

## Article

# Inkjet-Printed Interdigitated Capacitors for Sensing Applications: Temperature-Dependent Electrical Characterization at Cryogenic Temperatures down to 20 K

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**Abstract:** Microwave transducers are widely used for sensing applications in areas such as gas sensing and microfluidics. Inkjet printing technology has been proposed as a promising method for fabricating such devices due to its capability to produce complex patterns and geometries with high precision. In this work, the temperature-dependent electrical properties of an inkjet-printed single-port interdigitated capacitor (IDC) were investigated at cryogenic temperatures down to 20 K. The IDC was designed and fabricated using inkjet printing technology, while its reflection coefficient was measured using a vector network analyzer in a cryogenic measurement setup and then transformed into the corresponding admittance. The resonant frequency and quality factor (Q-factor) of the IDC were extracted as functions of the temperature and their sensitivity was evaluated. The results showed that the resonant frequency shifted to higher frequencies as the temperature was reduced, while the Q-factor increased as the temperature decreased. The trends and observations in the temperature-dependent electrical properties of the IDC are discussed and analyzed in this paper, and are expected to be useful in future advancement of the design and optimization of inkjet-printed microwave transducers for sensing applications and cryogenic electronics.

**Keywords:** inkjet printing; interdigitated capacitor; microwave measurements; material characterization; cryogenic temperatures



**Citation:** Gugliandolo, G.; Alimenti, A.; Latino, M.; Crupi, G.; Torokhtii, K.; Silva, E.; Donato, N.

Inkjet-Printed Interdigitated

Capacitors for Sensing Applications:

Temperature-Dependent Electrical

Characterization at Cryogenic

Temperatures down to 20 K.

*Instruments* **2023**, *7*, 20. [https://](https://doi.org/10.3390/instruments7030020)

[doi.org/10.3390/instruments7030020](https://doi.org/10.3390/instruments7030020)

Academic Editor: Antonio Ereditato

Received: 17 May 2023

Revised: 16 July 2023

Accepted: 17 July 2023

Published: 19 July 2023



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## 1. Introduction

In recent years, microwave transducers have gained widespread use in various fields, including sensing applications [1–3], dielectric materials characterization [4–6], microfluidics [7–9], and metrological applications [10–12]. The high sensitivity and fast response time of these devices make them suitable for measuring small changes in physical, chemical, and biological parameters. Microwave transducers designed in microstrip or coplanar technology have been widely adopted due to their features of easy fabrication, low cost, and low power requirements [13]. These devices inherit compatibility with wireless technology, as they operate in radio frequency bands, and offer the advantage of contactless measurements, which are highly desirable or even essential in various sensing applications [14–16]. Various geometries for microwave transducers have been proposed in the literature, such as disk resonators [17], ring resonators [18], split-ring resonators [19], spirals [1], rectangular patches [20], and interdigitated capacitors (IDCs) [3], with each geometry offering specific advantages for particular applications. Among the different geometries, IDCs have been widely studied and employed in sensing applications from low frequencies to the microwave range [3,21–23].

The concept of interdigital electrodes was first introduced by Nikola Tesla in 1891 [24], when he used parallel rectangular plates known as fingers to increase the total capacitance of an electrical condenser. This principle is used today in the fabrication of modern capacitors [25]. The theoretical expression for the estimation of the capacitance between coplanar strips was introduced in 1924 [26]. However, it was not until the 1960s that interdigital electrodes started to be extensively used for sensing applications [25]. Their simple design and ease of fabrication using such different technologies as screen printing, computer numerical control (CNC) micromachining, and inkjet printing, are among of the key advantages of IDCs.

The small size of IDCs makes them suitable for sensing applications and integration into microfluidic devices. In microfluidic systems, IDCs can be used as sensing elements to monitor the conductivity of the liquid under test or its dielectric properties, detecting and quantifying chemical and biological species [21,22,27,28]. These features make IDCs suitable for medical diagnostics and lab-on-a-chip applications [8,29].

In gas sensing, IDCs are frequently combined with a sensing material with dielectric properties that change in the presence of the target gas [3]. The surface of the IDC is coated with a sensing layer that reacts with the gas, causing changes in the electrical properties of the material, such as its effective complex permittivity. These changes can then be detected by the IDC as a variation in the total equivalent capacitance and this value can be related to the gas target concentration through a proper calibration procedure. IDCs have been employed in the microwave range to develop gas sensors for various gases, including oxygen [3], ammonia [30], and ethanol [31], among others.

Inkjet printing technology has emerged as a cost-effective and high-resolution alternative to more traditional fabrication techniques. It offers a contactless, maskless, and flexible approach to the deposition of conductive ink on substrates. Inkjet printing has been used in a wide range of applications, such as flexible electronics [32,33], on-paper devices [34], and microwave transducers [35]; however, one of the major limitations of this technology is related to the properties of the ink and how they can vary with time or operating conditions. For instance, the performance of inkjet-printed devices can be affected by factors such as the relative humidity or temperature [36], which can result in changes to the ink properties and lead to decreased device performance. Despite this issue, few studies have been conducted on the performance of inkjet-printed microwave transducers under extreme conditions, particularly at cryogenic temperatures. In this context, studying the behavior of inkjet-printed microwave transducers over a range of operating temperatures, particularly at cryogenic temperatures, can provide valuable insights into the performance and reliability of such devices. Moreover, a thorough understanding of the temperature-dependent electrical properties of microwave transducers can be helpful for optimization of the design and performance of high-performance microwave sensors.

In this study, we have focused on investigating the temperature-dependent electrical properties of an inkjet-printed interdigitated capacitor over a wide temperature range from 20 K to 300 K. The substrate material used for fabrication of the IDC was glass-reinforced epoxy laminate (FR4), which was characterized in terms of its electromagnetic properties. The obtained characterization results were used for the design of the prototype by employing computer-aided design (CAD) simulation software. The IDC prototype was fabricated by printing it onto the FR4 substrate using a Voltera V-One printer along with a commercial conductive ink, the Conductor 2 ink produced by Voltera [37]. A single-port IDC was designed using coplanar technology and its electrical properties were evaluated using a vector network analyzer (VNA), for which the device was placed inside a cryogenic measurement system. The aim of the study was to provide a more comprehensive understanding of the electrical behavior of inkjet-printed IDCs at cryogenic temperatures.

