and metallic materials and, accordingly, a large range of values on the Λ -parameter in Equation (1) are covered. Both the theoretical/numerical and experimental results are valid for cone scratching with $\beta = 22^\circ$; see Figure 2 for the indenter geometry, and classical elastoplasticity is relied upon as a foundation for the constitutive behavior.

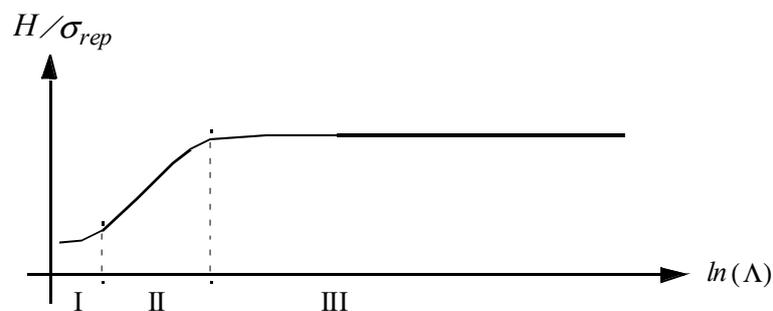


Figure 1. Sketch of the characteristic behavior at indentation [4]. The indentation hardness H divided by the representative stress σ_{rep} is plotted against the non-dimensional strain parameter Λ . The different levels at indentation are indicated.

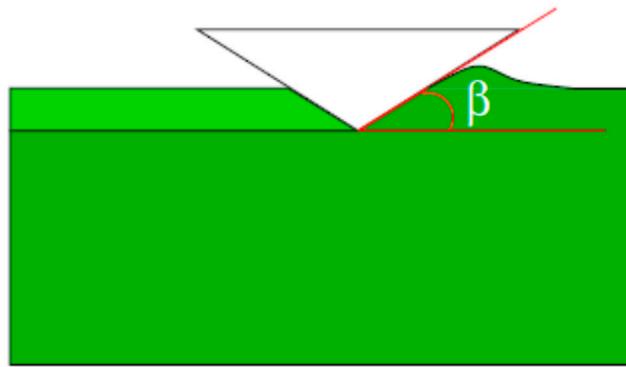


Figure 2. Schematic of the scratching problem. The angle β defines the geometry of the conical stylus, where $\beta = 22^\circ$. The scratch direction is from left to right.

2. Theoretical Background

The problem at issue concerns scratching using a sharp rigid conical stylus with $\beta = 22^\circ$ in Figure 2. The specimen is considered large and composed of a monolithic material. In the presentation below h is the normal indentation depth. Quasi-static and steady-state conditions, as well as classical elastoplasticity, are assumed and based on these restrictions; the problem is self-similar with no characteristic length present. Consequently, global indentation properties are independent of indentation depth h during the loading sequence of a scratch test on homogeneous materials. Here and in the rest of the paper the apparent coefficient of friction is defined as, cf. e.g., [19],

$$\mu_0 = \mu_i + \mu_p. \quad (4)$$

The decomposition in Equation (4) between an adhesive/interfacial part and a plowing part was originally suggested by Bowden and Tabor [24] but has also been verified numerically by Felder and Bucaille [15]. Furthermore, the plowing part of the apparent coefficient of friction depends on the mechanics of the scratch problem and it is possible to write

$$\mu_p = \mu_p(\Lambda). \quad (5)$$

From analytical considerations, cf. [25,26], it is possible to derive approximate estimates on $\mu_p(\Lambda)$. In short, it can be shown that at ideal rigid plastic behavior,

$$\mu_p(\Lambda) = 0.26 \quad (6)$$

at cone scratching with $\beta = 22^\circ$.

In order to arrive at a more complete description of the Λ -dependence of μ_p also including strain-hardening plasticity, Larsson [23] studied the scratch problem using, as mentioned above, finite element methods. Regarding the constitutive details in [23], classical rate-independent elastoplasticity, accounting for large deformations, is assumed to prevail and the strain hardening behavior was based on a power law assumption according to

$$\sigma(\varepsilon_p) = \sigma_Y + \sigma_0(\varepsilon_p)^n, \quad (7)$$

where $\sigma(\varepsilon_p)$ is the flow stress, σ_Y is the initial yield stress, ε_p is, as previously stated, the effective plastic strain, and n is the hardening exponent. The representative stress value is then defined as

$$\sigma_{\text{rep}} = \sigma(\varepsilon_{\text{rep}}), \quad (8)$$

where ε_{rep} is the effective strain at a representative value of the plastic strain.