The rest of the article is organized into four sections: Section 2 below comprises three subsections detailing the characterization of the FR4 material, the design and fabrication of the IDC prototype, and the cryogenic test setup; Section 3 outlines the main results obtained

from the tests; then the article concludes with Section 4, which provides a summary of the main findings and their significance.

## 2. Materials and Methods

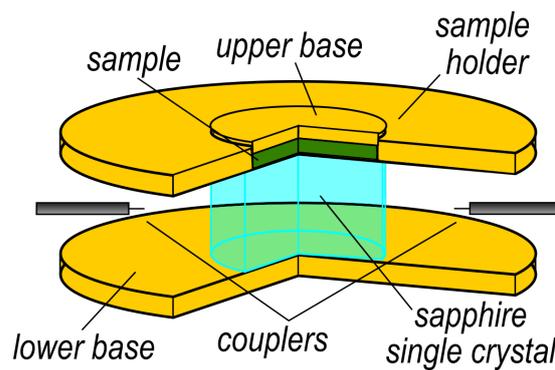
This section of the paper provides a detailed account of the materials and methods employed to investigate the temperature-dependent electrical properties of the inkjet-printed interdigitated capacitor. The section is divided into three subsections, each of which outlines a specific aspect of the investigation. The first subsection focuses on the characterization of the FR4 material used as a substrate for the IDC under test. The dielectric constant and loss tangent of the FR4 material were determined through experimental measurements, and were subsequently used in the device design process. This process is discussed in the second subsection, which provides a comprehensive description of the IDC prototype, including its geometric parameters and fabrication process through inkjet printing technique. The design of the IDC prototype was carried out using Matlab software; the details of this design process are included in this subsection as well. The third and final subsection provides a detailed description of the cryogenic measurement system used to perform the temperature-dependent electrical measurements of the IDC prototype. The measurement system was designed to operate over a wide temperature range and equipped with a VNA to measure the electrical properties of the IDC at different operating conditions.

### 2.1. Material Characterization

The initial step in the IDC design process involves the characterization of the electromagnetic properties of the substrate material being used. In this case, it was a commercial FR4 provided by the Voltera company. The dielectric constant values of FR4 substrates differ significantly depending on the material's anisotropy and resin content [38]. Therefore, direct measurement is necessary for an appropriate device simulation. The response of a dielectric material to electromagnetic fields is defined by the complex dielectric permittivity  $\epsilon_0\tilde{\epsilon} = \epsilon_0(\epsilon' - i\epsilon'')$ , where  $\epsilon_0$  is the dielectric constant of vacuum  $\epsilon' = \text{Re}(\tilde{\epsilon})$ ,  $\epsilon'' = -\text{Im}(\tilde{\epsilon})$ , the ratio  $\tan \delta = \epsilon''/\epsilon'$  is known as the loss tangent, and  $i = \sqrt{-1}$  [39]. Therefore, in order to design the resonator geometry parameters accurately it is essential to measure the complex dielectric permittivity of the FR4 substrate used for the IDC substrate.

Microwave measurement techniques employed for material characterization can be categorized into two main groups [40], resonant and non-resonant. The former is well known for its high sensitivity, while the latter, although possessing lower sensitivity and accuracy, provides the significant benefit of being wideband [40]. At the operational frequencies of the designed IDC, the dielectric characteristics of the FR4 are only slightly frequency-dependent [41]; thus, the wideband nature of transmission/reflection measurement techniques does not offer a significant advantage. In contrast, the high sensitivity of resonant methods can be particularly useful for achieving accurate IDC design. Therefore, a dielectric-loaded resonator was exploited to characterize the FR4 substrate, as shown in [6,42] and depicted in Figure 1.

The measurement technique used in this work is known as the volume perturbation technique [40]. In this method, a portion of the internal volume of an electromagnetic resonant structure is first replaced with a reference material and then with the material being investigated. By examining the variation in the quality factor  $Q$  and resonance frequency  $f_0$ , it is possible to estimate the electromagnetic properties of the unknown material [40]. Air is utilized as the reference material here due to its dielectric constant being approximately  $\tilde{\epsilon} \approx 1 - i0$ . It has been demonstrated in previous studies [6,40,42] that  $\Delta f_0$  is proportional to  $\epsilon'$ , while  $\Delta(Q^{-1})$  is proportional to the loss tangent  $\tan \delta$ . Here,  $\Delta x$  indicates the variation in the parameter  $x$  with respect to a reference value (in this case, air).



**Figure 1.** Sketch of the microwave dielectric-loaded resonator used to measure the dielectric constant of the FR4 substrate. The lower and upper bases are made of brass, as is the sample holder. The dielectric crystal employed in the measurement is a sapphire cylinder of 5 mm height and 8 mm diameter. The  $TE_{011}$  mode is generated by means of coaxial cables that are ended with magnetic loops and operate at approximately 12.5 GHz. The dielectric sample is loaded beneath the upper base, and the resonator is closed by placing a 500 g weight on it. For further details, please refer to [42].

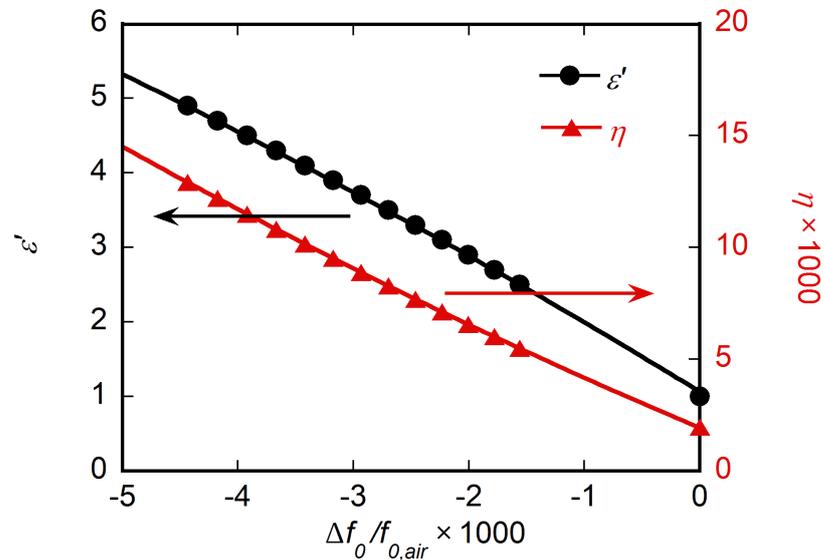
To avoid a significant reduction in  $Q$  and resulting loss of sensitivity due to expected large losses of FR4, only a small volume of the resonator was replaced with this material. The optimization of the sample volume with respect to its expected electromagnetic properties is discussed in detail in [42]. In this study, an FR4 fragment with dimensions  $15.0 \times 15.0 \times 1.6 \text{ mm}^3$  was loaded into the resonator, as shown in Figure 1. For evaluation of the measurement repeatability, the reflection and transmission scattering ( $S$ -) parameters of the resonator were measured twenty times, with the sample being dismounted for each measurement.  $Q$  and  $f_0$  were determined by fitting the acquired  $S$ -parameters using the method described in [43]. It was determined that a single measurement per mounting was sufficient, as a statistical analysis of multiple measurements indicated that the uncertainties on  $Q$  and  $f_0$  were comparable to those obtained through the fitting procedure and analyzing the fit residuals from a single measurement [43].

To obtain the reference measurement, we replaced the same volume of the FR4 sample with air. In this configuration, a ring with the same thickness as the sample under investigation was printed and inserted into the resonator instead of the sample itself. This allowed the resonator geometry to be maintained while introducing a controlled volume of air. Using the FR4 sample, a quality factor of  $Q_{\text{sample}} = 2413.8(17)$  and a resonance frequency of  $f_{0,\text{sample}} = 12.381122(21)$  GHz were obtained, where the values in parentheses represent the numerical value of the experimental standard deviations,  $s(Q_{\text{sample}})$  and  $s(f_{0,\text{sample}})$ , respectively, for the corresponding last digits of the reported results. The twenty repeated measurements with the air reference sample resulted in a quality factor of  $Q_{\text{air}} = 4998(13)$  and a resonance frequency of  $f_{0,\text{air}} = 12.434475(15)$  GHz.

The relative change in resonance frequency ( $\Delta f_0/f_{0,\text{air}} = (f_{0,\text{sample}} - f_{0,\text{air}})/f_{0,\text{air}}$ ) was compared with the calibration curves of dielectric permittivity ( $\epsilon'(\Delta f_0/f_{0,\text{air}})$ ) and  $\eta(\Delta f_0/f_{0,\text{air}})$ , with  $\eta$  being the sample filling factor [40] obtained through electromagnetic simulations of the resonating structure [42]. The calibration curves shown in Figure 2 were used to determine  $\epsilon'$  and  $\eta$ . The loss tangent ( $\tan \delta$ ) was then calculated as  $\tan \delta \approx \Delta(Q^{-1})/\eta$  as described in [42].

The FR4 substrate under analysis yielded  $\epsilon' = 4.75(8)$  and  $\tan \delta = 1.73(6) \times 10^{-2}$ , which are consistent with values reported in the literature [44]. The reported standard uncertainties were evaluated by propagating the distribution of simulated and measured quantities using the Monte Carlo method in compliance with [45,46]. The simulated  $f_0$  and  $\eta$  were assumed to have a normal distribution with 1% relative standard deviation. In each trial of the Monte Carlo simulation, the calibration curves  $\epsilon'(\Delta f_0/f_{0,\text{air}})$  and  $\eta(\Delta f_0/f_{0,\text{air}})$  were re-evaluated and fitted with a second-order polynomial, providing the uncertainties on the fitting parameters reported in the caption of Figure 2. The measured  $f_{0,\text{sample}}$  and

$f_{0,air}$  and their standard deviations allowed  $\epsilon'$  and  $\eta$  (as well as their uncertainties) to be obtained, from which  $\tan \delta$  (and its uncertainty) was calculated using the measured  $Q$  and  $s(Q)$ . In all,  $10^6$  Monte Carlo trials were performed [45].



**Figure 2.** Calibration curves  $\epsilon'(\Delta f_0/f_{0,air})$  (black dots—left scale) and  $\eta(\Delta f_0/f_{0,air})$  (red triangles—right scale), with  $\Delta f_0 = f_{0,sample} - f_{0,air}$ , i.e., the difference between the resonance frequency  $f_{0,sample}$  measured when the sample is loaded into the resonator and that of  $f_{0,air}$  measured with the reference (air). The continuous curves were obtained by a second-order polynomial fit:  $\epsilon' = -2.81(3) \times 10^4 (\Delta f_0/f_{0,air})^2 - 1.002(2) \times 10^3 \Delta f_0/f_{0,air} + 1.002(2)$  with  $1 - R^2 = 2.86 \times 10^{-6}$  and  $\eta = 72.95(7) (\Delta f_0/f_{0,air})^2 - 2.1400(3) \Delta f_0/f_{0,air} + 1.9720(4) \times 10^{-3}$  with  $1 - R^2 = 1.38 \times 10^{-8}$ . The uncertainty bars are within the dimensions of the symbols.