Based on the results from finite element simulations, and within the basic assumptions discussed above, Larsson [23] suggested that the plowing part of the apparent coefficient of friction is very well correlated with the Johnson parameter Λ according to

$$\begin{aligned}\mu_p(\Lambda) &= 0.03 + 0.072\ln\Lambda \text{ (Level II, } \Lambda \leq 20 \text{ approximately)} \\ \mu_p(\Lambda) &= 0.26 \text{ (Level III, } \Lambda \geq 50 \text{ approximately)},\end{aligned}\quad (9)$$

with the choice of representative stress in Equation (1) being

$$\begin{aligned}\sigma_{\text{rep}} &= \sigma(\varepsilon_p = 0.25) \text{ (Level II, } \Lambda \leq 20 \text{ approximately)} \\ \sigma_{\text{rep}} &= \sigma(\varepsilon_p = 0.39) \text{ (Level III, } \Lambda \geq 25 \text{ approximately)}.\end{aligned}\quad (10)$$

In the analysis leading to Equations (9) and (10), frictional effects are not included in the numerical approach. This is due to the decomposition suggested by Bowden and Tabor [24], explicitly presented in Equation (4), indicating that interfacial and plowing frictional effects are separable. Accordingly, μ_p is the same for any value on μ_i and performing the numerical simulations with $\mu_i = 0$ is then the obvious choice. Based on this discussion, it goes almost without saying that experimental verification of Equations (9) and (10) is crucial. This will also be attempted below, pertinent to materials described in the next section.

3. Experiments and Materials

In this section the experimental cone scratching results in [19] are detailed. Included in this discussion is also a description of the metallic and polymeric materials at issue in the present investigation.

In [19] cone scratching with an indenter angle $\beta = 22^\circ$ (see Figure 2) was performed on six different materials including three polymers and three metals. A diamond indenter was used in order to secure as much as possible a non-deforming stylus. The normal and tangential forces were recorded and, among other things, normal and scratch hardness were recorded.

It should be immediately stated that one of the materials, Aluminum 4120, showed anomalous behavior due to the formation of oxide on the aluminum surface. For that reason, this material is left out of the present comparison.

The material characteristics are shown in Table 1. The strain hardening behavior of the materials is specified in [19] and defined here by the values of the two representative stresses as specified in Equation (10).

Consequently, the experimental scratch results, pertinent to the plowing part of the apparent coefficient of friction, for the five materials are then compared with the predictions from Equations (9) and (10). This comparison will be presented and discussed below. It should be emphasized that the experimental scratch hardness values, presented and discussed below, are averages from the experiments and standard deviations are not shown for brevity. The latter feature is, however, clearly presented in [19].

Table 1. Material properties as presented by Wredenberg and Larsson [19] and utilized in the present analysis. The value of Young's modulus is given in GPa, while all stress quantities are given in MPa. The parameter Λ , as defined in Equation (1), is dimensionless.

Material	E	σ_Y	$\sigma(\varepsilon_p = 0.25)$	$\sigma(\varepsilon_p = 0.39)$	$\Lambda(\varepsilon_p = 0.25)$	$\Lambda(\varepsilon_p = 0.39)$
Vinyl Ester	3.5	108	108	108	15	15
PMMA	2.9	110	110	110	12	12
Epoxy	3.1	97	97	97	14	14
Al7050	70	560	724	816	42.5	38
Stainless steel	200	560	625	805	146	113

4. Results and Discussion

The theoretical/numerical predictions from [23], as detailed in Equations (9) and (10), will be compared with corresponding experimental results presented in [19]. In doing so, it seems appropriate to describe the theoretical/numerical results leading to the relations in Equation (9). This is done in Figures 3 and 4. In doing so, the representative stress measures in the figures are always the ones presented in Equation (10).

In Figure 3, relevant for level II contact, and Figure 4, relevant for level III contact, the results for the plowing part of the apparent coefficient of friction, μ_p , are shown as first presented in [23]. Results for different values of the power law exponent n , in Equation (7), are depicted together with the predictions from Equation (9). According to Equation (10), then, in Figure 3, $\sigma(\varepsilon_p = 0.25)$ is used as a representative stress and, correspondingly in Figure 4, $\sigma(\varepsilon_p = 0.39)$ is used. Clearly, except for very low values of $\ln\Lambda$ where the elastic level I contact effects are intervening, and at the transition between level II and level III contact, the different curves come together as a single universal curve with good accuracy. This universal curve is then represented by Equation (9), with the representative stresses as detailed in Equation (10). It should be noted in passing that the numerical results show good agreement with Equation (6) in the level III contact regime.