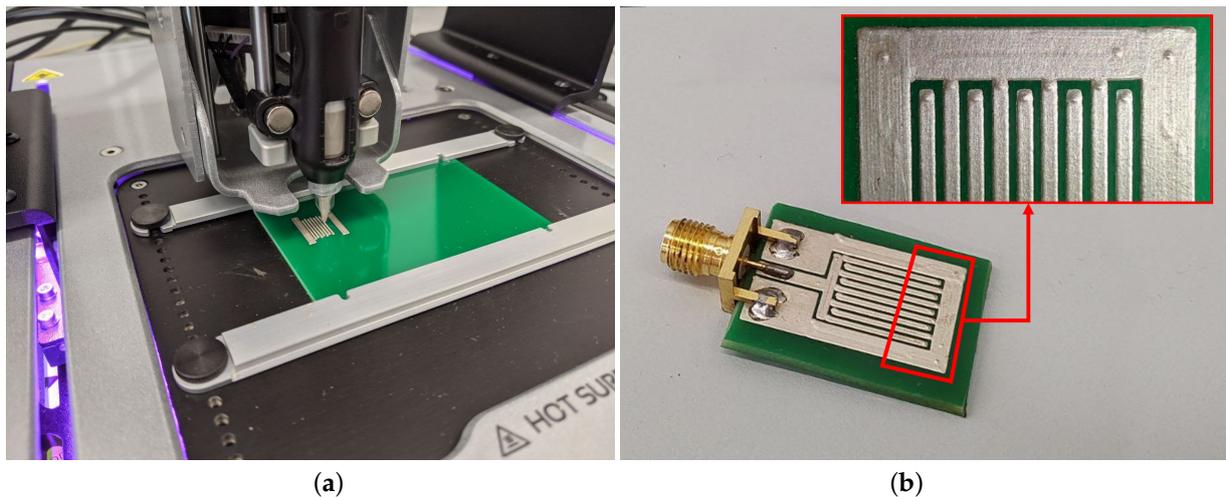
## 2.2. Inkjet Printed Interdigitated Capacitor: Design and Fabrication

The device under investigation is a one-port coplanar interdigitated capacitor. The design was carried out using MATLAB as a computer-aided design tool. In particular, the Antenna toolbox was used for accurate 3D modeling of the device in order to simulate its electrical response over a frequency range spanning from 100 MHz to 6 GHz. The device was made up of a 1.6-mm-thick FR4 substrate surrounded by a conductive layer. The material properties for the FR4 substrate were set according to the results obtained by the material characterization described in the previous subsection ( $\epsilon_r = 4.75$ ;  $\tan \delta = 0.0173$ ), while the conductive layer properties were set according to the nominal specification of the conductive ink used, i.e., a nominal thickness of 50  $\mu\text{m}$  and a nominal resistivity of  $1.265 \times 10^{-7} \Omega\cdot\text{m}$  [37].

A 2d drawing of the device is shown in Figure 3. The length, width, and spacing of the fingers are 11.5 mm, 0.5 mm, and 0.5 mm, respectively, while the coplanar feedline's slot, signal, and finite ground widths are 0.5 mm, 1.5 mm, and 5.5 mm, respectively. The device geometry was selected based on the printer's resolution (about 225  $\mu\text{m}$  for the selected ink) and need to maintain an overall size less than 1 inch square, i.e.,  $2.54 \times 2.54 \text{ mm}^2$ . Moreover, the IDC geometry was tuned by means of computer simulations in order to ensure a first resonance at about 1 GHz for the potential use of the IDC in sensing applications [47].

The Voltera V-One inkjet printing machine was used to fabricate the device. This printer is based on direct ink writing (DIW) dispensing technology with an open-loop ink pressure control. It allows tracks with a minimum width of 0.2 mm to be printed, and its resolution is 10  $\mu\text{m} \times 10 \mu\text{m} \times 1 \mu\text{m}$ , for the x, y, and z axes, respectively. A silver-based conductive ink, "Conductor 2" by Voltera [37], was deposited on the FR4 substrate to create the IDC pattern selected during the design process. The main properties of the employed ink are summarized in Table 1.

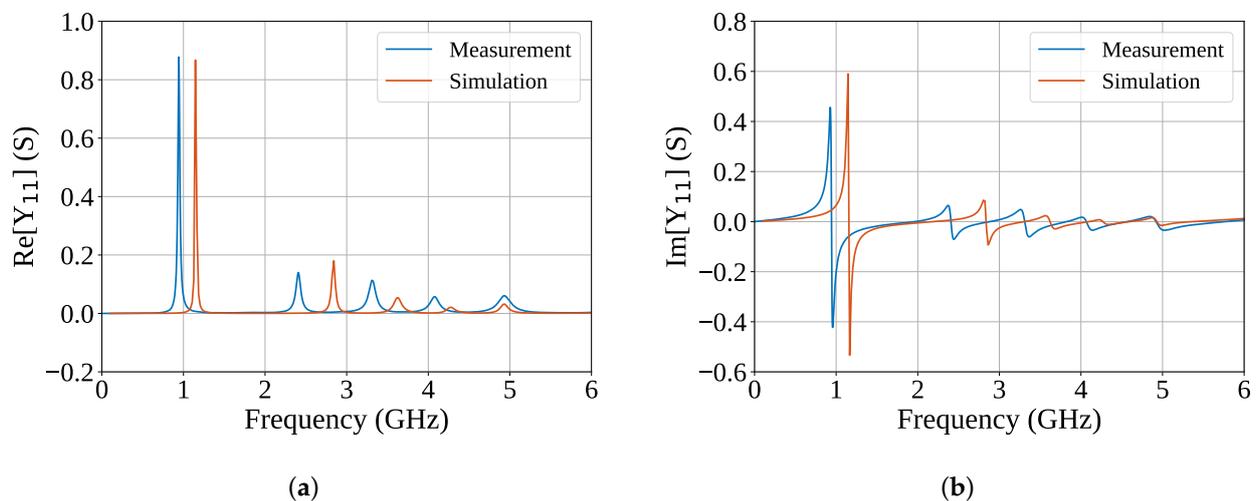




**Figure 4.** Photos of the IDC prototype: (a) Voltera V-One during the printing process and (b) fabricated prototype after curing and SMA connector soldering.

In this work, the device was characterized in terms of the admittance parameter  $Y$ , which was calculated from the reflection coefficient (i.e.,  $S_{11}$ ) measured with the VNA and converted using the appropriate formula [49].

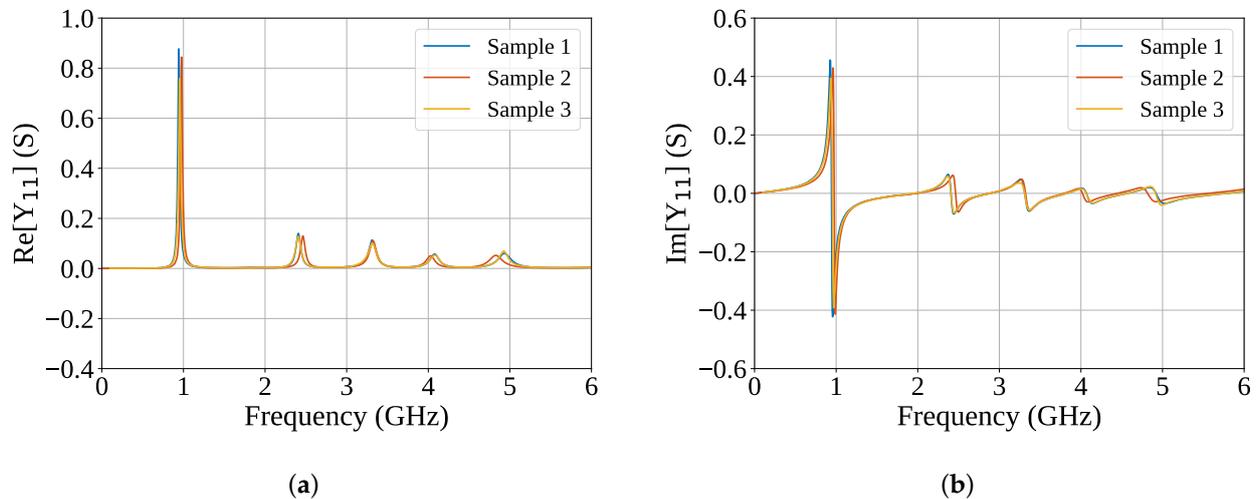
Figure 5 presents a comparison of the simulated and measured admittance comprising both the real and imaginary parts across a frequency range spanning from 100 MHz to 6 GHz. The plot exhibits multiple resonances wherein the real part of the admittance reaches a local maximum while the imaginary part is null. Despite a slight frequency shift of approximately 200 MHz towards the lower frequencies in the fabricated prototype, the measured results are in good agreement with the simulations. This deviation can be attributed to inevitable fabrication tolerances and trace spread after printing.



**Figure 5.** Measurement (blue) and simulation (orange) of the admittance of the IDC in the frequency range from 100 MHz to 6 GHz, depicting (a) real and (b) imaginary parts. The measured  $Y$  is calculated from the IDC reflection coefficient acquired with the VNA in the same frequency range.

To evaluate the reproducibility of IDC fabrication, three samples were printed on an FR4 substrate using the same Voltera V-One printer, each after separately performing the printer calibration procedure. The results in terms of admittance were then compared, as shown in Figure 6. As can be seen, the comparison revealed variations among the samples in the first and second resonance. The reproducibility was evaluated in terms of standard deviation, which was found to be 16 MHz and 33 MHz for the first and second resonances, respectively. These results indicate that the fabrication process was consistent

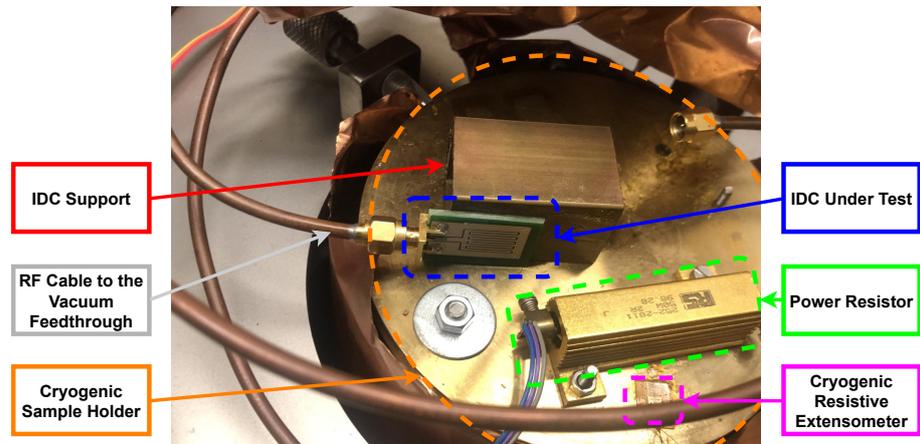
and reproducible, with only relatively small variations observed between the different samples. Furthermore, the results demonstrate the effectiveness of the printer calibration procedure in minimizing variations during the printing process.



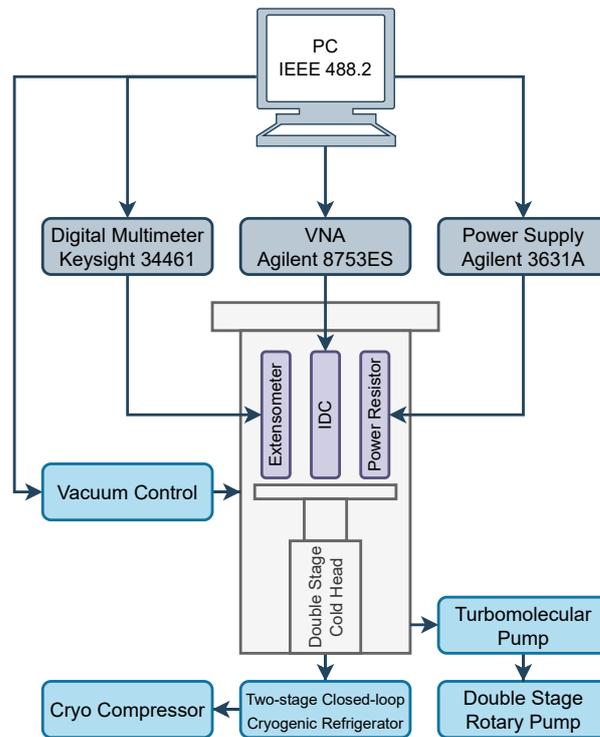
**Figure 6.** Study of the IDC fabrication method's reproducibility. Three samples were compared in terms of the admittance parameter for both the real and imaginary parts, with the respective results shown in (a) and (b).