The main purpose of the present study is to determine if the universal results in Figures 3 and 4, and specifically the relations presented in Equations (9) and (10), are sufficient to predict experimental results for the materials presented and described in Table 1. Consequently, the representative stresses determined previously and presented in Equation (10) are used to determine the experimental values on the plowing part of the apparent coefficient of friction, μ_p , for the materials in Table 1 and then compared with the predictions from Equations (9) and (10). The outcome of this comparison is detailed in Figures 5 and 6.

To start, in Figure 5 results for the plowing part of the apparent coefficient of friction are shown for level II contact conditions. This corresponds to the upper Equation (9), which is identified as the continuous straight line in the figure. Corresponding experimental results for the three polymers in Table 1 are also shown. Clearly, the agreement between predictions and experiments is very good except for the case of the polymeric material PMMA, where the predictions are somewhat high but within the standard deviation.

Corresponding results for level III rigid plastic contact are shown in Figure 6. Also, in this case the agreement between predictions and experimental results is good even though the explicit value on μ_p for Al7050 is somewhat high compared with the predictions. The difference can, however, be explained by experimental uncertainties but also due to the fact that the Λ -values for Al7050 are close to the transition zone between level II and level III contact. Indeed, if the upper relation in Equation (9) is used to predict the experimental results for Al7050 (instead of the lower relation in Equation (9) as in Figure 6) then almost perfect agreement is found between predictions and experimental results (in both cases $\mu_p = 0.3$ is found). This is not unusual behavior, though, when contact mechanics testing is correlated using the Johnson parameter Λ defined in Equation (1). Indeed, this feature is also present at straightforward normal indentation testing where hardness (normalized using a representative stress) values overshoot in the transition zone and then reach a constant value as level III is properly achieved, cf. [12]. This is the same behavior indicated by the experimental results for Al7050 in Figure 6. Obviously, this suggests that special care has to be taken when using the relations in Equation (9) at transitional contact (i.e., at transition between level II and level III contact); this would be a very interesting subject for further studies.

The results shown in Figures 5 and 6 are encouraging as they suggest that the relations shown in Equations (9) and (10) are sufficient (in most cases) for accurate prediction of the plowing part of the apparent coefficient of friction. The results also indicate that the plowing frictional behavior is well described from an analysis assuming no interfacial friction.

Also, thin film scratching has been studied a lot recently, cf. e.g., [27–38], and the investigations concern aspects such as fracture, delamination behavior (as described in [39,40]), the behavior of different global and local scratch variables, and alternative constitutive behavior. Even though thin films introduce a characteristic length into the problem, the outcome of the present study can be used at least as a guideline to analyze the mechanics of thin film scratching.

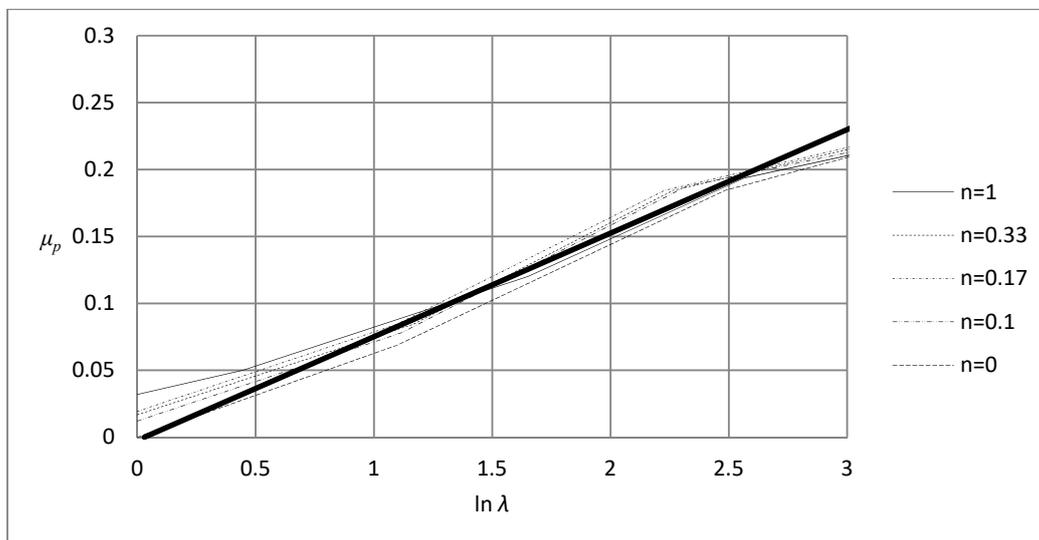


Figure 3. Plowing part of the apparent coefficient of friction, μ_p , as a function of $\ln \Lambda$, where Λ is the Johnson parameter defined in Equation (1). The representative stress is determined from $\sigma_{\text{rep}} = \sigma(\varepsilon_p = 0.25)$. The results are taken from Larsson [23]. The thick line represents upper Equation (9).