### 2.3. Experimental Setup

For its frequency- and temperature-dependent electrical characterization, the IDC was placed on a cryogenic thermal chuck inside of a stainless-steel chamber, as shown in Figure 7. The measurement setup included a cryogenic unit, a vacuum system, and an electrical characterization system. The cryogenic unit was used to cool the thermal chuck, acting as a two-stage closed-loop cryogenic refrigerator. The employed cryogenic system was a dual-stage CTI-Cryogenics recirculation helium system with a nominal temperature bottom limit of 10 K. To achieve the desired vacuum level, a Varian vacuum system consisting of a double-stage rotary pump and a turbomolecular pump was used. The vacuum level was monitored through thermocouple-based and cold cathode sensors. Figure 8 shows a schematic representation of the experimental setup. The electrical characterization system included an Agilent 3631A power supply connected to a power resistor that acted as an actuator for the chuck temperature control. A Keysight 34461 digital multimeter was used to measure the resistance of a cryogenic resistive extensometer, allowing the thermal chuck temperature to be evaluated. A VNA Agilent 8753ES was used to measure the IDC reflection coefficient. A frequency range from 100 MHz to 6 GHz was selected in this study, for a total of 1601 acquired points. Prior to the measurement process, a full two-port calibration procedure was performed on the VNA with a commercially available calibration kit using a short-open-load-through (SOLT) procedure. All components of the experimental setup were connected to a desktop computer through the IEEE 488.2 general-purpose interface bus (GPIB). This allowed the configuration of each instrument, as well as the chuck temperature, to be set using custom-made software.



**Figure 7.** Photo of the IDC prototype inside the stainless-steel chamber placed on the cryogenic thermal chuck. The power resistor and cryogenic resistive extensometer used for temperature actuation and control are depicted as well.

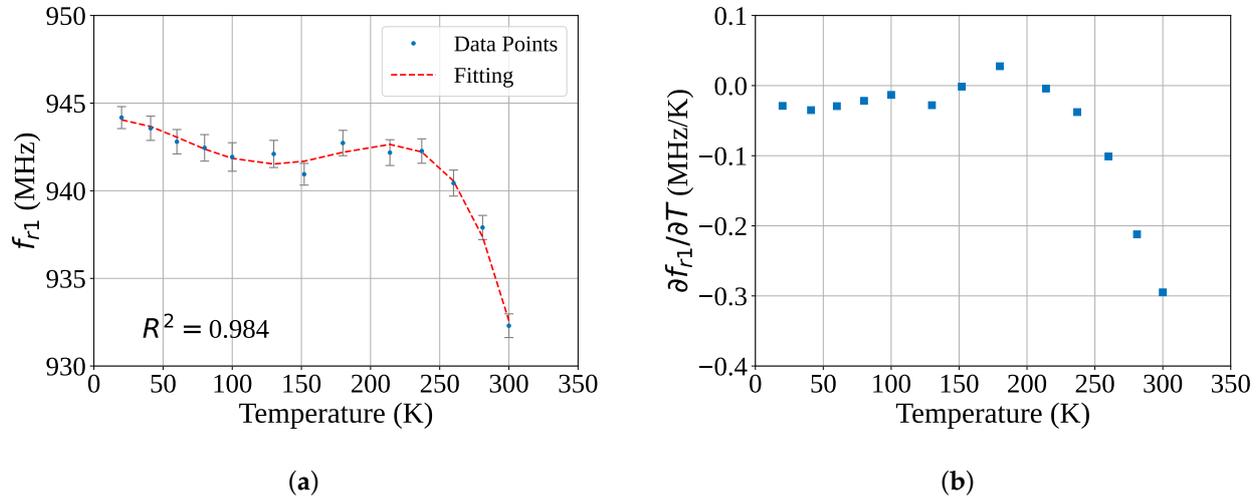


**Figure 8.** Schematic representation of the cryogenic measurement system.

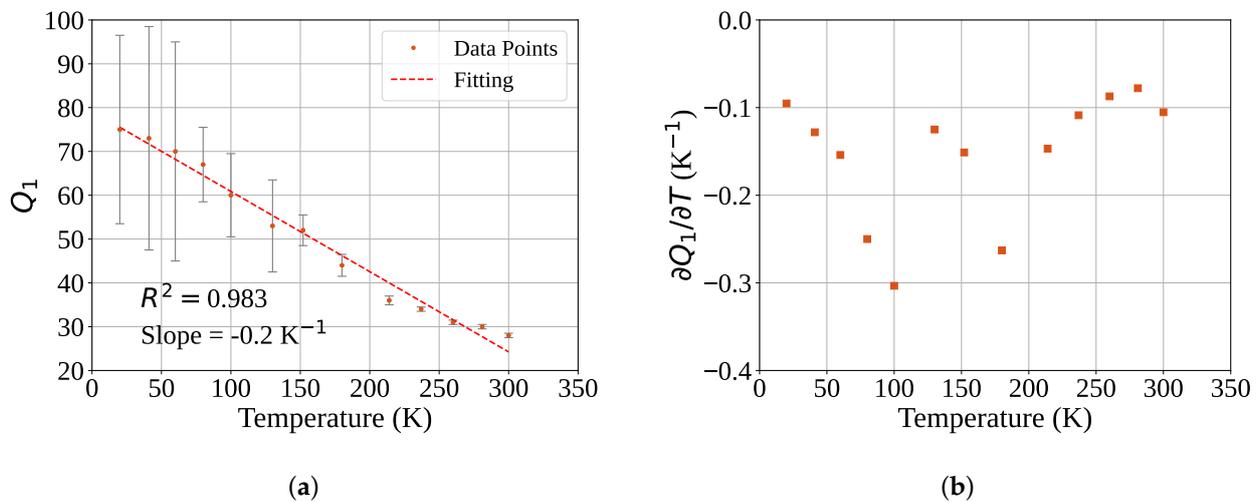
### 3. Results

After the IDC had been placed inside the cryogenic chamber and a high vacuum level achieved, the reflection coefficient was measured from room temperature (i.e., 300 K) down to 20 K. The reflection coefficient was then converted into the corresponding admittance using the conversion formula [49]. This study primarily focuses on the first two resonances in the  $Y$  parameter, i.e., those at 932 MHz and 2.388 GHz. Both resonances are affected by temperature change, with a frequency shift towards lower frequencies observed as the temperature decreases. The peaks sharpness increases as well, corresponding to an increase in the  $Q$ -factor. In order to more accurately estimate the resonant parameters (i.e., the resonant frequency and  $Q$ -factor), a fitting procedure based on a Lorentzian function [50] was applied to the admittance parameter.

The first studied resonance occurred at 932 MHz (at 300 K). Figures 9 and 10 show that the resonant frequency (i.e.,  $f_{r1}$ ) and Q-factor ( $Q_1$ ) are both temperature-dependent. As the temperature decreases, the resonant frequency shifts towards lower values while the Q-factor improves.



**Figure 9.** Temperature dependence of the resonance at 932 MHz: (a) depicts the relationship between the resonant frequency and the temperature; the trend is modeled with a fourth-order polynomial function with a  $R^2 = 0.984$  (the red dashed line), while the sensitivity towards the temperature is shown in (b).



**Figure 10.** Temperature dependence of the resonance at 932 MHz: (a) depicts the relationship between the Q-factor and the temperature; the trend can be considered linear with a good approximation ( $R^2 > 0.98$ ) (the red dashed line), while The sensitivity towards the temperature is shown in (b) and is about  $-0.2 \text{ K}^{-1}$ .

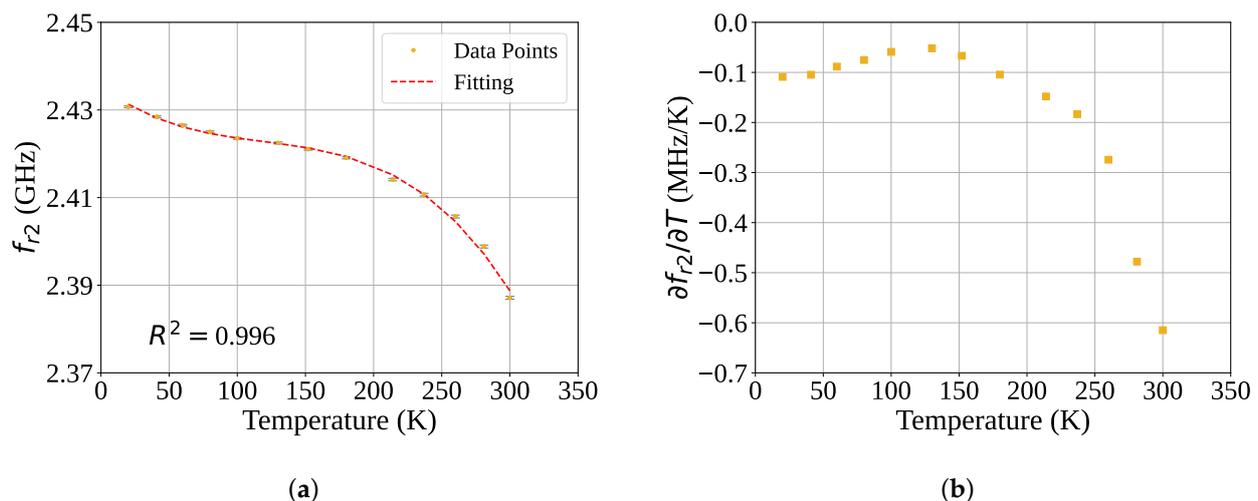
It should be noted that the error bars displayed in the figures correspond to the standard uncertainty of the measurements derived from both the Agilent 8753ES VNA and the fitting procedure. The measurement uncertainty of the VNA was determined using the VNA Uncertainty Calculator software from Keysight. As the temperature decreases, the sharpness of the peaks increases, which can negatively impact the accuracy of the fitting procedure. In fact, when the total number of acquired points is kept constant (at the maximum allowed by the VNA, i.e., 1601) and measurement conditions and VNA settings are both held constant as well, the peaks are inevitably described by a lower number of points due to a higher Q-factor. As a result, wider error bars are obtained.

For this resonance, the resonant frequency value does not exhibit a linear relationship with the temperature. Specifically, a linear trend is observed from 20 K to 100 K, followed by a constant range between 100 K and 250 K, then a sharp drop in resonant frequency for higher temperature values. This trend can be modeled with a fourth-order polynomial function, yielding an  $R^2$  value greater than 0.98. The sensitivity of  $f_{r1}$  with respect to temperature as determined by  $\partial f_{r1}/\partial T$  is shown in Figure 9b. In contrast, the Q-factor shows a more linear trend with temperature ( $R^2$  approximately equal to 0.98), and its values and sensitivity with temperature ( $-0.2 \text{ K}^{-1}$ ) are plotted in Figure 10b.