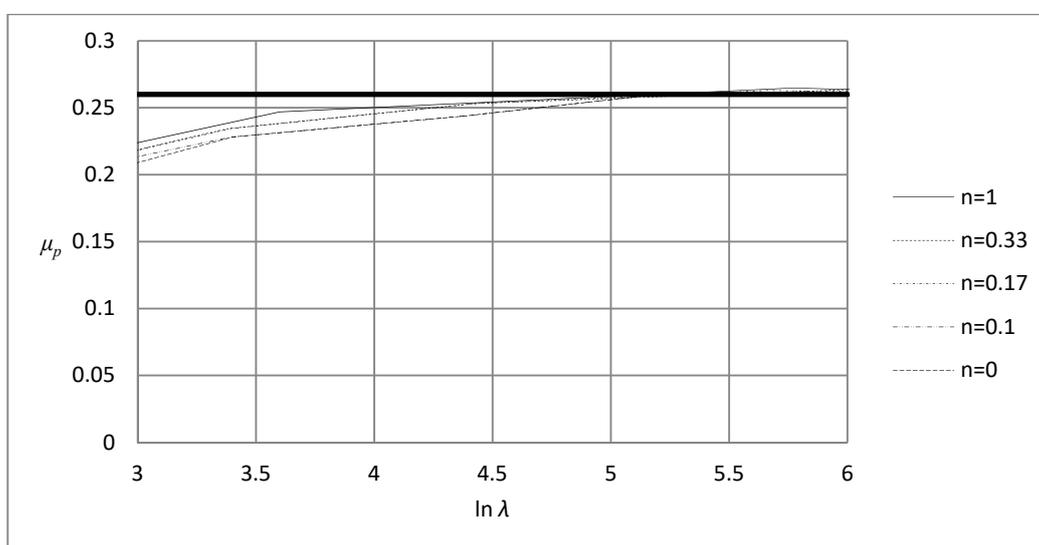


Figure 4. Plowing part of the apparent coefficient of friction, μ_p , as a function of $\ln \Lambda$, where Λ is the Johnson parameter defined in Equation (1). The representative stress is determined from $\sigma_{\text{rep}} = \sigma(\varepsilon_p = 0.39)$. The results are taken from Larsson [23]. The thick line represents lower Equation (9).