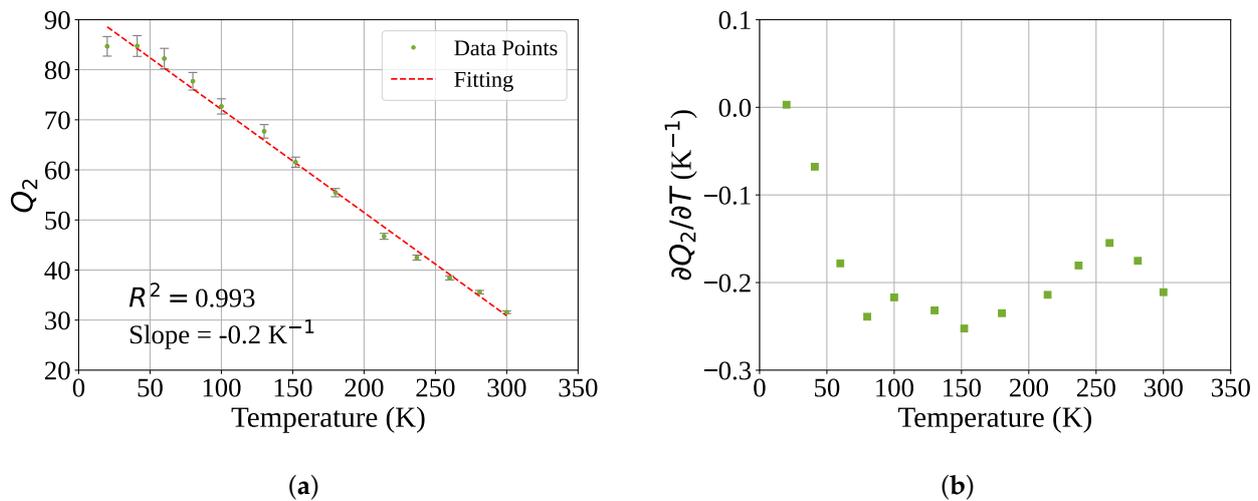
The second resonance, referred to as  $f_{r2}$ , is characterized by a monotonic trend from 2.435 GHz at 20 K to 2.388 GHz at 300 K. The relationship between  $f_{r2}$  and temperature can be accurately modeled with a third-order polynomial function with  $R^2 > 0.99$ . The sensitivity of  $f_{r2}$  to temperature ranges from  $-0.1 \text{ MHz/K}$  at 20 K to  $-0.58 \text{ MHz/K}$  at 300 K. The resonant frequency and its sensitivity to temperature are illustrated in Figure 11. Figure 12 displays the Q-factor of the second resonance as a function of the temperature in the cryogenic chamber. The trend can be approximated as linear with good accuracy ( $R^2 > 0.99$ ), and exhibits a sensitivity of approximately  $-0.2 \text{ K}^{-1}$ .

It can be observed that the resonances shift towards higher values as the temperature decreases towards cryogenic levels. This can be attributed to the thermal expansion effect of both the ink and the substrate, which results in a decrease in their size at lower temperatures, leading to a commensurate increase in the resonant frequency [51]. Another factor that may contribute to the resonant frequency shift is the temperature coefficient of the FR4 dielectric constant. The value of the dielectric constant decreases with lower temperatures [51], resulting in an increase in the resonant frequency. In fact, the IDC's capacitance is proportional to the effective dielectric constant between the fingers. By considering the inductive contribution of the IDC as constant ( $L_{eq}$ ) with the temperature change, the equivalent capacitance ( $C_{eq}$ ) decreases with the FR4 dielectric constant, thereby increasing the resonant frequency; this can be expressed as  $f_r = 1/(2\pi\sqrt{L_{eq}C_{eq}})$  [9].

Furthermore, an increase in the Q-factor is observed at cryogenic temperatures. The Q-factor is generally linked to Ohmic losses and dielectric losses [52]. Therefore, an increase in the Q-factor may be related to a rise in ink conductivity at cryogenic temperatures and/or a decrease in FR4 dielectric losses.



**Figure 11.** Temperature dependence of the resonance at 2.388 GHz: (a) depicts the relationship between the resonant frequency and the temperature; the trend is modeled with a third-order polynomial function with a  $R^2 = 0.996$  (the red dashed line), while the sensitivity towards the temperature is shown in (b).



**Figure 12.** Temperature dependence of the resonance at 2.388 GHz: (a) depicts the relationship between the Q-factor and the temperature; The trend can be considered linear with a good approximation ( $R^2 > 0.98$ ) (the red dashed line), while the sensitivity towards the temperature is shown in (b) and is about  $-0.2 \text{ K}^{-1}$ .

#### 4. Conclusions

The study presented in this article aimed to investigate the frequency-dependent and temperature-dependent electrical properties of an inkjet-printed interdigitated capacitor. After proper characterization of the employed substrate material, a prototype was designed and fabricated by means of an inkjet printer by depositing a silver-based commercial ink on a commercial FR4 substrate. The goal of this research was to gain a better understanding of the behavior of inkjet-printed IDCs at cryogenic temperatures. This is relevant to a variety of different fields, and is especially applicable to sensing applications. In order to achieve this goal, the IDC was placed in a cryogenic chamber and the resonant parameters (i.e., resonant frequencies and Q-factors) were evaluated over a temperature range of 300 K to 20 K. For improved accuracy in the estimation of these parameters, a fitting procedure based on a Lorentzian function was applied. Our experimental results demonstrate that the resonant frequency of the IDC shifted towards higher values as the temperature decreased to cryogenic levels. Additionally, the Q-factor of the IDC increased at cryogenic temperatures. This behavior can be ascribed to changes in the physical properties of both the conductive ink and the substrate as the temperature changes (e.g., thermal expansion, conductivity, and permittivity).

This study investigated the characteristics of an inkjet-printed IDC fabricated on a low-cost FR4 substrate. In future research, it would be valuable to expand this investigation to include inkjet-printed devices using higher-performance substrates more typical for use in the microwave frequency range. This would allow for an assessment of the ink's reliability on different substrates and at different frequency ranges. Additionally, it would be beneficial to conduct material characterization before and after cryogenic testing in order to identify any irreversible changes in the ink material. Finally, future investigations could explore the sensitivity of devices to additional external perturbations such as the presence of contaminants. A particular focus could be the impact of humidity on device performance, which could further enrich understanding of these devices' behavior under different environmental conditions.

**Author Contributions:** Conceptualization, N.D. and E.S.; methodology, N.D.; software, G.G. and K.T.; validation, G.C. and M.L.; formal analysis, G.G., A.A. and K.T.; investigation, G.G., A.A. and K.T.; data curation, G.C. and M.L.; writing—original draft preparation, G.G. and A.A.; writing—review and editing, G.G., A.A., M.L., G.C., K.T., E.S. and N.D.; supervision, E.S. and N.D. All authors have read and agreed to the published version of the manuscript.

**Funding:** This work has been partially funded by European Union (NextGeneration EU), through the MUR-PNRR project SAMOTHRACE (ECS00000022).

**Institutional Review Board Statement:** Not applicable.

**Informed Consent Statement:** Not applicable.

**Data Availability Statement:** Not applicable.

**Conflicts of Interest:** The authors declare no conflict of interest.

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