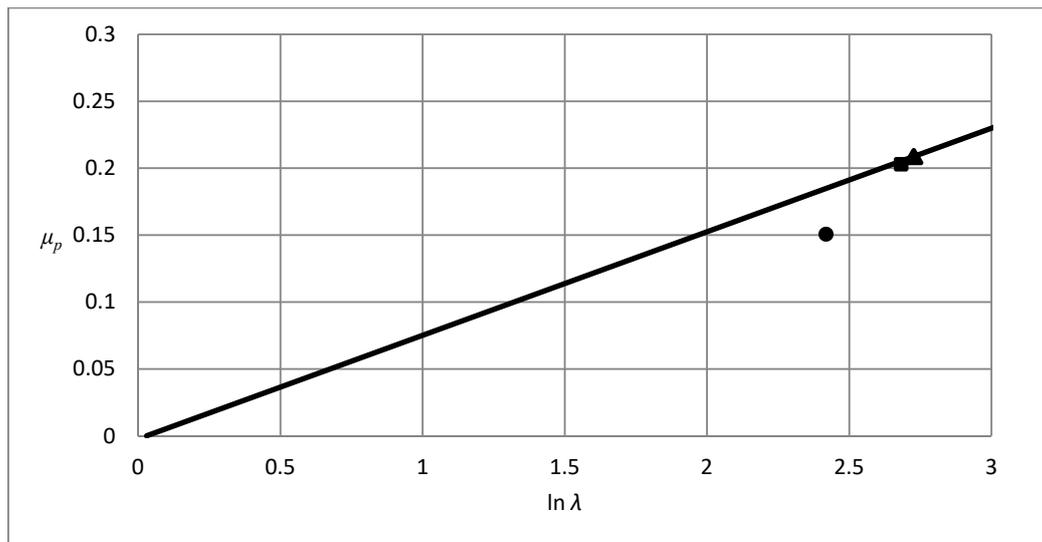


Figure 5. Plowing part of the apparent coefficient of friction, μ_p , as a function of $\ln \Lambda$, where Λ is the Johnson parameter defined in Equation (1). The representative stress is determined from $\sigma_{rep} = \sigma(\epsilon_p = 0.25)$. The thick line represents upper Equation (9). (●), experimental results for PMMA as detailed in Table 1. Results taken from Wredenberg and Larsson [19]. (■), experimental results for epoxy as detailed in Table 1. Results taken from Wredenberg and Larsson [19]. (▲), experimental results for vinyl ester as detailed in Table 1. Results taken from Wredenberg and Larsson [19]. The standard deviations are 0.035 (PMMA), 0.033 (epoxy) and 0.021 (vinyl ester).

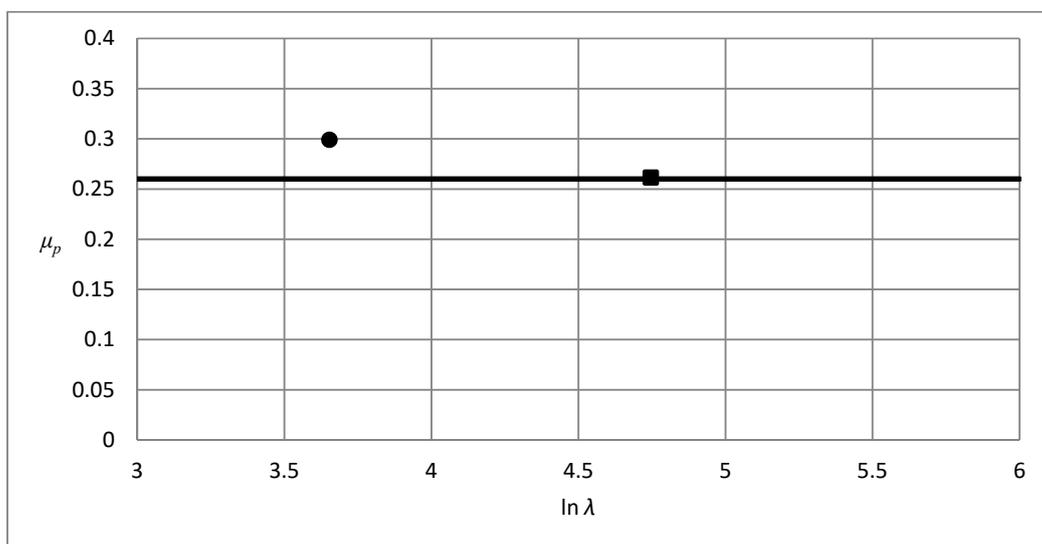


Figure 6. Plowing part of the apparent coefficient of friction, μ_p , as a function of $\ln \Lambda$ where Λ is the Johnson parameter defined in Equation (1). The representative stress is determined from $\sigma_{rep} = \sigma(\epsilon_p = 0.39)$. The thick line represents lower Equation (9). (●), experimental results for Al7050, as detailed in Table 1. Results taken from Wredenberg and Larsson [19]. (■), experimental results for stainless steel as detailed in Table 1. Results taken from Wredenberg and Larsson [19]. The standard deviations are 0.009 (Al7050) and 0.009 (stainless steel).

In the author's opinion, the most significant finding is the tools to predict (and analyze) the frictional behavior at scratch testing in both the elastoplastic and the rigid plastic contact regime. While this feature was previously rather well known in the case of metals, the present results also make it possible to account for the frictional effects of scratching for polymers. The applications of scratch testing could then include:

- Material characterization (together with similar relations for global scratch properties).
- Determination of the interfacial friction coefficient.
- Determination of the cracking behavior at scratching based on FEM calculations where the frictional behavior is properly accounted for.

5. Conclusions

Correlation of global quantities at scratch testing has been studied. In particular, this concerns the plowing part of the apparent coefficient of friction. It has been shown previously that within the framework of classical Mises elastoplasticity and quasi-static conditions, correlation can be achieved by using a combination of stresses at different levels of plastic strains to define representative quantities at scratching, accounting for the difference in mechanical behavior at elastoplastic and rigid plastic scratching. This correlation is detailed in explicit relations and these relations are compared to the experimental results for a wide range of materials ranging from polymers to metals. In general, good agreement is found between the two sets of results.

Funding: This research was funded by Swedish Research Council grant number 621-2005-5803

Conflicts of Interest: The author declares no conflict of interest

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